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Lotka-Volterra system and KCC theory: Differential geometric structure of competitions and predations

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Abstract

We consider the differential geometric structure of competitions and predations in the sense of the Lotka-Volterra system based on KCC theory. For this, we visualise the relationship between the Jacobi stability and the linear stability as a single diagram. We find the following. (I) Ecological interactions such as competition and predation can be described by the deviation curvature. In this case, the sign of the deviation curvature depends on the type of interaction, which reflects the equilibrium point type. (II) The geometric quantities in KCC theory can be expressed in terms of the mean and Gaussian curvatures of the potential surface. In this particular case, the deviation curvature can be interpreted as the Willmore energy density of the potential surface. (III) When the equations of the system have nonsymmetric structure for the species (e.g. a predation system), each species also has nonsymmetric geometric structure in the nonequilibrium region, but symmetric structure around the equilibrium point. These findings suggest that KCC theory is useful to establish the geometrisation of ecological interactions.

Keywords: Lotka-Volterra system, KCC theory, Competitions, Predations, Jacobi stability, Curvatures

1. Introduction

2 The geometrisation of nature has been the subject of theoretical interest.
3 This includes the use of geometric concepts and techniques in the natural

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4 sciences, such as the general theory of relativity in physics (e.g. [1]). In bi-
5 ology, the geometrisation of nature is also of scientific and practical interest.
6 For example, the theory of Kosambi-Cartan-Chen (KCC theory: the geomet-
7 ric theory of the stability of dynamic systems [2, 3, 4]) has been applied to
8 several biological problems, such as production in the Volterra method [5, 6],
9 the Volterra-Hamilton system [7, 8, 9], Tyson's model for the cell division
10 cycle [10] and the robustness of biological systems [11].

11 The general theory of relativity is well known to geometrize the interac-
12 tion between masses. Can we also geometrize competitions and predations,
13 i.e. the interaction between living things? The main purpose of this pa-
14 per is to consider the differential geometric structure of the Lotka-Volterra
15 system based on KCC theory. We are concerned with (I) the nature of the
16 competitions (intraspecific vs. interspecific) and (II) predations between the
17 prey and the predator. Even though these are the typical ecological inter-
18 actions described by the Lotka-Volterra system, few geometric studies have
19 investigated these concrete cases.

20 The structure of this paper is as follows. In Section 2, we give a brief
21 review of KCC theory. It has been applied to various nonlinear dynamic sys-
22 tems (e.g. [12, 13, 14, 15, 16, 17]) because it clarifies their intrinsic properties
23 using differential geometric concepts such as "connections" and "curvatures".
24 In Section 3, we consider the geometric structure of the competitions and pre-
25 dations described by the Lotka-Volterra system. In this analysis, we use the
26 equation for the differential geometric quantities in terms of the Jacobian
27 matrix of the linearised system. Based on this result, we discuss the ecolog-
28 ical meaning of the geometric expression for the interactions between living
29 things in Section 4. This will give new insights into the effect of the curva-
30 tures of the potential surface on the stability of the system. We also discuss
31 the nonequilibrium case and the bifurcations in KCC theory. Section 5 is
32 devoted to the conclusions.

33 2. Brief review of KCC theory and Jacobi stability

34 In this section, we give a brief review of KCC theory and Jacobi sta-
35 bility based on previous papers [7, 15]. Throughout this paper, Einstein's
36 summation convention is used.

37 Let M be a real smooth n -dimensional manifold, and (TM, π, M) be its
38 tangent bundle, where $\pi : TM \rightarrow M$ is a projection from the total space TM
39 to the base manifold M . A point $x \in M$ has local coordinates (x^i) , where

40 $i = 1, \dots, n$. The local chart of a point in TM is denoted by (x^i, \dot{x}^i) , where
 41 t is time (regarded as an absolute invariant) and $\dot{x}^i = dx^i/dt$.

42 Let us consider the path equation

$$\ddot{x}^i + g^i(x, \dot{x}) = 0, \quad (1)$$

43 where $g^i(x, \dot{x})$ is a smooth function. The trajectory $x^i(t)$ of the system (1)
 44 is changed to nearby trajectories according to $\bar{x}^i = x^i + u^i \delta\tau$, where u^i is a
 45 vector field and $\delta\tau$ is a small parameter. In this case, Eq. (1) becomes the
 46 variational equation for the limit $\delta\tau \rightarrow 0$:

$$\ddot{u}^i + \frac{\partial g^i}{\partial x^j} u^j + \frac{\partial g^i}{\partial \dot{x}^j} \dot{u}^j = 0. \quad (2)$$

