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# Intricate Reaction Pathways on $\text{CH}_3\text{NH}_3\text{PbI}_3$ Photocatalysts in Aqueous Solution Unraveled by Single-Particle Spectroscopy

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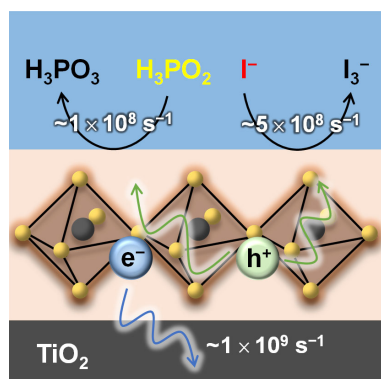
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**ABSTRACT.** Organic–inorganic hybrid perovskites such as MAPbI<sub>3</sub> (MA<sup>+</sup> = CH<sub>3</sub>NH<sub>3</sub><sup>+</sup>) have emerged as promising materials for solar cells and light-emitting devices. Despite their poor stability against moisture, perovskites work as hydrogen-producing photocatalysts or photosensitizers in perovskite-saturated aqueous solutions. However, the fundamental understanding of how chemical species or support materials in the solution affect the dynamics of the photogenerated charges in perovskites is still insufficient. In this study, we investigated the photoluminescence (PL) properties of MAPbI<sub>3</sub> nanoparticles in aqueous media at the single-particle level. A remarkable PL blinking phenomenon, along with significant decreases in the PL intensity and lifetime compared to those in ambient air, suggested temporal fluctuations in the trapping rates of photogenerated holes by chemical species (I<sup>-</sup> and H<sub>3</sub>PO<sub>2</sub>) in the solution. Moreover, electron transfer from the excited MAPbI<sub>3</sub> to Pt-modified TiO<sub>2</sub> proceeds in a concerted fashion for photocatalytic hydrogen evolution under the dynamic solid–solution equilibrium condition.

## TOC GRAPHICS

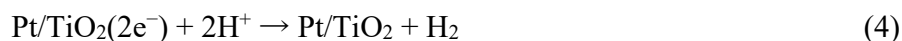
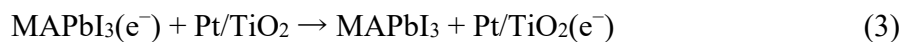
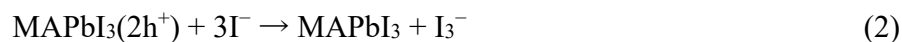


Metal halide perovskites such as MAPbX<sub>3</sub> (MA<sup>+</sup> = CH<sub>3</sub>NH<sub>3</sub><sup>+</sup>, X<sup>-</sup> = Cl<sup>-</sup>, Br<sup>-</sup>, or I<sup>-</sup>) and CsPbX<sub>3</sub> are expected to be utilized in photovoltaic and optoelectronic devices owing to their favorable properties, which include a strong absorption of visible light, long-length carrier diffusion, and a high photoluminescence (PL) quantum yield (QY).<sup>1-5</sup> Recently, these perovskites have been applied to various photochemical reactions.<sup>6</sup> For instance, Xu et al. reported that CsPbBr<sub>3</sub>/graphene oxide (GO) acted as a photocatalyst for CO<sub>2</sub> reduction in ethyl acetate.<sup>7</sup> Yuan et al. found that ligand-capped CsPbBr<sub>3</sub> photocatalyzed stereoselective C–C bond formation under visible light irradiation.<sup>8</sup> Huang et al. reported, for the first time, the oxidation of alcohols to aldehydes using FAPbBr<sub>3</sub> (FA<sup>+</sup> = CH(NH<sub>2</sub>)<sup>+</sup>).<sup>9</sup>

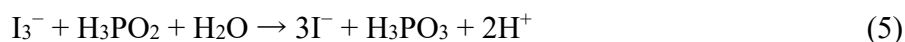
In general, most perovskites decompose easily when in contact with H<sub>2</sub>O or polar solvents.<sup>10,11</sup> However, Park et al. showed that MAPbI<sub>3</sub> is stable in aqueous solution by achieving a dynamic equilibrium between perovskite precipitates and a saturated solution containing their precursors.<sup>12</sup> Under this condition, the perovskites work as hydrogen (H<sub>2</sub>)-producing photocatalysts by themselves (with Pt modification) or as photosensitizers when combined with Pt-modified TiO<sub>2</sub> (Pt/TiO<sub>2</sub>) for H<sub>2</sub> evolution via HX splitting.<sup>13</sup> In addition, reduced GO, TiO<sub>2</sub>, nickel carbide (Ni<sub>3</sub>C), and black phosphorus have been employed as electron transport supports or co-catalysts to improve the photocatalytic activities.<sup>14-16</sup>

In aqueous solution systems, the electrons (e<sup>-</sup>) and holes (h<sup>+</sup>) generated in MAPbI<sub>3</sub> upon photoexcitation reduce protons (on materials such as Pt/TiO<sub>2</sub>, where the deposited Pt nanoparticles act as H<sub>2</sub>-producing co-catalysts) to produce H<sub>2</sub> and oxidize iodide ions (I<sup>-</sup>) to eventually produce triiodide ions (I<sub>3</sub><sup>-</sup>), as described by eqs 1–4.





