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Quantum Representation of Affine Weyl Groups and Associated Quantum Curves

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Abstract. We study a quantum (non-commutative) representation of the affine Weyl group mainly of type $E_8^{(1)}$, where the representation is given by birational actions on two variables x, y with q -commutation relations. Using the tau variables, we also construct quantum “fundamental” polynomials $F(x, y)$ which completely control the Weyl group actions. The geometric properties of the polynomials $F(x, y)$ for the commutative case is lifted distinctively in the quantum case to certain singularity structures as the q -difference operators. This property is further utilized as the characterization of the quantum polynomials $F(x, y)$. As an application, the quantum curve associated with topological strings proposed recently by the first named author is rederived by the Weyl group symmetry. The cases of type $D_5^{(1)}$, $E_6^{(1)}$, $E_7^{(1)}$ are also discussed.

Key words: affine Weyl group; quantum curve; Painlevé equation

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1 Introduction

Quantization of the Painlevé equations (or isomonodromic deformations more generally) and their discrete variations is an important problem. Recently, this subject attracts various interests due to its relation to conformal field theories, gauge theories and topological strings. Despite some interesting pioneering works [5, 6, 7, 9, 19, 37], there remain many problems to be studied especially on the quantization of the discrete Painlevé equations. One of the main problems is to establish the quantization compatible with the geometric formulation in [29, 49].¹ Such a study is expected to clarify various developments mentioned above from a geometric viewpoint of *quantum curves*.

Recently, in the study of topological strings, certain quantum curves related to the affine Weyl group of type $D_5^{(1)}$, $E_6^{(1)}$, $E_7^{(1)}$, $E_8^{(1)}$ were obtained [41]. The quantum curves were obtained by combining previous classical results in [3, 33] and an empirical observation for quantization of the classical multiplicities [36] (as discussed later in Section 3). Our main motivation is to formulate a quantum representation of the affine Weyl groups to provide a solid basis for the study of these quantum curves and the corresponding quantum q -difference Painlevé equations. Among others, our work enables the derivation of these quantum curves from the first principle.² As discussed in the last section, we expect that our work will clarify the group structure of various related physical theories. In particular, we hope to relate directly our tau functions

¹In Appendix B, we give a short summary for the classical cases.

²Recently, the elliptic quantum curve for the E-string theory is obtained in [10].

fully equipped with the group structure to the partition functions in topological strings in the future.

The contents of this paper is as follows. In the remaining part of this section, we recall some basic results on the representation of the affine Weyl group $W(E_8^{(1)})$ in the commutative case, focusing on polynomials (which we call fundamental or F -polynomials) generated by the Weyl group actions. In Section 2, a natural quantization of the representation of $W(E_8^{(1)})$ is formulated. The quantization of the F -polynomials is associated to q -difference operators and we study a crucial non-logarithmic property of it in Section 3. In Section 4, we show the main theorem which characterizes the quantum F -polynomials. In Section 5, applying the constructions, we give a characterization of the quantum curve of type E_8 . In Section 6, we give a bilinear form of the Weyl group actions. Section 7 is for summary and discussions. In Appendix A, the similar constructions are obtained for the cases of $D_5^{(1)}$, $E_6^{(1)}$ and $E_7^{(1)}$. In Appendix B, the relation of the classical Weyl group representation in Section 1 to the standard representations used in the q -Painlevé equations is summarized.

In order to explain the problem of this paper more explicitly, we recapitulate some basic facts on a birational representation of the affine Weyl group of type $E_8^{(1)}$, $W(E_8^{(1)}) = \langle s_0, s_1, \dots, s_8 \rangle$ defined by the Dynkin diagram:

$$\begin{array}{c} s_0 \\ | \\ s_1 - s_2 - s_3 - s_4 - s_5 - s_6 - s_7 - s_8. \end{array}$$

All the results in this section are known in literature (see [51] for example) up to a change of parametrization, hence we omit the proofs.