47 The covariant form of (2) is given by

$$\frac{D^2 u^i}{Dt^2} = P_j^i u^j, \quad (3)$$

48 where $D(\dots)/Dt$ is a covariant differential defined by

$$\frac{Du^i}{Dt} = \frac{du^i}{dt} + N_j^i u^j, \quad (4)$$

49 N_j^i is a coefficient of the nonlinear connection

$$N_j^i = \frac{1}{2} \frac{\partial g^i}{\partial \dot{x}^j}, \quad (5)$$

50 P_j^i is the deviation curvature tensor

$$P_j^i = -\frac{\partial g^i}{\partial x^j} + \frac{\partial N_j^i}{\partial x^k} \dot{x}^k - G_{jk}^i g^k + N_k^i N_j^k, \quad (6)$$

51 and G_{jk}^i is a Finsler (Berwald) connection

$$G_{jk}^i = \frac{\partial N_j^i}{\partial \dot{x}^k}. \quad (7)$$

52 The first term of (6): $\partial g^i / \partial x^j$ is the curvature when all the coefficients
 53 of connections become zero. In this paper, we tentatively call it the zero-
 54 connection curvature

$$Z_j^i = \frac{\partial g^i}{\partial x^j}. \quad (8)$$

55 As we will see in Section 4.2, the zero-connection curvature corresponds to
 56 the Gaussian curvature of the potential surface.

57 The deviation curvature tensor P_j^i gives the stability of whole trajectories
 58 via the following theorem [5]: *The trajectories of the system (1) are Jacobi*
 59 *stable if and only if the real parts of the eigenvalues of P_j^i are strictly negative*
 60 *everywhere, and Jacobi unstable otherwise.* In particular, the trajectories
 61 of the one-dimensional system are Jacobi stable when $P_1^1 \leq 0$, and Jacobi
 62 unstable when $P_1^1 > 0$.

63 3. Differential geometric quantities for competitions and preda- 64 tions

65 In this section, we derive the deviation curvature for competitions and
 66 predations in the sense of the Lotka-Volterra system. In previous studies, a
 67 linear stability analysis has often been applied to the Lotka-Volterra system.
 68 Linear stability analysis is the theory of local stability around a point on
 69 the tangent space, which means that the behaviour of the nonlinear dynamic
 70 systems is described by the tangent bundle. In this case, the equation (1) is
 71 a first-order differential equation with respect to \dot{x}^i , and the Jacobi stability
 72 equation (2) is reduced to an equation in a linear stability theory (e.g. [17]).
 73 Therefore, the Jacobi stability gives a more global stability than the linear
 74 stability (see also [13, 14]).

75 In linear stability analysis, the Jacobian matrix of the linearised system
 76 plays an important role. To clarify the relationship between the Jacobi sta-
 77 bility and the linear stability in a linearised system, we express the geometric
 78 quantities of KCC theory in terms of the Jacobian matrix of the linearised
 79 system (see also [11, 18]).

80 3.1. Differential geometric quantities in terms of the Jacobian matrix

81 We consider a vector field described by

$$\dot{x}^i = f^i(x), \quad (9)$$

82 where $i = 1, 2, \dots, n$ and f^i denote a given function. This can be approxi-
 83 mated by a linear system around an equilibrium point x_0^i using the relation
 84 $x^i = x_0^i + \xi^i$, where ξ^i is a small quantity. That is,

$$\dot{\xi}^i = J_j^i(x_0)\xi^j, \quad (10)$$

85 where $J_j^i(x_0)$ is the Jacobian matrix of f^i .

86 In this paper, we consider the two-dimensional case ($i = 1, 2$):

$$\dot{\xi} = J(x_0)\xi, \quad (11)$$

87 where

$$\xi = \begin{pmatrix} \xi^1 \\ \xi^2 \end{pmatrix}, \quad (12)$$

$$J(x_0) = J = \begin{pmatrix} \partial_1 f^1(x_0) & \partial_2 f^1(x_0) \\ \partial_1 f^2(x_0) & \partial_2 f^2(x_0) \end{pmatrix}, \quad (13)$$