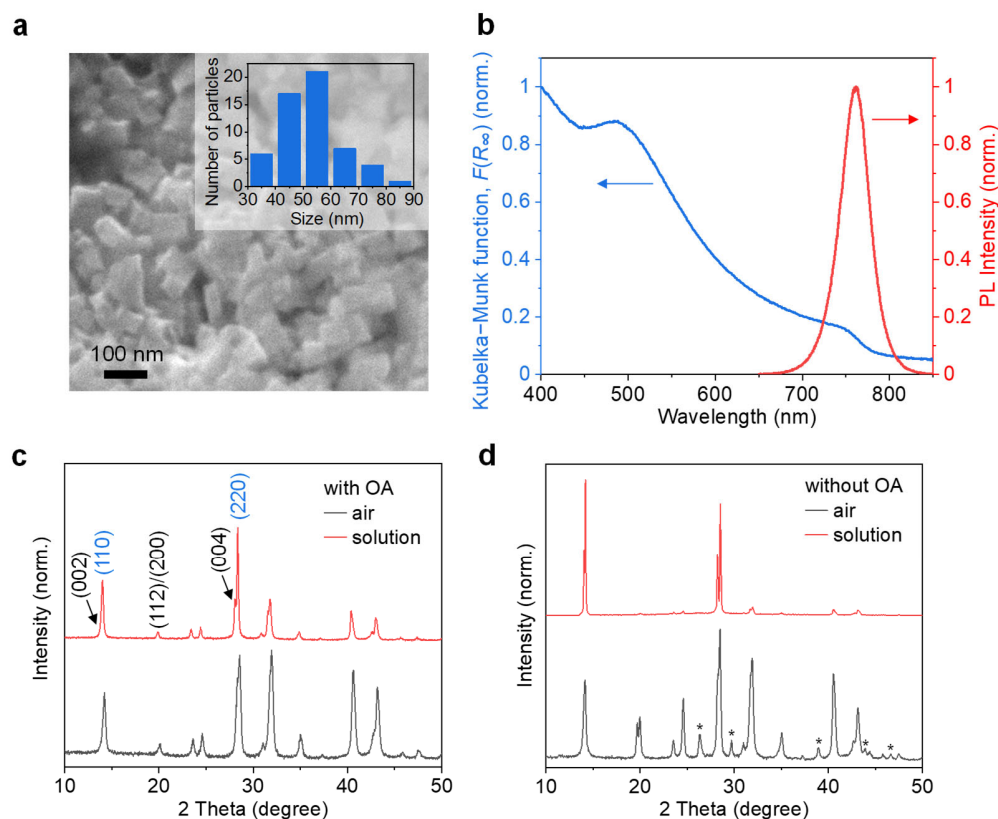
$\text{I}_3^-$  strongly absorbs visible light, which inhibits the light absorption by  $\text{MAPbI}_3$ . To achieve stable  $\text{H}_2$  production, phosphinic acid ( $\text{H}_3\text{PO}_2$ ) was added to the solution to reduce  $\text{I}_3^-$  to  $\text{I}^-$  (eq 5).<sup>12</sup>



The time scale of interfacial charge transfer and lifetime of photogenerated charge carriers are significantly influenced by the structural properties (size, shape, degree of crystallinity, point defects, impurities, etc.) of semiconductor materials and determine their photocatalytic performance.<sup>17</sup> However, quantitative information related to the structural stability and charge carrier dynamics of perovskites in aqueous media is not yet available. A proper understanding of these fundamental principles could lead to their high and stable photocatalytic activities.

In this study, we elucidated the structural and charge carrier dynamics of  $\text{MAPbI}_3$  nanoparticles in aqueous solution systems by analyzing their PL properties at the single-particle level. Single-molecule (single-particle) fluorescence spectroscopy allowed us to detect the emissions from a very few samples, ultimately individual molecules or particles, under a fluorescence microscope. It has been employed to solve the dynamic structural and chemical conversions over organo-lead halide perovskites, which are usually hidden in ensemble- and time-averaging bulk experiments.<sup>18,19</sup>

For single-particle observations, cuboid-shaped MAPbI<sub>3</sub> nanoparticles were synthesized by adding a poor solvent (dichloromethane) to a solution of precursors (MAI and PbI<sub>2</sub>) in a good solvent ( $\gamma$ -butyrolactone) at room temperature, followed by mixing with a solution of capping ligands (*n*-octylamine, OA) (Figure 1a).<sup>20</sup> The steady-state diffuse reflectance spectrum of the nanoparticles in toluene shows a continuous absorption band due to interband transition with an edge at ~765 nm (Figure 1b). The maximum PL emission appeared at ~765 nm, indicating that the observed PL originates from the radiative recombination of photogenerated electrons and holes in the MAPbI<sub>3</sub> (Figure 1b).<sup>21,22</sup> The powder X-ray diffraction (XRD) pattern indicates that the OA-capped MAPbI<sub>3</sub> nanoparticles possess a tetragonal structure (Figures 1c and S1).<sup>23</sup> From the Scherrer equation,<sup>24</sup> the crystallite size for the (220) plane was calculated to be ~31 nm. This value is reasonable compared with the statistical data of short-axis lengths (~53 nm) obtained from field emission scanning electron microscopy (FE-SEM) images (Figure 1a). The results of further structural and elemental analyses were reported in our previous paper.<sup>20</sup>



