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# Infrared Absorption of $\text{Zn}_{0.5}\text{Cd}_{0.5}\text{S}$ Photocatalyst Bandgap-Excited under Aqueous Environment

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## ABSTRACT

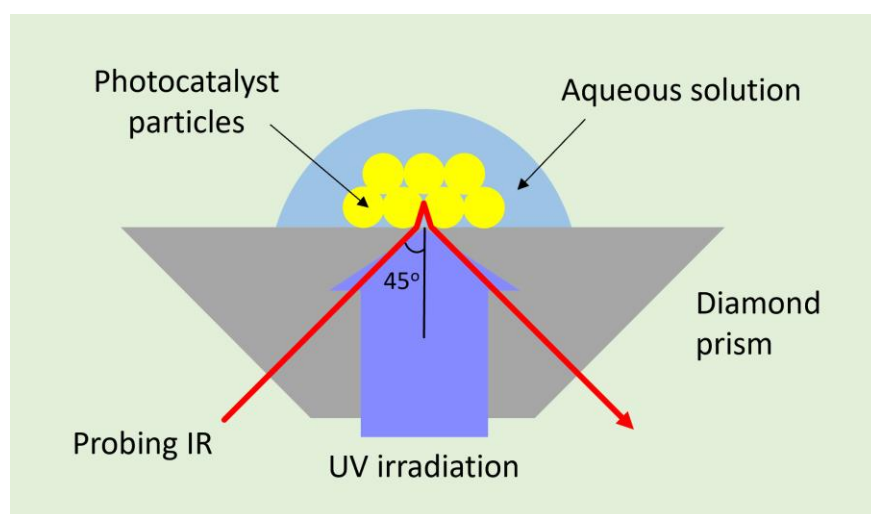
The adoption of infrared and near-infrared spectrometry technique for mechanistic and kinetic studies on the photocatalytic water splitting reaction has been steadily increasing over the years. Herein, a transition metal sulfide photocatalyst,  $\text{Zn}_{0.5}\text{Cd}_{0.5}\text{S}$  (ZCS), which has been proven to be an efficient hydrogen ( $\text{H}_2$ ) evolution photocatalyst, was investigated through attenuated total reflection FTIR (ATR-FTIR). Electronic absorption in the wavenumber region of 500 – 3000  $\text{cm}^{-1}$  was detected while the photocatalyst was irradiated by UV light source in a pure water environment. The source of absorption was further characterized by employing scavenging solutions, and it was found that the contribution from photogenerated electrons and holes coincidentally appeared in the same wavenumber window. This is likely to be unprecedented as prior studies generally were not able to observe electronic absorption by holes, or the holes absorption was present close to the visible wavenumber window instead. Nevertheless, the shape of absorption peak associated to holes was noticeably distinctive from that of electrons when their normalized spectra were compared, allowing easy identification of the nature of the detected absorption. These observations suggested that electrons and holes trap states were present in ZCS, and it is proposed that the trap states could be innately present due to the formation of physical defects during chemical synthesis, or they were instantaneously created upon photoexcitation where polarons were manifested by the disruption of ionic equilibrium in valence band (VB) and conduction band (CB) through injection of foreign photogenerated electrons or holes.

## INTRODUCTION

Since the pioneering report by Honda and Fujishima,<sup>1</sup> the conversion of water into hydrogen ( $H_2$ ) and oxygen ( $O_2$ ) using semiconductor photocatalysts has continued to garner substantial interest over the past few decades.  $H_2$  holds arguable more value as compared to  $O_2$  due to it boasting a high energy content of 142 kJ/g, much higher than that of conventional fossil fuels.<sup>2</sup> The combustion of  $H_2$  for energy generation also does not leave behind any carbon footprint, deeming it to be a promising source of green fuel in near future. Through photocatalytic water splitting, the process of generating  $H_2$  is also much more environmentally friendly as compared to other developed methods such as steam reforming and electrolysis, since water and sunlight are the only two resources needed alongside the photocatalyst to kickstart the reaction. Nevertheless, this photocatalytic approach still has a glaring downside of insufficient efficiency for real-life applications at the present state of technology, despite several impressive breakthroughs reported in recent years.<sup>3,4</sup> In view of this, the adoption of effective analytical techniques in hope of grasping more of the underlying kinetics and mechanisms behind the photocatalytic reaction becomes imperative, so that a stronger foundation can be established for the ensuing design and modification of photocatalysts.

In-situ FTIR analysis is a technique frequently used for the detection of chemical reaction intermediates as well as unraveling of reaction pathway and mechanisms. For the analysis of water splitting reaction, things can become more intricate as the water medium tends to absorb the probing infrared (IR) light which greatly lowers the signal strength of the resultant spectrum. Attenuated total reflection (ATR) in a diamond prism comes in handy in this scenario because of its simplicity in accommodating the measurement of photocatalyst particles suspended in water (Figure 1). A number of reports have demonstrated the use of this technique for the detection of vibrational absorption arising from the intermediates involved in photocatalytic water splitting.<sup>5-8</sup> However, there is another important phenomenon that can be observed

through ATR-FTIR, but it has been relatively less investigated, and that is the electronic absorption spurred by photoexcited charge carriers in photocatalysts. Through comprehending the absorption signal induced by the charge carriers, scientific information such as the presence and nature of shallow states trapping the charge carriers as well as the decay kinetics of the carriers can be unveiled.<sup>9</sup> Hence it is believed that this piece of information should not be neglected.