88 and $\partial_i = \partial(\dots)/\partial x^i$. The simultaneous differential equation (11) can be
 89 rewritten as a second-order ordinary differential equation. When we consider
 90 the coordinate system $(\xi^i, \dot{\xi}^i)$, we have the following equation for $i = 1$:

$$\ddot{\xi}^1 - \text{tr}[J]\dot{\xi}^1 + \det[J]\xi^1 = 0. \quad (14)$$

91 This is a particular case of (1) for $g^1 = -\text{tr}[J]\dot{\xi}^1 + \det[J]\xi^1$. Therefore, Eqs.
 92 (5), (8) and (6) give the nonlinear connection, the zero-connection curvature
 93 and the deviation curvature of the linearised system, respectively:

$$N = -\frac{1}{2}\text{tr}[J], \quad (15)$$

$$Z = \det[J], \quad (16)$$

$$P = \frac{1}{4}\{(\text{tr}[J])^2 - 4\det[J]\} = N^2 - Z, \quad (17)$$

94 where $N_1^1 = N$, $Z_1^1 = Z$ and $P_1^1 = P$ (we use the same abbreviation through-
 95 out this paper). From Eq. (7), the Finsler connection vanishes in this lin-
 96 earised system (the nonvanishing case will be shown in Section 4.3). Now,
 97 Eqs. (15), (16) and (17) show that the geometric quantities of the linearised
 98 system can be easily calculated when the Jacobian matrix of the system is ob-
 99 tained. Eq. (17) has already been derived in previous papers (e.g. [11, 18]).
 100 In the next section, we derive the differential geometric quantities describing
 101 competitions and predations from these equations.

102 *3.2. Lotka-Volterra competition system*

103 As an example of (9), we consider the Lotka-Volterra competition system
 104 for two-species: $x^1 = x$ and $x^2 = y$. The standard form of the system is
 105 given by (e.g. [19])

$$\dot{x} = r_1 x \left(1 - \frac{x + a_1 y}{k_1}\right), \quad (18)$$

$$\dot{y} = r_2 y \left(1 - \frac{y + a_2 x}{k_2}\right), \quad (19)$$

106 where r_i is the natural growth rate, k_i is the carrying capacity and a_i is the
 107 competition coefficient. These parameters are all positive.

108 This system can be approximated by a linear system around an equilib-
 109 rium point (x_0, y_0) . From (13), the Jacobian matrix of the system is

$$J = \begin{pmatrix} r_1 - 2(r_1/k_1)x_0 - a_1(r_1/k_1)y_0 & -a_1(r_1/k_1)x_0 \\ -a_2(r_2/k_2)y_0 & r_2 - 2(r_2/k_2)y_0 - a_2(r_2/k_2)x_0 \end{pmatrix}. \quad (20)$$

110 Next, we consider the deviation curvature in two cases: 1) two species co-
 111 exist and 2) only one species survives. Because the standard form for the
 112 competition does not allow extinction of the two species except for under the
 113 trivial condition, we ignore this case.

114 In the case when the two species coexist, i.e. $x_0 \neq 0$ and $y_0 \neq 0$, we have
 115 $k_1 = x_0 + a_1 y_0$ and $k_2 = y_0 + a_2 x_0$. Therefore, (20) becomes

$$J = \begin{pmatrix} A_1 & a_1 A_1 \\ a_2 A_2 & A_2 \end{pmatrix}, \quad (21)$$

116 where $A_1 = -r_1 x_0 / k_1$ and $A_2 = -r_2 y_0 / k_2$. From (15) and (16), the nonlinear
 117 connection and the zero-connection curvature are given by $N = -(A_1 +$
 118 $A_2)/2$ and $Z = (1 - a_1 a_2) A_1 A_2$, respectively. Then, from (17), the deviation
 119 curvature for the two species coexisting is

$$P = \frac{1}{4} \{ (A_1 - A_2)^2 + 4a_1 a_2 A_1 A_2 \}. \quad (22)$$

120 We discuss the ecological meaning of (22) in Section 4.1.

121 In the case when only x survives, i.e. $x_0 = k_1$ and $y_0 = 0$, (20) becomes

$$J = \begin{pmatrix} -r_1 & -a_1 r_1 \\ 0 & C \end{pmatrix}, \quad (23)$$

where $C = r_2(1 - a_2 k_1/k_2)$. In this case, since $N = -(C - r_1)/2$ and $Z = -r_1 C$, the deviation curvature for only x surviving is

$$P = \frac{1}{4}(C - r_1)^2 + r_1 C. \quad (24)$$

For the Jacobi matrix to be stable, P should be negative or zero. This requires at least $C \leq 0$, i.e. $k_2 - a_2 k_1 \leq 0$. This condition also guarantees linear stability, as the zero isoclines method indicates (e.g. [19]).

3.3. Lotka-Volterra predation system

We consider the Lotka-Volterra predation system for the prey $x^1 = x$ and predator $x^2 = y$. The standard form of the system is given by (e.g. [19, 20])

$$\dot{x} = rx - axy, \quad (25)$$

$$\dot{y} = bxy - cy, \quad (26)$$

where r is the natural growth rate of the prey, a and b are coefficients of predation and c is the natural death rate of the predator. These parameters are all positive.