Proposition 1.1. *Define the algebra automorphism s_0, \dots, s_8 on parameters $h_1, h_2, e_1, \dots, e_{11}$ and variables $x, y, \sigma_1, \sigma_2, \tau_1, \dots, \tau_{11}$ as*

$$\begin{aligned} s_0 &= \left\{ e_{10} \rightarrow \frac{h_2}{e_{11}}, e_{11} \rightarrow \frac{h_2}{e_{10}}, h_1 \rightarrow \frac{h_1 h_2}{e_{10} e_{11}}, x \rightarrow x \frac{1 + y \frac{h_2}{e_{10}}}{1 + y e_{11}}, \right. \\ &\quad \left. \tau_{10} \rightarrow (1 + y e_{11}) \frac{\sigma_2}{\tau_{11}}, \tau_{11} \rightarrow \frac{\sigma_2}{\tau_{10}} \left(1 + y \frac{h_2}{e_{10}} \right), \sigma_1 \rightarrow (1 + y e_{11}) \frac{\sigma_1 \sigma_2}{\tau_{10} \tau_{11}} \right\}, \\ s_1 &= \{e_8 \leftrightarrow e_9, \tau_8 \leftrightarrow \tau_9\}, \quad s_2 = \{e_7 \leftrightarrow e_8, \tau_7 \leftrightarrow \tau_8\}, \\ s_3 &= \left\{ e_1 \rightarrow \frac{h_1}{e_7}, e_7 \rightarrow \frac{h_1}{e_1}, h_2 \rightarrow \frac{h_1 h_2}{e_1 e_7}, y \rightarrow \frac{1 + x \frac{e_7}{h_1}}{1 + \frac{x}{e_1}} y, \right. \\ &\quad \left. \tau_1 \rightarrow \left(1 + x \frac{e_7}{h_1} \right) \frac{\sigma_1}{\tau_7}, \tau_7 \rightarrow \frac{\sigma_1}{\tau_1} \left(1 + \frac{x}{e_1} \right), \sigma_2 \rightarrow \frac{\sigma_1 \sigma_2}{\tau_1 \tau_7} \left(1 + \frac{x}{e_1} \right) \right\}, \\ s_4 &= \{e_1 \leftrightarrow e_2, \tau_1 \leftrightarrow \tau_2\}, \quad s_5 = \{e_2 \leftrightarrow e_3, \tau_2 \leftrightarrow \tau_3\}, \quad s_6 = \{e_3 \leftrightarrow e_4, \tau_3 \leftrightarrow \tau_4\}, \\ s_7 &= \{e_4 \leftrightarrow e_5, \tau_4 \leftrightarrow \tau_5\}, \quad s_8 = \{e_5 \leftrightarrow e_6, \tau_5 \leftrightarrow \tau_6\}. \end{aligned} \tag{1.1}$$

Then these actions give a birational representation of the affine Weyl group $W(E_8^{(1)})$ on the field of rational functions $\mathbb{C}(h_i, e_i, x, y, \sigma_i, \tau_i)$.

The representation is based on a special configuration of 11 points on $\mathbb{P}^1 \times \mathbb{P}^1$ (see Figure 1). For the blow-up X of $\mathbb{P}^1 \times \mathbb{P}^1$ at the 11 points p_i ($i = 1, \dots, 11$), the Picard lattice $P = \text{Pic}(X)$ is generated by $H_1, H_2, E_1, \dots, E_{11}$, with the only non-vanishing intersection pairings being $H_1 \cdot H_2 = H_2 \cdot H_1 = 1$, $E_i \cdot E_i = -1$. The actions (1.1) are closed on subfields $\mathbb{C}(h_i, e_i)$ and $\mathbb{C}(h_i, e_i, x, y)$. The restriction on $\mathbb{C}(h_i, e_i)$

$$s_0 = \left\{ e_{10} \rightarrow \frac{h_2}{e_{11}}, e_{11} \rightarrow \frac{h_2}{e_{10}}, h_1 \rightarrow \frac{h_1 h_2}{e_{10} e_{11}} \right\}, \quad s_1 = \{e_8 \leftrightarrow e_9\}, \quad s_2 = \{e_7 \leftrightarrow e_8\},$$

$$s_3 = \left\{ e_1 \rightarrow \frac{h_1}{e_7}, e_7 \rightarrow \frac{h_1}{e_1}, h_2 \rightarrow \frac{h_1 h_2}{e_1 e_7} \right\}, \quad s_4 = \{e_1 \leftrightarrow e_2\}, \quad s_5 = \{e_2 \leftrightarrow e_3\},$$

$$s_6 = \{e_3 \leftrightarrow e_4\}, \quad s_7 = \{e_4 \leftrightarrow e_5\}, \quad s_8 = \{e_5 \leftrightarrow e_6\},$$