**Figure 1.** Diamond prism used for ATR spectroscopy.

To-date, most photocatalysts that had their electronic absorption signals studied through IR spectroscopy were oxide-based semiconductors, for instance,  $\text{TiO}_2$ ,<sup>10–16</sup>  $\text{SrTiO}_3$ ,<sup>17,18</sup> and  $\text{NaTaO}_3$ .<sup>19–21</sup> This is not surprising as oxide-based photocatalysts are known to demonstrate good activity, high robustness while also being comparably easy to synthesize. In fact, one of the most efficient water splitting photocatalytic compounds reported hitherto is based on  $\text{SrTiO}_3$  with the support of co-catalysts  $\text{Rh/Cr}_2\text{O}_3$  and  $\text{CoOOH}$ , which in turn resulting in a groundbreaking quantum efficiency of near unity.<sup>4</sup> Despite that, majority of the oxide-based semiconductors suffer from possessing a large bandgap, rendering them active only under illumination of short wavelength UV light.

From this perspective, sulfide-based semiconductors emerge as promising photocatalysts since they generally carry bandgaps below 2.6 eV, with the exception of ZnS.<sup>22</sup> Although the shortcoming of sulfides being prone to photocorrosion is undeniable, researchers have demonstrated many times the potential of sulfides as candidates for photocatalytic water splitting.<sup>23–32</sup> Among them,  $\text{Zn}_x\text{Cd}_{1-x}\text{S}$  turns out to be a popular choice owing to its tunable bandgap based on the molar ratio between Zn and Cd, which is convenient depending on the premise of photocatalytic reactions. Moreover, it has also been reported that through appropriate conditioning of the hydrothermal synthesis process, homojunctions of the zinc blende and wurtzite phases provoked by lattice disruption can be generated. It has been claimed that the presence of these homojunctions essentially created a step-like structure within its hybridized states which is conducive to prolonging the lifetime of charge carriers, thus giving rise to an enhancement in activity.<sup>33–36</sup> Nonetheless, despite being on the hot radar of photocatalysis, the kinetic and mechanistic aspects of the carriers fostered by illumination of  $\text{Zn}_x\text{Cd}_{1-x}\text{S}$  have rarely been investigated. Herein, the change in IR absorption in  $\text{Zn}_{0.5}\text{Cd}_{0.5}\text{S}$  (ZCS) during UV light illumination was observed in several aqueous solutions, and the nature of the states and their formation mechanism were deliberately speculated and reported.

## METHODS

**Photocatalyst preparation.** ZCS was synthesized through a precipitation-hydrothermal approach slightly modified from a previous report.<sup>25</sup> 4 mmol of zinc acetate dihydrate (99%, FUJIFILM Wako Pure Chemicals) and 4 mmol of cadmium acetate dihydrate (98%, FUJIFILM Wako Pure Chemicals) were dissolved completely in 70 mL of ultrapure water, before slowly adding in 10 mmol of thioacetamide (FUJIFILM Wako Pure Chemicals). After 30 mins of vigorous stirring, 10 mL of hydrazine monohydrate (97%, FUJIFILM Wako Pure Chemicals) was added in a dropwise manner and the mixture is further allowed to be stirred for 30 mins. Next, the resultant yellowish solution was transferred into an airtight Teflon-lined stainless-

steel autoclave, which was then placed in an oven for the hydrothermal treatment. The heating temperature was set as 190 °C and the reaction was conducted for 24 hours. After allowing the reactant to be cooled down to room temperature, the precipitate was retrieved by centrifugation and was washed two times each using ultrapure water and ethanol. The washed sample was then dried overnight at 60 °C in a drying oven.

To validate the formation of ZCS, the crystallographic structure of the resultant sample was analyzed through XRD using Rigaku Smartlab diffractometer with Cu K $\alpha$  radiation. The diffraction pattern displayed multiphase characteristics, where peaks indexed to the cubic zinc blende phase and hexagonal wurtzite phase were observed and marked accordingly in Figure S1. The morphology of ZCS was also inspected through FE-SEM (JEOL JSM-7500F), and it was found that the ZCS particles have an average diameter of 20-80 nm, as seen in Figure S2. Both characterizations tallied with the results acquired in the previous works.<sup>25,34</sup>

**ATR spectroscopy.** The ATR assembly used in this study comprises of a diamond prism manufactured by Jasco. As presented earlier in Figure 1, the illustrated configuration is different from the vast majority of other reports in similar field of study, such that the excitation-triggering light source is illuminated on the same surface probed by IR beam. The incident angle of the probing IR was fixed at 45° from the normal of the reflection plane, whereas the UV light was guided perpendicularly to the plane from the bottom. Both light pass through the 1.8 mm diameter reflection plane and fall directly on the sample placed on top of the plane. The depths of IR (wavenumber: 1000 cm<sup>-1</sup>) light penetration into water was estimated to be 1.5  $\mu$ m based on the refractive indices of diamond and water which are 2.38 and 1.32 respectively.<sup>37</sup> The assembly was set in a Fourier transform infrared spectrometer (FT/IR-6600, Jasco), in which a He-Ne laser light was used to irradiate the reflection plane for interferometer calibration. In this study, IR absorption in the 9000 – 500 cm<sup>-1</sup> wavenumber region was measured (results were only presented from 6000 – 500 cm<sup>-1</sup> instead as there were