This system can be approximated by a linear system around an equilibrium point (x_0, y_0) . From (13), the Jacobian matrix of the system is

$$J = \begin{pmatrix} r - ay_0 & -ax_0 \\ by_0 & bx_0 - c \end{pmatrix}. \quad (27)$$

Next, we consider the deviation curvature in two cases: 1) the two species coexist and 2) extinction. Because the standard form of the predation does not include the case of only x or y surviving, we ignore this case.

In the case when the two species coexist, i.e. $x_0 \neq 0$ and $y_0 \neq 0$, we have $x_0 = c/b$ and $y_0 = r/a$. Therefore, (27) becomes

$$J = \begin{pmatrix} 0 & -ac/b \\ br/a & 0 \end{pmatrix}. \quad (28)$$

140 From (15) and (16), the nonlinear connection and the zero-connection cur-
 141 vature are given by $N = 0$ and $Z = cr$. Then, from (17), the deviation
 142 curvature of coexistence is

$$P = -cr. \quad (29)$$

143 We discuss the ecological meaning of (29) in Section 4.1.

144 In the case of extinction, i.e. $x_0 = y_0 = 0$, (27) becomes

$$J = \begin{pmatrix} r & 0 \\ 0 & -c \end{pmatrix}. \quad (30)$$

145 In this case, since $N = -(r - c)/2$ and $Z = -cr$, the deviation curvature of
 146 extinction is

$$P = \frac{1}{4}(r - c)^2 + rc. \quad (31)$$

147 This is always positive, i.e. the extinction state is the Jacobi unstable.

148 4. Discussion

149 4.1. Geometric structure of competitions and predations: Jacobi stability and 150 linear stability

151 In this paper, we have considered the geometric meaning of competition
 152 and predation based on KCC theory. In the case of the two species coexisting,
 153 the deviation curvatures of competition and predation were given by

$$P = \frac{1}{4}\{(A_1 - A_2)^2 + 4a_1a_2A_1A_2\} > 0, \quad (32)$$

$$P = -cr < 0, \quad (33)$$

154 where $A_1 = -r_1x_0/k_1$ and $A_2 = -r_2y_0/k_2$. These equations show that the
 155 deviation curvature of the competition and predation are always positive
 156 and negative, respectively. That is, the competition and the predation have
 157 essentially distinct geometric structures.

158 The ecological meaning of this result can be understood if we return to
 159 the equations for the geometric quantities in terms of the Jacobian matrix,
 160 which were derived in Section 3.1:

$$N = -\frac{1}{2}tr[J], \quad (34)$$

$$Z = det[J], \quad (35)$$

$$P = \frac{1}{4}\{(tr[J])^2 - 4det[J]\}. \quad (36)$$

161 As indicted in Section 2, the left term of Eq. (36) is related to the Jacobi
162 stability, i.e.

$$\begin{cases} \text{Jacobi stable for } P \leq 0, \\ \text{Jacobi unstable for } P > 0. \end{cases} \quad (37)$$

163 On the other hand, the right term of Eq. (36) is related to the linear stability.
164 Let Δ be a discriminant of the characteristic polynomial of J : $\Delta = (tr[J])^2 -$
165 $4det[J]$. In this case, the classification of an equilibrium point as shown on a
166 single diagram such as Fig. 1(a) is well known (e.g. [21]). From (34) and (35),
167 this diagram gives the relationship between the Jacobi stability and the linear
168 stability, as shown in Fig. 1(b). Fig.1(b) shows that the deviation curvature
169 corresponds to the discriminant, i.e. the Jacobi stability distinguishes the
170 node-saddle system and the spiral-centre system, as pointed out by Sabău
171 [11, 18]. Fig. 1(b), however, also shows that the linear stability is related
172 to the nonlinear connection but not to the deviation curvature. That is, the
173 system is

$$\begin{cases} \text{linear stable for } N > 0, \\ \text{linear unstable for } N < 0. \end{cases} \quad (38)$$

174 The coexistence of predation systems often results in periodic variations
175 in populations (e.g. [22, 23]). This corresponds to the spirals in the linear
176 stability theory, so the deviation curvature of the predations should be nega-
177 tive at least. This agrees with Eq. (33). On the other hand, the coexistence
178 state of the competition system is not periodic, i.e. no nodes exist. There-
179 fore, the deviation curvature of the competition should be positive at least,
180 in agreement with Eq. (32).

181 The zero-connection curvature of the competitions derived in Section 3.2
182 ($Z = (1 - a_1a_2)A_1A_2$) allows geometric interpretation of the intra- and inter-
183 species competition as follows. When the interspecific competition is more
184 intense than the intraspecific competition, i.e. $a_1a_2 > 1$, we have negative Z .