**Figure 1.** (a) FE-SEM image of OA-capped MAPbI<sub>3</sub> nanoparticles and (inset) particle size distribution. (b) Steady-state diffuse reflectance and PL spectra of OA-capped MAPbI<sub>3</sub> nanoparticles dispersed in toluene. (c,d) XRD patterns of MAPbI<sub>3</sub> nanoparticles synthesized with (c) and without (d) OA. The particles were dried in ambient air (lower) or immersed in MAPbI<sub>3</sub>-saturated HI-H<sub>3</sub>PO<sub>2</sub> solution (upper). Asterisks indicate characteristic peaks of hydrated species.

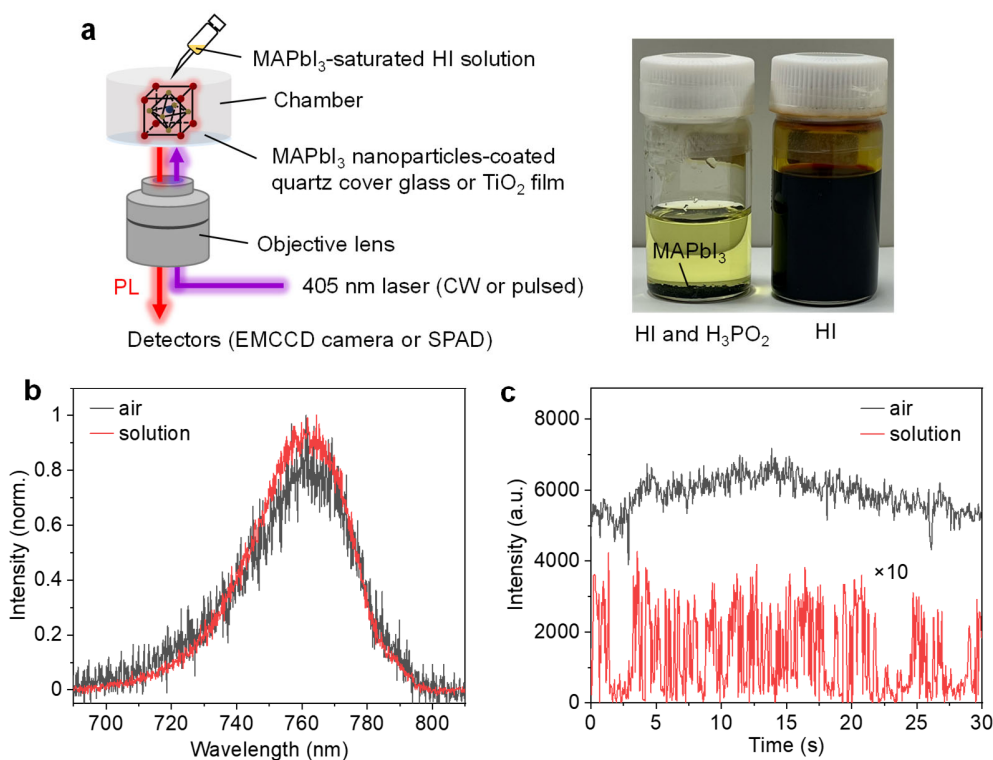
To examine the stability of the perovskite structure in aqueous solution, we exposed MAPbI<sub>3</sub> nanoparticles to MAPbI<sub>3</sub>-saturated HI solutions with and without H<sub>3</sub>PO<sub>2</sub> and then measured the XRD patterns. As shown in Figures 1c and S2, only the diffraction peaks indexed to the (110) and (220) planes were significantly enhanced after treatment. The crystallite size parallel to the (220) plane increased to ~66 nm on average, suggesting that the MAPbI<sub>3</sub> crystals tend to grow selectively along the [002] direction in the solution (Figure S1). This anisotropic growth may have been due to the adsorption of OA and/or excess MA<sup>+</sup> in solution on the (110) facets, leading to a reduction

in their surface energy.<sup>25-27</sup> On the other hand, after being exposed to ambient air for a while, the MAPbI<sub>3</sub> nanoparticles synthesized without OA exhibited several additional diffraction peaks, including those at  $\sim 26.3^\circ$ ,  $\sim 29.7^\circ$ , and  $\sim 39.0^\circ$ , which could be assigned to MAPbI<sub>3</sub> hydrates (Figure 1d).<sup>28,29</sup> Interestingly, only the peaks of pure MAPbI<sub>3</sub> were observed after exposing the MAPbI<sub>3</sub> hydrates to the MAPbI<sub>3</sub>-saturated HI solution. This result demonstrates that the MAPbI<sub>3</sub> structure did not decompose in the aqueous solution, but rather improved the crystallinity as a result of the dynamic equilibrium. The capping ligands work not only in obtaining homogeneous nanoparticles with a controlled size and shape, as is evident from FE-SEM image of MAPbI<sub>3</sub> nanoparticles synthesized without OA (Figure S3), but also in suppressing the hydration of MAPbI<sub>3</sub>, presumably through a hydrogen bonding interaction with surface iodide atoms.<sup>30</sup> Fourier transform infrared (FTIR) spectra confirmed the presence of the OA molecules on the surface of MAPbI<sub>3</sub> nanoparticles even after immersing in MAPbI<sub>3</sub>-saturated HI-H<sub>3</sub>PO<sub>2</sub> solution (Figure S4).