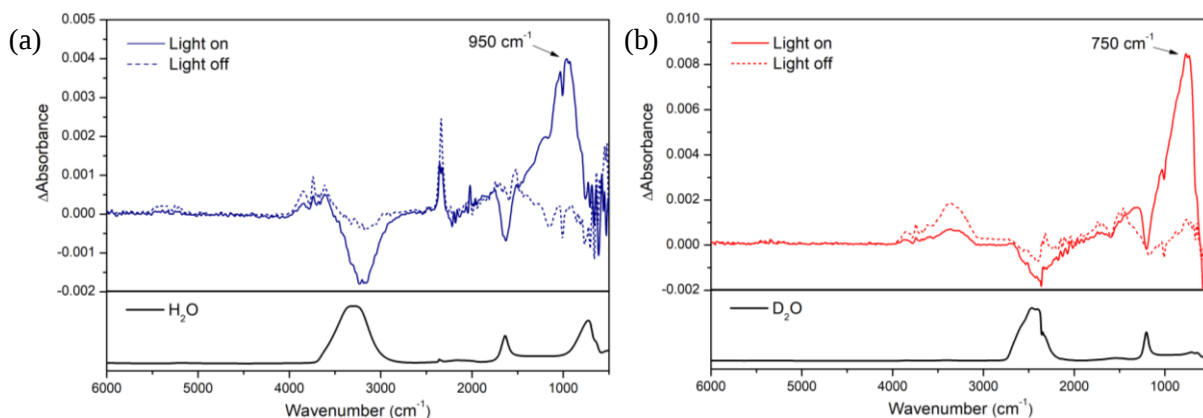
no responses detected in 9000 – 6000  $\text{cm}^{-1}$  region), and for this purpose, a ceramic light source and a mercury-cadmium-telluride (MCT) detector were deployed. With a wavenumber resolution of 16  $\text{cm}^{-1}$ , 64 rounds of measurements were collected and averaged to obtain the final spectrum. The total acquisition time was about 30 seconds per spectrum.

For each set of IR absorption measurement, the sample was first placed onto the reflection plane such that it is sufficient to cover the light irradiation window completely. 300  $\mu\text{L}$  of liquid was dropped onto the sample, followed by lowering the press to tighten the sample for a better result acquisition. A spectrum was then collected in the dark to serve as the background. Next, a light-emitting diode (LED) with 365 nm center wavelength (M365L3, Thorlabs) which was used as the light source for bandgap excitation was switched on to irradiate the sample for 1 minute, allowing the signals to reach a steady state. The UV light power density on the reflection plane was calibrated using a photodiode sensor (PD-300, Ophir) to be 3600  $\text{Wm}^{-2}$ . The change in absorption spectrum ( $\Delta\text{absorption}$ ) was then collected for evaluation. In between acquisition sets where different solution was used, the reflection plane was carefully wiped with ethanol and conditioned with the next solution to minimize the undesirable effect from any remaining impurities. Note that all measurements were conducted in open air, and the concentrations of sacrificial reagents, namely L(+)-Ascorbic acid (99.6%, FUJIFILM Wako Pure Chemicals) and  $\text{FeCl}_3$  (99% iron (III) chloride hexahydrate, Junsei Chemical), were standardized to be 0.05 M.

## RESULTS AND DISCUSSION

**IR absorption of ZCS in  $\text{H}_2\text{O}$  and  $\text{D}_2\text{O}$ .** ZCS immersed in pure water was first placed on the ATR crystal and analyzed for its changes in absorbance ( $\Delta\text{absorbance}$ ) spectrum with UV light illumination. From the resultant spectra presented in Figure 2 (a), it was observed that an absorption peak emerged at the 2000-750  $\text{cm}^{-1}$  region when the excitation UV light source was

turned on. This peak quickly subsided once the light was turned off. The shape of this peak appeared to be asymmetrical with the absorption maximum at about  $950\text{ cm}^{-1}$ . These observations point towards a deduction where the change in absorption band was brought upon by the charge carriers that were photogenerated when ZCS was subjected with UV light. As ZCS synthesized for this study possesses a bandgap of  $2.36\text{ eV}$ ,<sup>25</sup> theoretically it is not possible for it to absorb the probing IR light. However, if a shallow trap state is present in ZCS during light irradiation, the charge carriers which are temporarily trapped before recombination occurs can promptly utilize the probing IR light energy and re-excite back to a higher level of unoccupied energy band. This phenomenon is thus reflected in the IR  $\Delta$ absorbance spectrum, in which the energy level of such trap state in the photocatalyst can be estimated through the wavenumber position of the absorption band. Further details will be discussed comprehensively later in this manuscript. Hereby, the term ‘electronic absorption’ is introduced to represent this phenomenon, and it will be used recurrently throughout this manuscript.



**Figure 2.**  $\Delta$ Absorbance spectra of ZCS immersed in (a)  $\text{H}_2\text{O}$  and (b)  $\text{D}_2\text{O}$ . The solid and dotted lines represent the spectra acquired with and without light irradiation, respectively. The absorption bands of pure  $\text{H}_2\text{O}$  and  $\text{D}_2\text{O}$  were also included at the bottom of (a) and (b) correspondingly as a reference.