185 However, when the intraspecific competition is more intense, i.e. $a_1 a_2 < 1$,
 186 we have positive Z . Because $P > 0$ and $N > 0$ in the competition system,
 187 the latter case ($Z > 0$) is Jacobi unstable and in the linear stable region of
 188 Fig. 1(b), i.e. the stable nodes. The former case ($Z < 0$) is in the Jacobi
 189 unstable region, i.e. the saddle points. These results show that the sign of
 190 geometric quantities in KCC theory can describe the type of the ecological
 191 interactions.

192 4.2. Geometric structure of potential surface and stability

193 As indicated in the previous section, Eqs. (34), (35) and (36) are useful
 194 for analysing the geometric structure of the ecological system. The left-hand
 195 sides of these equations are geometric quantities, but the right hand sides are
 196 not. To fully geometrize the ecosystem, we must try to express these equa-
 197 tions as purely geometric relations. For this, we introduce the concept of a
 198 potential surface. Because ecological systems have often been analysed using
 199 the Lyapunov function, the results of this discussion are expected to show
 200 how the geometric restriction of the potential surface affects the stability of
 201 the ecological system.

202 We define the potential $V = V(x^i)$ for the vector field in Eq. (9) as

$$\dot{x}^i = f^i(x) = -\partial_i V(x). \quad (39)$$

203 In the two-dimensional case, the necessary and sufficient condition for the
 204 existence of the potential $V = V(x^1, x^2)$ is $\partial_2 f^1 = \partial_1 f^2$. In this case, the
 205 Jacobian matrix (13) becomes

$$J = \begin{pmatrix} -\partial_1 \partial_1 V & -\partial_2 \partial_1 V \\ -\partial_1 \partial_2 V & -\partial_2 \partial_2 V \end{pmatrix}. \quad (40)$$

206 Therefore, $tr[J]$ and $det[J]$ correspond to the mean curvature H and Gaussian
 207 curvature K of the potential surface, respectively.

$$tr[J] = -(\partial_1 \partial_1 V + \partial_2 \partial_2 V) = -2H, \quad (41)$$

$$det[J] = (\partial_1 \partial_1 V)(\partial_2 \partial_2 V) - (\partial_1 \partial_2 V)^2 = K. \quad (42)$$

208 Eqs. (34), (35) and (36) can therefore be expressed as purely geometric
 209 relations:

$$N = H, \quad (43)$$

$$Z = K, \quad (44)$$

$$P = H^2 - K. \quad (45)$$

210 Eqs. (43) and (44) show the geometric quantities of KCC theory: N and
 211 Z correspond to H and K in classic differential geometry, respectively. As
 212 Fig. 1 indicates, the classification of the equilibrium points is determined
 213 by N and Z . On the other hand, H and K are fundamental quantities that
 214 characterise the surface. Therefore, this correspondence describes the effect
 215 of the geometric restrictions on the stability of the system.

216 What is the meaning of the remaining equation (45)? When the potential
 217 surface is a smooth closed surface embedded in three-dimensional space, we
 218 can define the Willmore energy W as [24, 25]:

$$W = \int_V (H^2 - K) dA, \quad (46)$$

219 where dA is the area form of V . In this case, Eqs. (45) and (46) show that
 220 the deviation curvature of KCC theory corresponds to a kind of Willmore
 221 energy density per unit area.

222 When the potential surface is minimal, i.e. $H = 0$, we have $P = -K$.
 223 This shows that the sign of the deviation curvature (i.e. Jacobi stable or
 224 unstable) depends on the Gaussian curvature. Moreover, if we can define
 225 the Willmore energy, the integration of the deviation curvature is expressed
 226 only by the topological invariant, following the Gauss-Bonnet theorem. For
 227 example, this is the case of a velocity potential surface with zero vortex [26].
 228 In this case, Okubo-Weiss's Q-value Q corresponds to the deviation curvature
 229 P , since the Q-value is defined as the negative Gaussian curvature: $Q = -K$
 230 [27]. In fact, Okubo-Weiss's Q-value has been interpreted as the stability of
 231 the trajectory of Lagrangian particles immersed in a velocity field [28, 29].

232 Eq. (40) means that the eigenvalues of the Jacobi matrix are the principal
 233 curvature of the potential surface. Since the principal curvature in ordinary
 234 differential geometry is expressed in terms of real numbers, we assume a
 235 condition for the eigenvalues to be real numbers: $H^2 - K \geq 0$. From (45),
 236 this means that $P \geq 0$. In this case, the application of (44) and (45) to Fig.