The PL properties of the OA-capped MAPbI<sub>3</sub> nanoparticles were first investigated using a wide-field and confocal microscope system after depositing the particles on a quartz cover glass and immersing them in MAPbI<sub>3</sub>-saturated HI-H<sub>3</sub>PO<sub>2</sub> solution (Figure 2a). Figure 2b shows the typical PL spectra of single OA-capped MAPbI<sub>3</sub> nanoparticles (or few-particle aggregates) exposed to ambient air and the aqueous solution under 405 nm excitation. The peak wavelengths ( $760 \pm 5$  nm) observed under the two conditions were almost the same, suggesting that no significant structural changes such as hydration or particle size decrease<sup>31</sup> occurred during the observation period. Notably, approximately half the bare MAPbI<sub>3</sub> nanoparticles in ambient air exhibited the PL peak wavelengths same as those ( $760 \pm 5$  nm) of OA-capped ones, while the others possess the peaks blue-shifted by approximately 10–20 nm (0.02–0.04 eV), possibly due to hydration (Figure S5).<sup>32,33</sup> After immersion in the solution, however, no such blue shift was observed. These results



imply the dissolution and recrystallization of hydrated MAPbI<sub>3</sub> nanoparticles, leading to the formation of MAPbI<sub>3</sub> with a pristine crystal structure under dynamic equilibrium, in accordance with the XRD results (Figure 1d). Another possible explanation is that blue-shifted PL was almost completely quenched due to further hydration or decomposition. Thus, OA-capped MAPbI<sub>3</sub> nanoparticles were employed herein for the long-term PL observation of single particles.

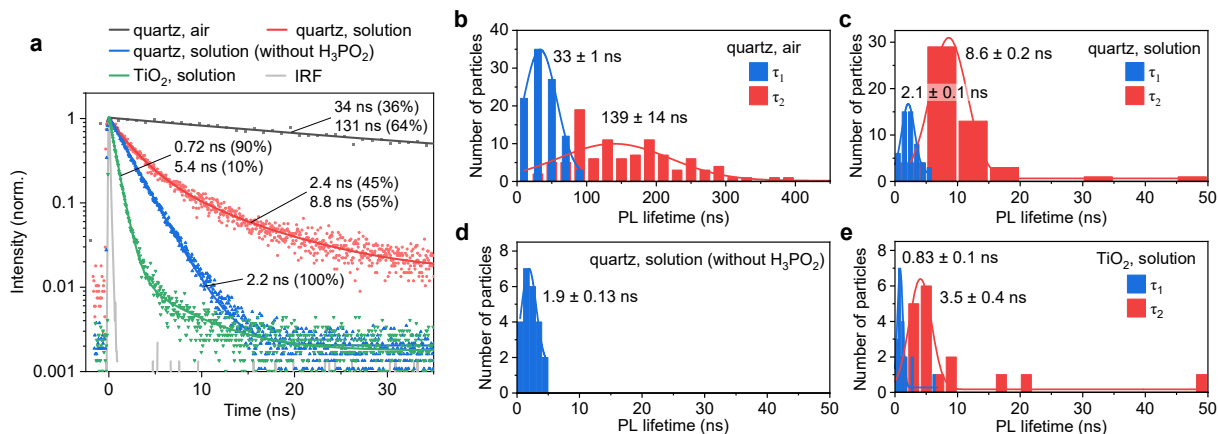


**Figure 2.** (a) Schematic of the PL microscopy measurement system. Photograph shows the MAPbI<sub>3</sub>-saturated HI aqueous solutions with (left) and without (right) H<sub>3</sub>PO<sub>2</sub> that reduces I<sub>3</sub><sup>-</sup> to I<sup>-</sup>. (b,c) Typical PL spectra (b) and intensity trajectories (c) of single OA-capped MAPbI<sub>3</sub> nanoparticles dispersed on quartz cover glasses. The samples were exposed to ambient air (black) or to MAPbI<sub>3</sub>-saturated HI-H<sub>3</sub>PO<sub>2</sub> solution (red).

After being exposed to the MAPbI<sub>3</sub>-saturated HI-H<sub>3</sub>PO<sub>2</sub> solution, most of the particles exhibited the so-called PL blinking, where the particle emitted in an intermittent on-and-off fashion (Figures 2c, S6, and Supporting Movies S1 and S2).<sup>34</sup> This blinking phenomenon is commonly observed

for single semiconductor nanostructures including metal halide perovskites<sup>18,35–38</sup> and mostly originates from the suppression of radiative recombination due to the stochastic trapping of charge carriers by surface defects or the Auger recombination, which quenches the emission via the transfer of exciton energy to the third carrier in the particle.<sup>39,40</sup> The intriguing blinking features of the MAPbI<sub>3</sub> nanoparticles will be discussed later.