Besides the apparent peak in  $2000\text{--}750\text{ cm}^{-1}$  region, negative peaks superimposed at  $\sim 3180$  and  $\sim 1650\text{ cm}^{-1}$  were also noticed. These can be associated to the inevitable bleaching effect

imposed by UV illumination on the vibrational absorption of water, namely O-H stretching mode and H-O-H bending mode correspondingly. Similar observations were also reported in prior works.<sup>19,20</sup>

To further testify this, the same measurement was conducted by using D<sub>2</sub>O in place of H<sub>2</sub>O, and the result is showed in Figure 2 (b). As expected, negative peaks were observed at lower wavenumber of  $\sim 2480$  and  $1200\text{ cm}^{-1}$  where the O-D stretching and D-O-D bending mode are typically found. An additional positive peak also emerged at the same region as the O-H stretching mode ( $\sim 3250\text{ cm}^{-1}$ ), and it can be ascribed to either the remaining water component adsorbed on ZCS, or atmospheric water vapor infringing through the probing volume. On the other hand, electronic absorption was also clearly observed in D<sub>2</sub>O, however interestingly the band has a peak at slightly lower wavenumber in D<sub>2</sub>O as compared to when ZCS was immersed in H<sub>2</sub>O ( $750\text{ cm}^{-1}$  vs  $950\text{ cm}^{-1}$ ). Since the same ZCS was used in both cases and it is highly unlikely for D<sub>2</sub>O to induce any changes in the electronic structure of ZCS, the depth of trap states should technically remain the same. Therefore, the charge carriers should also have similar electronic properties and should not result in such distinctive changes in the peak position.

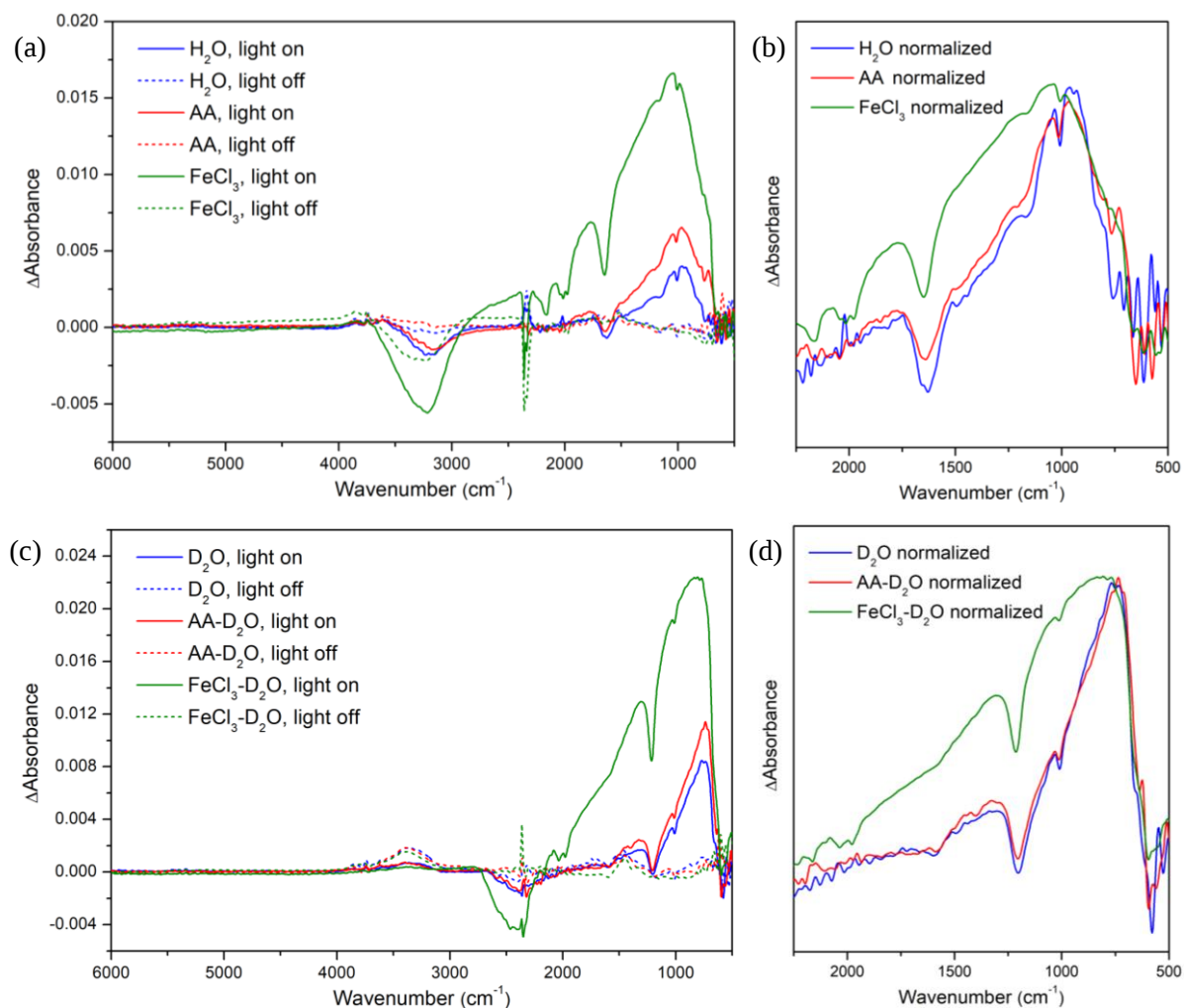
Considering that liquid H<sub>2</sub>O has a third characteristic absorption band of librational motions at  $\sim 750\text{ cm}^{-1}$ , it is suspected that intense bleaching also occurred at this position. As a result, negative peak similar to the ones at  $3180$  and  $1650\text{ cm}^{-1}$  was coincidentally superimposed onto the electronic absorption peak of ZCS in H<sub>2</sub>O, consequently instilling a ‘peak shift’ towards  $950\text{ cm}^{-1}$ . In contrast, since the libration band of D<sub>2</sub>O is at an even lower wavenumber region outside of the range in this work, the full electronic absorption band was displayed without any interruption from bleaching. Since the peaks in both H<sub>2</sub>O and D<sub>2</sub>O resemble with each other, it is believed that both responses arose from the same source, and thus both spectra can be used interchangeably depending on one’s need. Meanwhile, the nature of these carriers, whether are

they electrons or holes, is not possible to be identified by merely using water or D<sub>2</sub>O as the liquid medium. Scavenger solutions are employed next with the aim to configure them.