1 gives the following correspondence:

- (1) $P > 0$,
 - (1-1) $K > 0$ (elliptic) \longleftrightarrow nodes,
 - (1-2) $K < 0$ (hyperbolic) \longleftrightarrow saddle,
 - (1-3) $K = 0$ (parabolic) \longleftrightarrow non-isolated fixed points,
- (2) $P = 0 \longleftrightarrow$ degenerate nodes.

These correspondences show that the deviation curvature P and the Gaussian curvature K are related to the type of equilibrium point. On the other hand, Eqs. (38) and (43) show that the mean curvature H is related to the stability of the equilibrium point: the system is stable for $H > 0$ and unstable for $H < 0$.

We now discuss the effect of the geometric restriction on the stability. The Hartman-Grobman theorem (e.g. [21]) states that *the local phase portrait near a hyperbolic equilibrium point is topologically equivalent to the phase portrait of the linearisation; in particular, the stability type of the equilibrium point is faithfully captured by the linearisation*. When the equilibrium point is hyperbolic, the Gaussian curvature of the potential surface is negative, so Eq. (45) shows that the deviation curvature is always positive, i.e. the system is always Jacobi unstable. According to the theorem, this is captured by the linearisation. For example, when the potential surface is the minimal surface ($H = 0$ and $K < 0$), the stability type of the equilibrium point is captured by the linearisation. However, if the equilibrium point is non-hyperbolic, the Gaussian curvature is more than zero, so the sign of the deviation curvature (i.e. the Jacobi stability) is not uniquely determined. For example, when the potential surface is elliptic ($K > 0$), the stability type cannot be uniquely determined.

4.3. Nonequilibrium case of predations and competitions

In this paper, we have studied small perturbations at the equilibrium point. However, KCC theory can treat small perturbations at any point. In this section, we discuss the deviation curvature in the nonequilibrium case for the predation and the competition systems.

First, we consider the predation system. To use Eq. (6) based on Eq. (1), we rewrite Eqs. (25) and (26) as the equation for the predator $y = x^2$ by vanishing $x = x^1$:

$$\ddot{y} + \gamma = 0, \quad (47)$$

with

$$\gamma = -\frac{\dot{y}^2}{y} + (ay - r)\dot{y} + acy^2 - cry, \quad (48)$$

where $\gamma = g^2$. To derive this equation, we assume that $y \neq 0$, so the following case is that of coexistence. From (5), (7) and (8), we have $N_2^2 = (1/2)\partial\gamma/\partial\dot{y} = -\dot{y}/y + (ay - r)/2$, $G_{22}^2 = -1/y$ and $Z_2^2 = \dot{y}^2/y^2 + a\dot{y} + 2acy - cr$. Therefore, from (6), the deviation curvature of the predator in the nonequilibrium case can be obtained:

$$P_2^2 = -\frac{a}{2}\dot{y} - acy + \frac{a^2}{4}\left(y - \frac{r}{a}\right)^2. \quad (49)$$

This is derived from the equation for the predator (47). In a similar fashion, we can derive the deviation curvature of the prey in the nonequilibrium case:

$$P_1^1 = \frac{b}{2}\dot{x} - brx + \frac{b^2}{4}\left(x - \frac{c}{b}\right)^2. \quad (50)$$

The comparison between Eq. (49) and Eq. (50) shows that the predator and prey have distinct geometric structure in the nonequilibrium case because the original equations (25) and (26) are not symmetric for the predator and prey. However, the predator and prey in the equilibrium case ($y = y_0 = r/a$ and $x = x_0 = c/b$) have the same geometric structure, $P_1^1 = P_2^2 = -cr$, in agreement with Eq. (29). This implies that the equilibrium point can be interpreted as the point in which the two deviation curvatures become equal, although the original equations have a nonsymmetric variable structure. In other words, the nonsymmetric structure inherent in the system develops over the geometric structure in the nonequilibrium region. As noted in Sections 4.1 and 4.2, the geometric structure of the dynamic system is related to the stability of the system, which means that the system has a different stability structure for each species in the nonequilibrium region.

Next, we consider the competition system (18) and (19). In a similar fashion to that described above, we can obtain the deviation curvature of the competition in the nonequilibrium case:

$$P_1^1 = \left(-\frac{1}{2k_1} + \frac{1}{a_1k_2} - \frac{a_2}{2k_2} \right) \dot{x} + \left(\frac{1}{4k_1^2} + \frac{1}{a_1k_1k_2} - \frac{3a_2}{2k_1k_2} + \frac{a_2^2}{4k_2^2} \right) x^2 + \left(\frac{1}{2k_1} - \frac{1}{a_1k_2} + \frac{a_2}{2k_2} \right) x + \frac{1}{4}, \quad (51)$$

where we assume that $x \neq 0$. Moreover, we assume that $r_1 = r_2 = 1$ because the natural growth rate is not related to the stability analysis. The deviation curvature for species y has the same structure because the original equations (18) and (19) are symmetric for x and y . When the two species coexist $x = x_0 = (k_1 - a_1k_2)/(1 - a_1a_2)$, in agreement with Eq. (22). When only x survives, $x = x_0 = k_1$, in agreement with Eq. (24). These results mean that nonlinear analysis gives the same result as linear analysis around the equilibrium point.