To investigate the charge carrier dynamics of MAPbI<sub>3</sub> nanoparticles, time-resolved PL measurements were carried out on individual particles. The PL decay profiles were fitted by a single or double exponential decay function to determine the PL lifetimes and their contributions (Figure 3a). Under the air condition, the faster component ( $\tau_1$ ) (~33 ns) is assigned to the trapping of charge carriers by surface defects, and the slower component ( $\tau_2$ ) (~139 ns) is assigned to the radiative recombination of charge carriers through intraparticle diffusion (Figures 3b and S7, and Table 1).<sup>2,41</sup> The wide distributions of  $\tau$  values would be mainly attributed to the variation in the particle size because a linear relationship between  $\tau_1$  and  $\tau_2$  obtained for the same particle in air (Figure S8) implied that both processes had a size-dependent characteristic.<sup>20</sup> After being exposed to the MAPbI<sub>3</sub>-saturated HI-H<sub>3</sub>PO<sub>2</sub> solution, the PL lifetimes decreased by one order of magnitude or more ( $\tau_1$  of ~2.1 ns and  $\tau_2$  of ~8.6 ns) (Figure 3c). This result is consistent with the weakened PL and implies that there were additional nonradiative pathways caused by the chemical species in the aqueous solution. For the bare nanoparticles in ambient air,  $\tau_1$  was shortened to ~4.2 ns (Figure S9), possibly as a result of partial hydration, while  $\tau_2$  (~32 ns) was very close to the  $\tau_1$  value (~33 ns) of the OA-capped sample, which was attributed to defect-mediated recombination, as mentioned above. In MAPbI<sub>3</sub>-saturated HI-H<sub>3</sub>PO<sub>2</sub> solution, much shorter components (~0.4 ns) appeared for the bare nanoparticles, suggesting that the hole transfer was accelerated through direct interactions with I<sup>-</sup> and H<sub>3</sub>PO<sub>2</sub>.



**Figure 3.** (a) PL decay profiles (dots) observed for typical single MAPbI<sub>3</sub> nanoparticles during 405-nm pulsed laser irradiation. The solid lines show the nonlinear least-squares fittings to exponential functions. The whole profile obtained for the particle on quartz in air is given in Figure S7. (b–e) Histograms of lifetimes for individual nanoparticles dispersed on quartz cover glasses (b–d) and a TiO<sub>2</sub>-coated quartz cover glass (e). The samples were exposed to ambient air (b) or to MAPbI<sub>3</sub>-saturated HI solution with (c,e) or without (d) H<sub>3</sub>PO<sub>2</sub>. The solid lines show the nonlinear least-squares fittings to Gaussian functions.

**Table 1.** Average PL Lifetimes of OA-Capped MAPbI<sub>3</sub> Nanoparticles and Their Related Reaction Processes

conditions	$\langle\tau_1\rangle$ (reaction process)	$\langle\tau_2\rangle$ (reaction process)
quartz, air	33 ns (e <sup>-</sup> /h <sup>+</sup> trapping by surface defects)	139 ns (radiative recombination)
quartz, solution	2.1 ns (h <sup>+</sup> transfer to I <sup>-</sup> )	8.6 ns (h <sup>+</sup> transfer to H <sub>3</sub> PO <sub>2</sub> )
TiO <sub>2</sub> , solution	0.83 ns (e <sup>-</sup> transfer to TiO <sub>2</sub> )	3.5 ns (h <sup>+</sup> transfer to I <sup>-</sup> /H <sub>3</sub> PO <sub>2</sub> )

To identify the chemical species that induced nonradiative charge-transfer processes, we prepared a MAPbI<sub>3</sub>-saturated HI solution without H<sub>3</sub>PO<sub>2</sub> and carried out the same measurements. As demonstrated in Figure 3d, only one component (~1.9 ns) was observed, suggesting that the fast (~2.1 ns) and slow (~8.6 ns) components in the presence of H<sub>3</sub>PO<sub>2</sub> are assignable to hole transfers to I<sup>-</sup> and H<sub>3</sub>PO<sub>2</sub>, respectively. The observed overall rate is given as  $k_{\text{obs}} = \langle n \rangle \times k'_{\text{HT}}$ , where  $\langle n \rangle$  is the average number of bound acceptors per particle, and  $k'_{\text{HT}}$  is the intrinsic hole transfer rate. The  $\langle n \rangle$  value is basically dependent on the concentration and adsorption constant (i.e., the interaction between the MAPbI<sub>3</sub> and hole acceptors) of the hole acceptors in solution at low surface coverage values. The molar ratio of I<sup>-</sup>:H<sub>3</sub>PO<sub>2</sub> in the solution was 3.3:1.0. Considering the pK<sub>a</sub> values of HI (-10) and H<sub>3</sub>PO<sub>2</sub> (1.23)<sup>42</sup> and the pH of the solution (approx. -0.7), the H<sub>3</sub>PO<sub>2</sub> was almost fully protonated. Therefore, I<sup>-</sup> ions were preferentially adsorbed on the cationic sites on the surface of the MAPbI<sub>3</sub> because of electrostatic interactions, which effectively captured holes. A further discussion of  $k'_{\text{HT}}$  within the framework of the Marcus theory is provided in Supporting Note S1.