**IR absorption in scavenger solutions.** Figure 3 (a) presents the spectra obtained when ZCS was immersed in solutions of ascorbic acid (AA) and iron (III) chloride (FeCl<sub>3</sub>), in which they served as hole scavenger and electron scavenger respectively. The spectrum observed in pure water is included for comparison purposes. In AA aqueous solution, the photogenerated holes are prone to be consumed by AA, resulting in a prolonged lifetime of photoexcited electrons which are temporarily residing either in the conduction band or in the shallow trap states present in ZCS. This increases the amount of electrons that can absorb probing IR light and in turn translating to a larger  $\Delta$  absorbance as compared to when pure water was used, as elucidated in Figure 3 (a). From this observation, it can be tentatively deduced that this peak is associated to photoexcited electrons instead of holes, since the holes should have been consumed by AA.

When FeCl<sub>3</sub> solution was employed, surprisingly an electronic absorption peak with greatly enhanced intensity – much stronger as compared to AA and water – appeared at similar low wavenumber region when the illuminating light source was turned on. This contradicts with the deduction that has just been made, as the presence of FeCl<sub>3</sub> as an electron scavenger is expected to remove most of the photoexcited electrons and thus an electronic absorption band associated to these electrons is bound to not occur. In a previous study conducted by Fu et al.,<sup>19</sup> it was also reported that the usage of methanol and Na<sub>2</sub>SO<sub>3</sub> (both holes scavenger) in similar assessment of their Sr-doped NaTaO<sub>3</sub> photocatalyst resulted in an enhancement of such electronic absorption peak, but when FeCl<sub>3</sub> and NaIO<sub>3</sub> (both electrons scavenger) were used, the peak completely disappeared, proving that the peak that they acquired can be evidently ascribed to the bandgap-excited electrons.

Through a closer look into Figure 3 (a), it was observed that the absorption band formed in  $\text{FeCl}_3$  solution has a noticeable wider shoulder as compared to the bands acquired in AA and water. The shape difference becomes more apparent by normalizing the absorption peaks in the region of  $500 - 2900 \text{ cm}^{-1}$  to the same maximum as plotted in Figure 3 (b). The absorption bands obtained in AA and water appear to have similar shape to each other, whereas the band in  $\text{FeCl}_3$  appears relatively broader. Based on this observation, it is suspected that the distinctive band feature is an implication of the difference in nature of the photogenerated charge carriers, such that the broader absorption band represents the holes, while the narrower band is a contribution from electrons. It is not the first time this phenomenon is discovered; in fact, a report by Yamakata et al. showed that the width of electronic absorption bands assigned to photogenerated electrons and holes are evidently different in  $\text{LaTiO}_2\text{N}$  photocatalyst.<sup>17</sup> Despite so, the bands exhibited by electrons and holes usually fall onto two separate regions, and this is a huge contrast to the findings from the present study. Considering the distinctive properties of different photocatalyst, it is proposed that ZCS demonstrated an unprecedented electronic absorption where the peaks associated to photogenerated electrons and holes coincidentally appeared at the same wavenumber window albeit having dissimilar spectral widths. Moreover, as the band shape in pure water resembles to that when AA was used, this also infers that the behavior of pure water when ZCS was employed as photocatalyst inclines towards it being a holes scavenger, despite that theoretically water can be reduced and oxidized simultaneously by ZCS due to the bandgap of ZCS enveloping both water reduction and oxidation potentials. Essentially, it obliquely implies that the oxidation of water by ZCS is more prone to occur in a pure water environment without any external scavengers. This is an interesting deduction since ZCS has rarely been studied on its water oxidation activity, hence other supporting characterizations will need to be done in the future to validate this speculative statement.



**Figure 3.**  $\Delta$ Absorbance spectra of ZCS immersed in sacrificial solutions prepared with (a) H<sub>2</sub>O and (c) D<sub>2</sub>O. (b) and (d) are the respective normalized spectra of (a) and (c). The solid and dotted lines represent the spectra acquired with and without light irradiation correspondingly.

To verify the consistency of the present findings, similar measurements were performed with the D<sub>2</sub>O counterparts of each scavenger, namely AA dissolved in D<sub>2</sub>O (AA-D<sub>2</sub>O) and FeCl<sub>3</sub> dissolved in D<sub>2</sub>O (FeCl<sub>3</sub>-D<sub>2</sub>O). The result compiled in Figure 3 (c) showed a same trend as that when water was used, where pure D<sub>2</sub>O has the lowest peak intensity, followed by AA-D<sub>2</sub>O and FeCl<sub>3</sub>-D<sub>2</sub>O accordingly. Without the bleaching effect of UV on the vibrational absorption of water which was elaborated earlier, all peaks were present at a lower wavenumber region of  $\sim 750 \text{ cm}^{-1}$ . By comparing the normalized peaks in Figure 3 (d), the band obtained in FeCl<sub>3</sub>-D<sub>2</sub>O also tallied with the prior set of results that were acquired in water-based solutions, in