4.4. Extended predation system: bifurcations and KCC theory

The stability according to linear theory plays an important role in bifurcation theory. How is the Jacobi stability of KCC theory related to bifurcation theory? Here we consider a simple example.

To consider bifurcation, we modify Eqs. (25) and (26) as

$$\dot{x} = rx \left(1 - \frac{x}{k} \right) - \frac{axy}{1 + hx}, \quad (52)$$

$$\dot{y} = \frac{bxy}{1 + hx} - cy, \quad (53)$$

where the new parameter k is the carrying capacity of the prey, and h is the saturation of the predation rate. When $b > ch$, the equilibrium population of the coexistence case is $(x_0, y_0) = (c/(b - ch), (r/a)(1 - x_0/k)(1 + hx_0))$. According to linear stability analysis, the critical point is $x_C = (1/2)(k - 1/h)$, and the system is linearly stable for $x_0 > x_C$ and linearly unstable for $x_0 < x_C$.

We reconsider this bifurcation from the viewpoint of KCC theory. The Jacobian matrix of Eqs. (52) and (53) is

$$J = \begin{pmatrix} \frac{2rhx_0}{k(1 + hx_0)}(x_C - x_0) & \frac{-ax_0}{1 + hx_0} \\ \frac{by_0}{(1 + hx_0)^2} & 0 \end{pmatrix}. \quad (54)$$

319 From (15), (16) and (17), we have

$$N = -\frac{rhx_0}{k(1+hx_0)}(x_C - x_0), \quad (55)$$

$$Z = \frac{abx_0y_0}{(1+hx_0)^3}, \quad (56)$$

$$P = N^2 - Z. \quad (57)$$

320 Eq. (55) shows that the sign of the nonlinear connection is changed due
 321 to the bifurcation. However, Eq. (57) shows that the bifurcation does not
 322 affect the sign of the deviation curvature. In KCC theory, the sign of the
 323 deviation curvature has been of considerable concern because this is related
 324 to the Jacobi stability. However, this result means that sign of the other
 325 geometric quantities such as N and Z should receive attention because it is
 326 expected to be related to the bifurcations. For example, Fig. 1 implies that
 327 the nonlinear connection is related to bifurcations such as the Hopf type, and
 328 the zero-connection curvature is related to bifurcations such as the saddle-
 329 node type. As a subject of future study, it would be interesting to consider
 330 how differential geometric structures are related to the type of bifurcation.

331 5. Conclusions

332 Our main conclusions are as follows.

- 333 1. The geometric quantities of KCC theory can be expressed in terms of
 334 the Jacobian matrix of the linearised system. This clarifies the rela-
 335 tionship between the Jacobi stability and the linear stability, as shown
 336 in Fig. 1.
- 337 2. The competitions and predictions described by Lotka-Volterra system
 338 show a distinct geometric structure. This reflects the type of equilib-
 339 rium point in the sense of the ecology.
- 340 3. When the potential of a dynamic system can be defined, the nonlin-
 341 ear connection and the zero-connection curvature correspond to the
 342 mean and Gaussian curvature of the potential surface, respectively. In
 343 this case, the deviation curvature is a function of only one geometric
 344 concept: the curvature. This means that the geometric structure of
 345 the potential surface restricts not only the linear stability but also the
 346 Jacobi stability.

- 347 4. When the equations of the system have nonsymmetric structures for the
348 species (e.g. the predation system), each species also has nonsymmet-
349 ric geometric quantities in the nonequilibrium region, but symmetric
350 quantities around the equilibrium point. This means that KCC theory
351 is useful for considering the development of the nonsymmetric system.
- 352 5. In KCC theory, the deviation curvature has been of considerable con-
353 cern. However, as the extended predation system implies, other geo-
354 metric quantities such as the nonlinear connection and the zero-connection
355 curvature are also important because they may be related to bifurcation
356 theory.