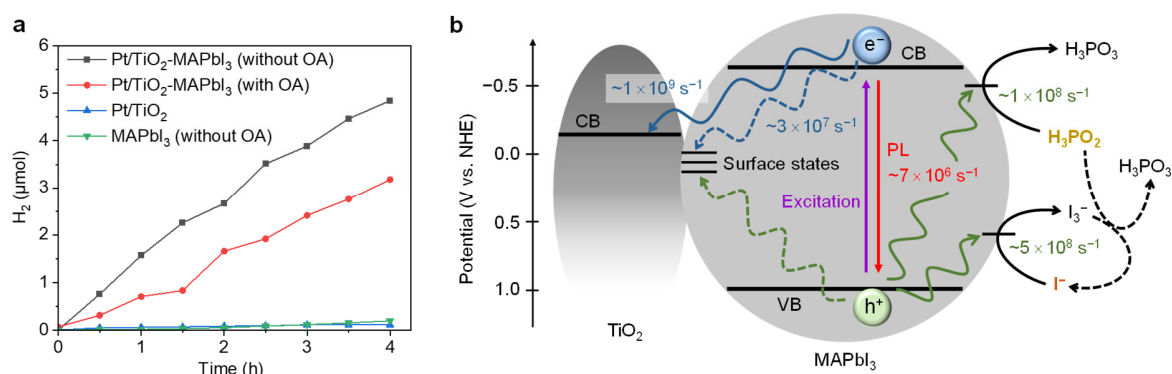
To gain further insight into the design of the photocatalytic H<sub>2</sub> production system, we examined the PL properties of OA-capped MAPbI<sub>3</sub> nanoparticles on the TiO<sub>2</sub> film formed on the quartz cover glass (Figure S10). Figure S11 shows the PL trajectories captured under the air and solution conditions. In air, the MAPbI<sub>3</sub> nanoparticles exhibited blinking, but were quickly photobleached. Because the conduction band (CB) of TiO<sub>2</sub> is more positive than that of MAPbI<sub>3</sub>, photogenerated electrons in MAPbI<sub>3</sub> were quickly transferred to the TiO<sub>2</sub>, which may result in the accumulation of excess holes that escaped the recombination processes, thus leading to the decomposition of the perovskite structure.<sup>43,44</sup> This result also indicates that holes in the valence band (VB) of MAPbI<sub>3</sub> are transferred to I<sup>-</sup> or H<sub>3</sub>PO<sub>2</sub>, which plays a crucial role in maintaining the overall charge balance

in the perovskite to avoid undesired photodegradation and catalyze the HI splitting reaction for H<sub>2</sub> generation.<sup>13</sup>

Time-resolved measurements revealed that the PL lifetimes of the OA-capped MAPbI<sub>3</sub> nanoparticles on TiO<sub>2</sub> were much shorter than those on quartz when exposed to MAPbI<sub>3</sub>-saturated HI-H<sub>3</sub>PO<sub>2</sub> solutions (Figures 3a, c, and e). This is because electron transfer to the TiO<sub>2</sub> is the dominant competitive process. From these results, we can assign the faster component (<1 ns) to electron transfer to TiO<sub>2</sub> and the slower component (3.5 ns) to hole transfer to I<sup>-</sup> and/or H<sub>3</sub>PO<sub>2</sub> (Figure 3e). The electron/hole transfer rates can be calculated from  $k_{ET/HT} = \langle \tau_{\text{acceptor}} \rangle^{-1}$ , where  $\langle \tau_{\text{acceptor}} \rangle$  is the specific lifetime component for the electron/hole transfer to the acceptors (Table 1). The hole transfer rate for I<sup>-</sup> ( $4.8 \times 10^8 \text{ s}^{-1}$ ) is approximately one-tenth of the hole transfer rate to the hole transport layers (Spiro-OMeTAD and PEDOT:PSS) of perovskite solar cells.<sup>45,46</sup> This is possibly because the charge transfer between a solid and solid interface is more efficient than that between liquid and solid interfaces owing to stronger electronic couplings with hole acceptors. Improving the hole transfer rate would be crucial to improve the stability of perovskites and thus enhance their photocatalytic activities.