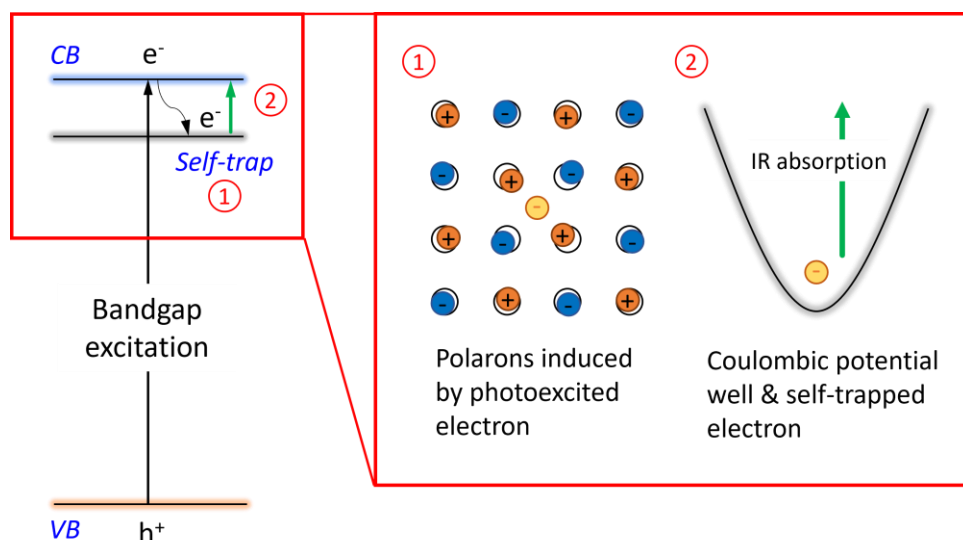
which it stood out with a much broader spectral shoulder among all spectra. As a side note, small peaks that appeared at  $\sim 3250\text{ cm}^{-1}$  and  $\sim 1450\text{ cm}^{-1}$  can be ascribed to O-H stretching mode arising from  $\text{H}_2\text{O}$  adsorbed on ZCS and H-O-D bending mode from the trace amount of HDO in the  $\text{D}_2\text{O}$  solution, respectively.<sup>38</sup>

Additionally, through Figure 3 (a) and (c), the electronic absorptions of ZCS in both variations of  $\text{FeCl}_3$  solutions ( $\text{H}_2\text{O}$  and  $\text{D}_2\text{O}$ -based) exhibited the strongest intensity, far ahead of the spectra taken in AA (AA- $\text{D}_2\text{O}$ ) or  $\text{H}_2\text{O}$  ( $\text{D}_2\text{O}$ ). It has been widely accepted that a strong absorption in this context signifies the relatively high quantity of charge carriers that are able to absorb the corresponding IR energy, even though the factors that affect this quantity can vary from case to case. Here, it is apparent that the contribution from photogenerated holes was substantially higher and this could be related to the scavenging ability of  $\text{FeCl}_3$  when paired with ZCS. In general, certain classes of photocatalytic materials have their ‘preferred’ scavenging solution where the resultant product yields are maximized due to the excellent scavenging abilities of the respective solutions. For example, thioacetamide (TEOA) as the best-in-slot (BIS) holes scavenger for  $\text{g-C}_3\text{N}_4$  polymeric photocatalyst,<sup>39–41</sup>  $\text{Na}_2\text{S}/\text{Na}_2\text{SO}_3$  or AA as the BIS holes scavenger for sulfide-based photocatalysts in which one may provide better outcome over another depending on the co-catalysts used,<sup>42–45</sup>  $\text{Na}_2\text{S}_2\text{O}_8$  as electron scavenger for ferrites such as  $\text{NiFeO}_x$  and  $\text{BiFeO}_3$ ,<sup>46,47</sup> and  $\text{AgNO}_3$  or  $\text{FeCl}_3$  as the more generic go-to electron scavenger in many of the studies using the popular oxygen evolution photocatalysts,  $\text{WO}_3$  and  $\text{BiVO}_4$ .<sup>48–54</sup> As the usage of ZCS for oxygen evolution is rather scarce, the combination of ZCS and  $\text{FeCl}_3$  as the photocatalyst-electron scavenger pair has most likely not been assessed before. Considering the occasion where  $\text{FeCl}_3$  has great synergy with ZCS and can effectively draw away the photoexcited electrons in ZCS, this will leave a substantial amount of holes in the valence band (VB) of ZCS that can thus contribute to strong electronic absorption during ATR-FTIR assessment as evidently portrayed in Figure 3 (a) and (c).

**IR Absorption mechanism.** Upon breaking down of the electronic absorption spectra that were successfully observed on ZCS in 3 different liquid environments, it is also important to dive deeper into the mechanism of these absorptions. As mentioned briefly before, these IR absorptions that occurred during light irradiation indicates the presence of shallow trap states. The photogenerated electrons (or holes) that were trapped momentarily in the trap states absorbed the probing IR beam to get re-excited to CB (or VB), hence exhibiting a response in the IR absorption spectrum. From a classic semiconductor point of view, trap states are often caused by the presence of physical defects which are introduced through modifications, for instance interstitial or substitutional doping of foreign atoms and stripping of local atoms resulting in atomic vacancies. Aside from that, it is also generally accepted that natural crystalline defects are inevitable during materials synthesis especially in wet-chemical procedures. In this study, the pristine ZCS photocatalyst used was synthesized via hydrothermal method without any form of additional modifications, thus it is safe to assume that any physical trap states present were results of inevitable natural defects that were induced during synthesis. In fact, in accordance with the observations in Section 3.2, these natural defects have fostered traps in close vicinity of both VB and CB as electronic absorption associated to both electrons and holes were observed.