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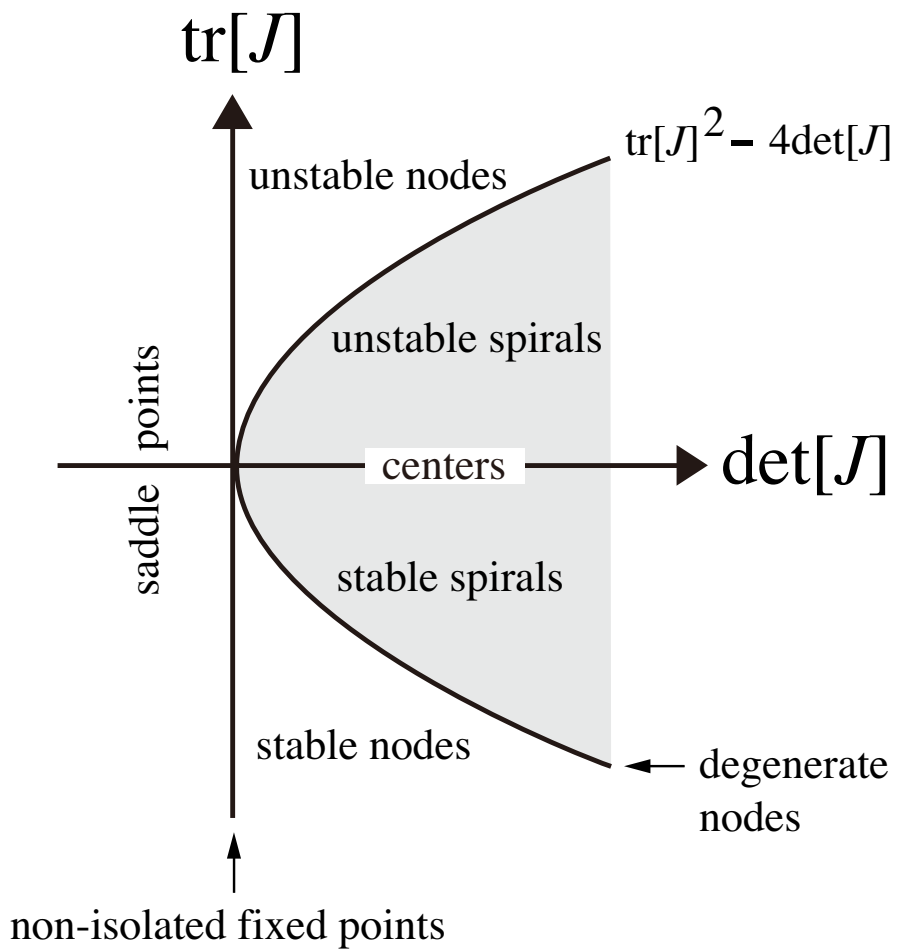
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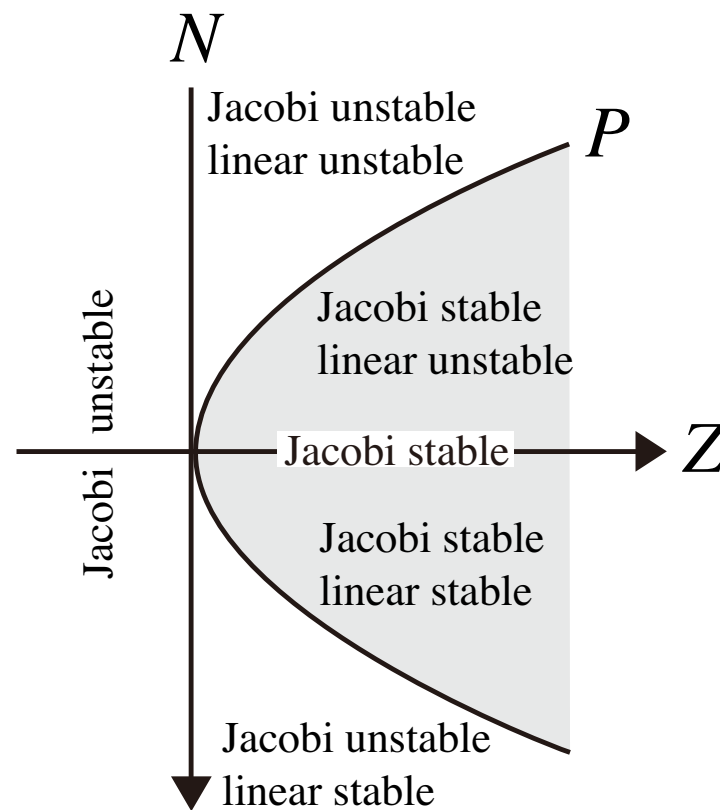
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428 **Figure captions**

429
430 **Fig. 1.** Type and stability of the equilibrium points. (a) The well known
431 diagram expressed in terms of the Jacobian J . (b). The corresponding
432 diagram expressed in terms of the geometrical quantities of the KCC theory.
433 Note that the N -axis is reversed because N is given by negative $Tr[J]$.
434



(a)



(b)