is nothing but the natural linear actions on the Picard lattice written in the multiplicative notation: $h_i = \exp H_i$, $e_i = \exp E_i$. When $x = y = 0$, the actions on σ_i , τ_i are just copies of the actions on h_i , e_i . In terms of the parameters h_i , e_i the points p_1, \dots, p_{11} can be parametrized as

$$p_i = (-e_i, 0) \quad (i = 1, \dots, 6), \quad p_i = \left(-\frac{h_1}{e_i}, \infty\right) \quad (i = 7, 8, 9),$$

$$p_{10} = \left(\infty, -\frac{e_{10}}{h_2}\right), \quad p_{11} = \left(0, -\frac{1}{e_{11}}\right).$$

This parametrization is compatible under the actions of the Weyl group $W(E_8^{(1)})$.

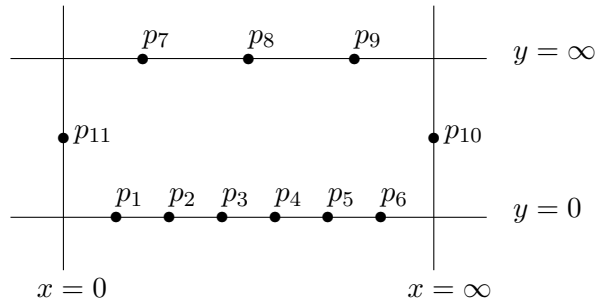


Figure 1. Configuration of the 11 points.

For an algebraic curve in X , its homological data $\lambda = (d_i, m_i)$ (i.e., the bidegree (d_1, d_2) and the multiplicity m_i at the i -th point p_i) can be represented by an element of P as

$$\lambda = d_1 H_1 + d_2 H_2 - m_1 E_1 - \dots - m_{11} E_{11}. \quad (1.2)$$

Sometimes, to represent the data $\lambda = (d_i, m_i)$, we use a multiplicative notation

$$e^\lambda = \frac{h_1^{d_1} h_2^{d_2}}{e_1^{m_1} \dots e_{11}^{m_{11}}}, \quad \tau^\lambda = \frac{\sigma_1^{d_1} \sigma_2^{d_2}}{\tau_1^{m_1} \dots \tau_{11}^{m_{11}}}.$$

We call variables σ_i , τ_i the *tau* variables (or *tau* functions). The tau functions are the main objects in the theory of isomonodromic deformations [26], and their representation-theoretical formulation was initiated in [46]. Although quantum curves as well as their classical analogs can be discussed in the subfield $\mathbb{C}(h_i, e_i, x, y)$ as in [41], we stress that the appropriate introduction of the variables σ_i , τ_i in equation (1.1) clarifies the structure of the Weyl group actions largely since it reduces the problem of rational functions of x, y into that of polynomials (see the last remark in this section). Indeed, the basic fact on the representation (1.1) is the following holomorphic property which is related to the *singularity confinement* (see [17] and references therein) and the *Laurent phenomenon* [14].

Proposition 1.2. *For any $w \in W(E_8^{(1)})$, the action of w on variables τ_i ($i = 1, \dots, 11$) is given by*

$$w(\tau_i) = \phi_{w,i}(x, y) \tau^\lambda, \quad \tau^\lambda = \frac{\sigma_1^{d_1} \sigma_2^{d_2}}{\tau_1^{m_1} \dots \tau_{11}^{m_{11}}},$$

where $\lambda = (d_i, m_i)$ is determined by $w(e_i) = e^\lambda = \frac{h_1^{d_1} h_2^{d_2}}{e_1^{m_1} \dots e_{11}^{m_{11}}}$, and $\phi_{w,i}(x, y)$ is a polynomial associated with the degree/multiplicity data $\lambda = (d_i, m_i)$. Moreover, regardless of the above

construction using the action of w , the polynomial $\phi_{w,i}(x, y)$ can be recovered by the geometric conditions specified by the data $\lambda = (d_i, m_i)$ uniquely up to a normalization. Hence we can denote $\phi_{w,i}(x, y)$ by $F_\lambda(x, y)$.

Remark 1.3. The curves $C: F_\lambda(x, y) = 0$ are transforms of the exceptional curve E_i under the birational actions $w \in W(E_8^{(1)})$, hence the curves C are rational and rigid.