Apart from the physical defect theory, the origin of the trap states can also be interpreted from the viewpoint of polarons. Here, the injection of charges into CB (or removal from VB) during photoexcitation instigated temporal dislocation of ionic atoms from their equilibrium positions due to polarizing effect, and this invokes Coulombic potential well that traps the photoexcited electrons (or holes). Such quasiparticle constituting of the foreign charged particle and its interaction with surrounding atoms is commonly known as a polaron. Contrasting to the physical defect theory where the trap states already exist in the crystal structure of ZCS prior to photoexcitation, the polaronic effect theory inferred that the traps are instantaneously self-

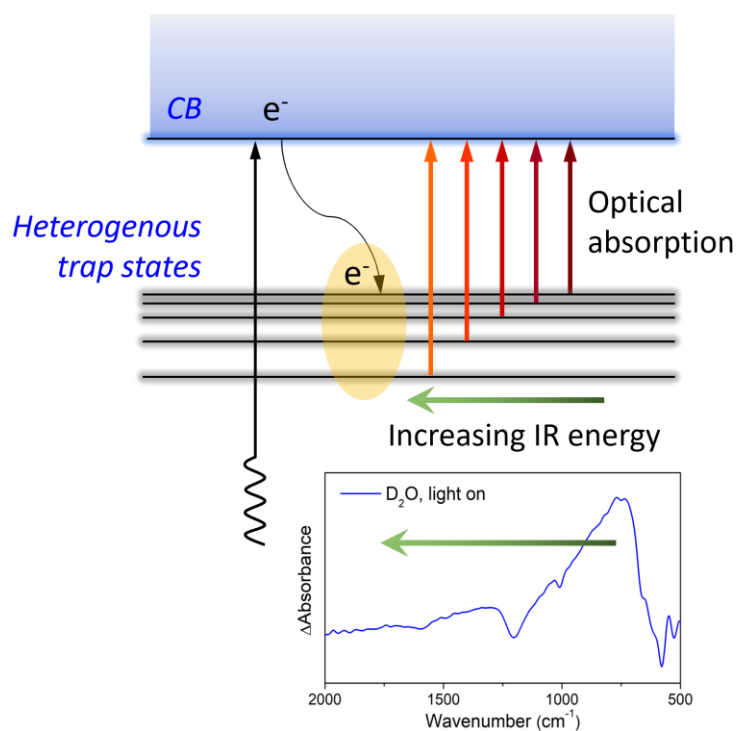
induced by the photogenerated charge carriers the moment when bandgap excitation is triggered. These traps are expected to fade off when the excitation light is turned off as the ionic atoms will return to their original equilibrium state. A sample illustration elucidating this occurrence and the self-trapping mechanism for photoexcited electrons is hereby included in Figure 4.



**Figure 4.** Self-trapping of electrons due to polaronic effect during photoexcitation. This illustration gives an example from the electrons perspective. Similar mechanism is applicable to holes at VB.

Following the two possible mechanisms of trap states formation, the attention is then shifted towards a striking feature in all absorption bands that were acquired, that is the asymmetry of the peaks. When a trap state is postulated to be present in the crystal structure of a photocatalyst, it is often regarded to be a state positioned on a discrete energy level, such that any movement of electrons (holes) between the trap and VB (CB) will require absorption of energy exactly equivalent to the energy gap between the two levels. As such, theoretically a symmetrical peak should be reflected in the IR absorption spectrum which is not the case here as seen in Figure 3. In fact, the asymmetrical bands with peaks that are skewed towards the lower wavenumber

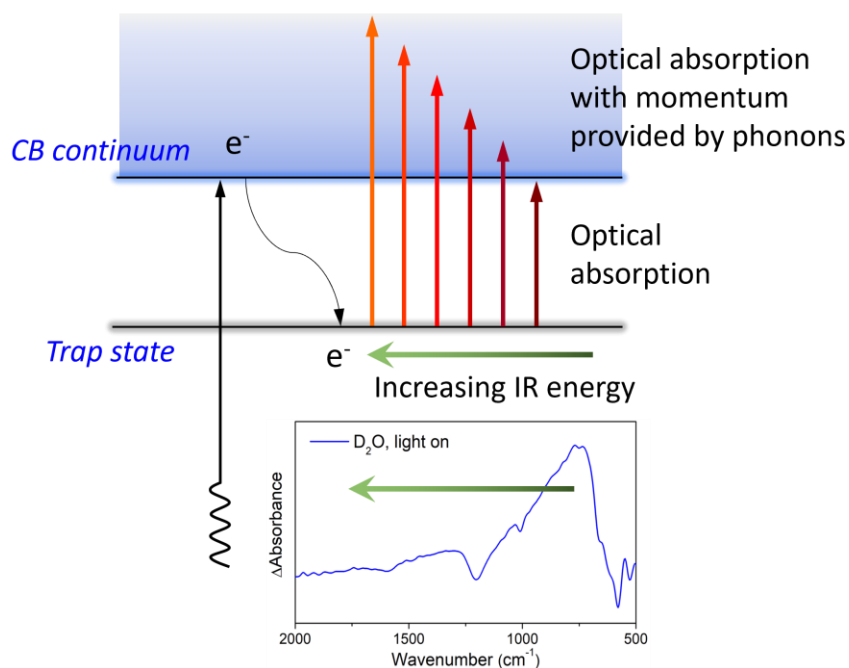
region suggest otherwise. Although the circumstance here can be ambiguous without further supporting evidence, one of the most probable reasons should be the formation of heterogeneous trap states, as depicted in Figure 5. For the case of physical traps, several levels of trap states can simply be formed just because the precise control of defects induction is usually not possible during synthesis. On another hand, when the polarons-induced traps are taken into consideration, different sizes of polarons may be manifested and it results in varied depth of Coulombic potential wells in which the charge carriers are trapped in. In either scenario, the energy level with the highest population of traps should be consistent with the position of the absorption peak, for instance 0.093 eV below the CB for electron traps (peak position of 750  $\text{cm}^{-1}$  from the narrower bands) and 0.10 eV above the VB for hole traps (peak position of 800  $\text{cm}^{-1}$  from the wider band), both enumerated from Figure 3 (c). As the energy gap increases, some trap states are also present, but the density of traps decreases correspondingly, therefore resulting in the skewed spectra.