The rate of electron transfer from OA-capped MAPbI<sub>3</sub> nanoparticles to TiO<sub>2</sub> ( $\langle \tau_1 \rangle^{-1} = 1.2 \times 10^9 \text{ s}^{-1}$ ) in MAPbI<sub>3</sub>-saturated HI-H<sub>3</sub>PO<sub>2</sub> solution is comparable to that reported for a MAPbI<sub>3</sub>/TiO<sub>2</sub> film in air.<sup>47</sup> The capping ligands or stabilizing agents can greatly improve the durability and physicochemical properties of nanoparticles in solution but often limit the accessibility of substrates on the surface and hinder the direct charge transfer between the reactants and active sites.<sup>48-50</sup> This is as true for the present systems, according to the fact that the rate of electron transfer from bare MAPbI<sub>3</sub> nanoparticles to TiO<sub>2</sub> in the solution was determined to be  $2.2 \times 10^9$

$s^{-1}$  (Figure S12). To examine the relationship between interfacial charge transfer dynamics and photocatalytic performance,  $H_2$  generation via visible-light-driven HI splitting was tested. As demonstrated in Figure 4a, MAPbI<sub>3</sub> nanoparticles both with and without OA showed stable  $H_2$  production by mixing with Pt-modified TiO<sub>2</sub> nanoparticles (Pt/TiO<sub>2</sub>) (Figure S13). Relative  $H_2$  production efficiencies of bare and OA-capped samples follow the relative rates of electron transfer (i.e.,  $2.2 \times 10^9 s^{-1}/1.2 \times 10^9 s^{-1} = 1.8$ ), revealing the critical role of surface-capping ligands in determining the photochemical activity of the MAPbI<sub>3</sub>/TiO<sub>2</sub> hybrid systems.



**Figure 4.** (a) Photocatalytic  $H_2$  evolution from the samples in MAPbI<sub>3</sub>-saturated HI- $H_3PO_2$  solutions under visible light irradiation (center wavelength = 565 nm; intensity =  $50 \text{ mW cm}^{-2}$ ). 10-mg MAPbI<sub>3</sub> nanoparticles synthesized with (red) or without (black) OA were mixed with 5-mg Pt/TiO<sub>2</sub> in 5 mL Ar-saturated solution. Pt/TiO<sub>2</sub> (5 mg; blue) and bare MAPbI<sub>3</sub> nanoparticles (10 mg; green) were used for control experiments. (b) Energy diagram of charge transfer processes on the OA-capped MAPbI<sub>3</sub>/TiO<sub>2</sub> system in aqueous solution.

Figure 4b illustrates the energy diagram and charge transfer processes of the MAPbI<sub>3</sub>/TiO<sub>2</sub> system in aqueous solution with representative reaction rates. First, electrons and holes are generated in the CB and VB of MAPbI<sub>3</sub>, respectively, upon excitation with photons with energies larger than the bandgap. The electrons are then transferred to the CB of TiO<sub>2</sub> located at  $-0.12 \text{ V}$  vs. NHE<sup>51</sup> ( $\sim 10^9 s^{-1}$ ), followed by hole transfer to I<sup>-</sup> in aqueous solution ( $10^8$ – $10^9 s^{-1}$ ). At the same

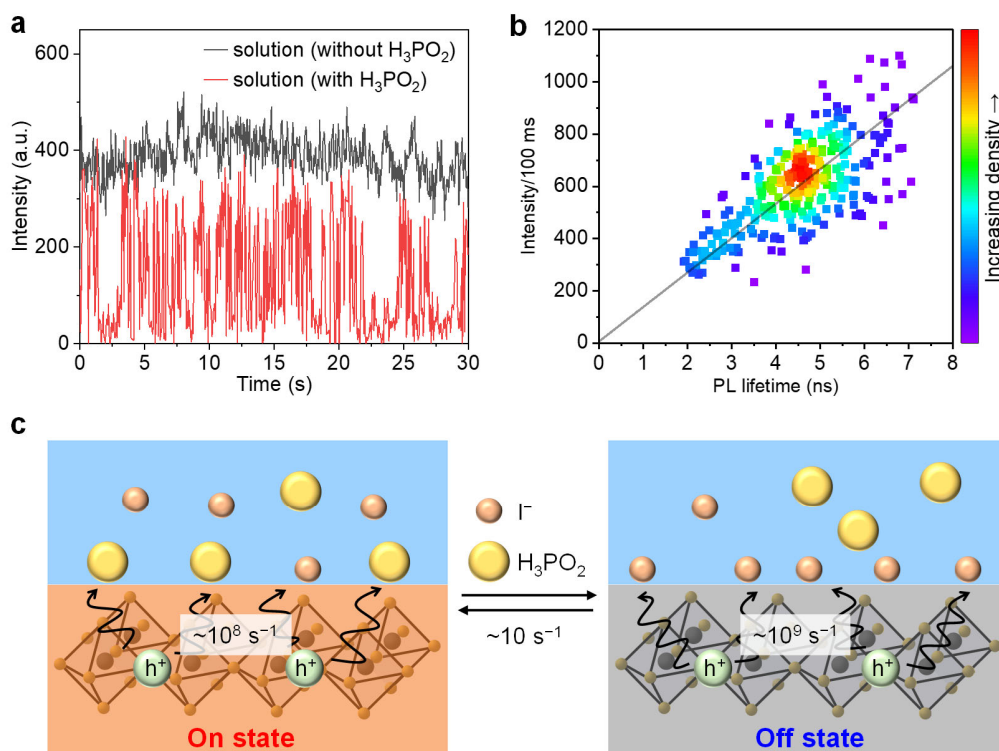
time, trap-mediated recombination, or hole transfer to  $\text{H}_3\text{PO}_2$  occurs at a slower rate ( $\sim 10^8 \text{ s}^{-1}$ ). Without  $\text{TiO}_2$ , the electrons in the  $\text{MAPbI}_3$  would be transferred to molecular oxygen adsorbed on the surface.<sup>43,52</sup> The surviving electrons and holes recombine radiatively ( $\sim 10^7 \text{ s}^{-1}$ ). Because these reaction rates temporally fluctuate, as suggested by the PL blinking, minor (slower) pathways are clearly recognized in single-particle experiments.