**Figure 5.** Electrons trapped in the heterogenous trap states absorbing different IR energy to be re-excited to CB. Similar interpretation can be applied for the holes at VB.

Another postulation for the skewed spectra lays in the absorption of additional IR energy that enables the trapped charge carriers to move from their spatially localized states into the unbound continuum states within the CB or VB while facilitated by phonons. This can occur regardless of whether the charges are accommodated in physical defects or the ground states of the polaronic potential wells, and the resultant optical absorption band usually appears asymmetric.<sup>55</sup> Normally, the initial states of these carriers are confined in real space and its momentum is limited. Along with the absorption of photon energy, there are chances for the optical transition of carriers from the trap states into the CB (or VB) continuum, albeit the probability is however dependent on the electron momentum required to reach the final state delocalized in space. Since this momentum is provided by phonons in the host lattice, statistically the phonons with higher energy are much lower in abundance as compared to phonons with lower energy, hence the chances for charge carriers to transition into a deeper unbound state are lower. From this context, the low energy threshold of the absorption is determined by the energy gap between the trap state and the bottom of the continuum state. This generates a peak near the threshold, and the absorption coefficient then reduces gradually when the photon energy goes above the threshold, generating a skewed spectrum as seen in Figure 3. A scheme demonstrating this phenomenon is included as Figure 6. Note that it is possible for the heterogeneity in trap states and transition of trapped charges into unbound continuum states to occur concurrently, but for the scope of this study, the two scenarios are elaborated separately to reduce the complexity. In fact, while it is well accepted that trap states can have varied effects on the resultant photocatalytic kinetics; for instance, some might be effective temporary traps that prolong the lifetime of electrons and holes and some might be serving as recombination centres instead; from the present scope of study, it is not possible to tell the effectiveness of the traps described in the two scenarios above in terms of enhancing the efficiency of photocatalytic water splitting reaction. More in-depth studies on the kinetics

of these traps and the combined effect of both phenomena will be a noteworthy topic of research in the future especially with the aid of computational methods such as density functional theory.



**Figure 6.** Electrons excited beyond the lowest unoccupied state in CB into the unbound continuum states facilitated by phonons. Similar interpretation can be applied for the holes at VB.

## CONCLUSIONS

In essence, the in-situ response of IR probing on ZCS immersed in aqueous environment during photoirradiation was successfully observed and reported for the first time. The  $\Delta$ Absorbance peak that appeared only during irradiation was ascribed to electronic absorption, which reflects the phenomenon of IR energy absorption by the photogenerated charge carriers that were momentarily trapped in the respective trap states. Through measurements in scavenging solutions, it was revealed that the contribution from electrons and holes could be distinguished based on the qualitative difference of the acquired peak shape, although coincidentally both contributions appeared in a similar wavenumber window, which was unheard of prior to this

study. Specifically, the electronic absorption by photogenerated electrons in ZCS was recognized by the relatively narrower absorbance band when compared with the absorption caused by holes. With the presence of these responses, it was divulged that trap states were present in close vicinity of both valence and conduction bands. The potential origins of these traps were also laid out accordingly based on two different viewpoints. From the perspective of classic semiconductive materials, natural defects that were inherently formed in the lattice of ZCS during the hydrothermal synthesis process could have introduced some physical trap states. On another hand, when the photogenerated charge carriers were injected into the valence or conduction bands, polarons that were instantaneously created have also caused disruption to the surrounding ionic atoms in their originally equilibrium states, and this led to the formation of Coulumbic potential wells that can entrap the carriers. Notably, based on the asymmetry in the observed absorbance band, it was suspected that there were some heterogeneities in the trap states, in which different density of traps were introduced at multiple levels, or the possibility that the carriers gained different magnitudes of momentum from the phonons to be projected into the unbound continuum states. All in all, this study opens a new avenue for the characterization of sulfide photocatalysts by highlighting the feasibility of ATR-FTIR to analyze the photogenerated charge carriers instantaneously during light irradiation.

## ASSOCIATED CONTENT

**Supporting Information.** The supporting information is available free of charge on the ACS publications website as a PDF.

XRD pattern (Figure S1) and FE-SEM image (Figure S2) of ZCS.

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The authors contributed to the development of the manuscript.

### Notes

The authors declare no competing financial interest.

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