Finally, we discuss the origin of the PL blinking observed for the OA-capped  $\text{MAPbI}_3$  nanoparticles in solution. There are two possibilities. The PL blinking might have been the result of stochastic charge trapping by surface defects and adsorbates or Auger recombination-assisted quenching. However, the contribution of Auger recombination was relatively small based on the fact that no remarkable blinking was observed when the excitation power was increased (Figure S14).<sup>53</sup> This conclusion is also supported by a statistical analysis of the blinking characteristics (see Supporting Note S2).

The OA-capped  $\text{MAPbI}_3$  nanoparticles in air showed no or less blinking, while those in aqueous solution exhibited many blinking cycles with on and off times ( $\tau_{\text{on}}$  and  $\tau_{\text{off}}$ ) on a time scale of milliseconds to seconds. These results suggest that the chemical species in the solution captured the charge carriers at a higher turnover rate. In addition, the nanoparticles in solution without  $\text{H}_3\text{PO}_2$  did not show distinct blinking at a frame rate of 20 fps, implying that  $\text{H}_3\text{PO}_2$  is essential for blinking on a time scale of tens of milliseconds to seconds (Figure 5a). The fluorescence lifetime-intensity distribution (FLID) map shows an almost linear relationship between them, indicating a temporal fluctuation of the hole transfer dynamics (Figure 5b).

Based on the results of the blinking analysis, we propose a possible blinking mechanism, as shown in Figure 5c.  $\text{MAPbI}_3$  nanoparticles emit PL (the on-state) when  $\text{H}_3\text{PO}_2$  molecules, which

are ineffective hole acceptors, temporarily occupy the adsorption sites on the MAPbI<sub>3</sub> surface. Meanwhile, the MAPbI<sub>3</sub> nanoparticles do not emit PL (the off-state) when I<sup>-</sup> ions, which are effective hole acceptors, are predominantly adsorbed on the surface. Our interpretation fits well with the model that the opening and closing of nonradiative channels in luminescent quantum dots are the result of fluctuations in the adsorbate binding to the particle surface.<sup>54</sup> This dynamic fluctuation is vital not only for the PL blinking phenomenon but also for the single turnover of the catalyzed conversion of substrates on solid (photo)catalysts.<sup>55,56</sup> A more complete understanding of such substrate-induced fluctuation events will provide fundamental insights into (photo)catalytic behaviors in dynamic equilibrium systems, which are often inaccessible in ensemble-averaged measurements.



**Figure 5.** (a) PL intensity trajectories of single OA-capped MAPbI<sub>3</sub> nanoparticles in MAPbI<sub>3</sub>-saturated HI solutions with (red) and without (black) H<sub>3</sub>PO<sub>2</sub>. (b) A typical FLID map obtained for



a single OA-capped MAPbI<sub>3</sub> nanoparticle in MAPbI<sub>3</sub>-saturated HI-H<sub>3</sub>PO<sub>2</sub> solution. The solid line indicates the fit of the linear relationship between them. (c) Mechanistic illustration of the blinking behavior under dynamic solid–solution equilibrium condition.

In conclusion, we successfully characterized the structural and PL properties of MAPbI<sub>3</sub> nanoparticles in aqueous solution. To the best of our knowledge, this is the first example of PL from metal halide perovskites in aqueous solutions. The perovskite structure is stable, especially with the aid of capping ligands, in the perovskite-saturated HI solution and recovered from decomposed ones, as revealed by the XRD data. Based on the remarkable PL blinking, as well as the decreased PL intensity and lifetimes for OA-capped MAPbI<sub>3</sub> nanoparticles in MAPbI<sub>3</sub>-saturated HI-H<sub>3</sub>PO<sub>2</sub> solution, we propose that the temporal fluctuation of the charge carrier trapping originates from the stochastic and competitive adsorption of hole acceptors, I<sup>-</sup> and H<sub>3</sub>PO<sub>2</sub>, where the former is a more effective species under the present conditions. We further note that interfacial electron transfer from excited MAPbI<sub>3</sub> to the TiO<sub>2</sub> support competes with these hole transfer processes and eventually determines the efficiency of photocatalytic H<sub>2</sub> evolution. Our single-particle approach for investigating the intricate nature of photocatalysis in dynamic solid–solution equilibrium will provide a design principle for a new type of (photo)catalytic system apart from conventional homogeneous and heterogeneous systems.

## ASSOCIATED CONTENT

**Supporting Information.** Experimental methods, additional results and discussion. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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## Notes

The authors declare no competing financial interest.

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