Example 1.4. For $w = s_{3,2,1,0,2,4,3}$ ($= s_3 s_2 s_1 s_0 s_2 s_4 s_3$), we have $e^\lambda := w(e_1) = \frac{h_1^2 h_2}{e_1 e_7 e_9 e_{10} e_{11}}$, and

$$F_\lambda(x, y) = \left(1 + \frac{e_1 e_7 e_9 e_{10} e_{11}}{h_1^2 h_2} x\right) \left(1 + \frac{1}{e_1} x\right) + e_{11} \left(1 + \frac{e_7}{h_1} x\right) \left(1 + \frac{e_9}{h_1} x\right) y. \quad (1.3)$$

For $w = s_{0,3,4,0,2,3,2,1,0,2,4,3}$, we have $e^\lambda := w(e_{11}) = \frac{h_1^2 h_2^2}{e_1 e_2 e_7 e_8 e_{10}^2 e_{11}}$, and

$$\begin{aligned} F_\lambda(x, y) &= \frac{x^2 \left(1 + \frac{h_2}{e_{10}} y\right)^2}{e_1 e_2} + x \left(1 + \frac{h_2}{e_{10}} y\right) \left\{ \left(\frac{1}{e_7} + \frac{1}{e_8}\right) \frac{h_1 h_2}{e_1 e_2 e_{10}} y + \left(\frac{1}{e_1} + \frac{1}{e_2}\right) \right\} \\ &\quad + (1 + e_{11} y) \left(1 + \frac{h_1^2 h_2^2}{e_1 e_2 e_7 e_8 e_{10}^2 e_{11}} y\right). \end{aligned} \quad (1.4)$$

Remark 1.5. We see that the variables k_1, k_2 defined by

$$k_1 = x \frac{\tau_{10}}{\tau_{11}}, \quad k_2 = y \frac{\tau_7 \tau_8 \tau_9}{\tau_1 \tau_2 \cdots \tau_6}, \quad (1.5)$$

are $W(E_8^{(1)})$ invariant. Hence, the rational actions of $w \in W(E_8^{(1)})$ on x, y can be determined by the polynomials corresponding to $w(\tau_1), \dots, w(\tau_{11})$.

2 Quantum representation

In the following, we use the same symbols $h_i, e_i, x, y, \sigma_i, \tau_i$ for the quantum (non-commutative) objects. This notation is economical and consistent with the commutative case since the latter can be recovered by taking the specialization $q = 1$.

Definition 2.1. Let \mathcal{K} be a skew (non-commutative) field on the variables $h_1, h_2, e_1, \dots, e_{11}, x, y, \sigma_1, \sigma_2, \tau_1, \dots, \tau_{11}$, where the non-trivial commutation relations are

$$yx = qxy, \quad \tau_i e_i = q^{-1} e_i \tau_i, \quad \sigma_1 h_2 = q h_2 \sigma_1, \quad \sigma_2 h_1 = q h_1 \sigma_2, \quad (2.1)$$

and other pairs are assumed to be commutative.

Remark 2.2. In view of the results in [38], where the construction of [47] is nicely quantized, it is natural to regard the variables σ_i, τ_i to be dual to the parameters h_i, e_i . Indeed, the q -commutation relations (2.1) can be concisely written as

$$\tau^\lambda e^\mu = q^{\lambda \cdot \mu} e^\mu \tau^\lambda,$$

using the intersection pairing $\lambda \cdot \mu = d_1 d'_2 + d_2 d'_1 - \sum_{i=1}^{11} m_i m'_i$ for $\tau^\lambda = \sigma_1^{d_1} \sigma_2^{d_2} / (\tau_1^{m_1} \cdots \tau_{11}^{m_{11}})$, $e^\mu = h_1^{d'_1} h_2^{d'_2} / (e_1^{m'_1} \cdots e_{11}^{m'_{11}})$ as well as $\lambda = (d_i, m_i), \mu = (d'_i, m'_i)$.

Under the non-commutative setting given above, there exists a natural quantization of Proposition 1.1.

Theorem 2.3. *On the skew field \mathcal{K} there exists a birational representation of the affine Weyl group $W(E_8^{(1)}) = \langle s_0, \dots, s_8 \rangle$ given exactly by the same equation as in equation (1.1).*

Proof. A direct computation (see also Remark 2.8). ■

Remark 2.4. We have fixed the operator ordering in equation (1.1) through the requirements of the Weyl group relations. Since the results seem to be consistent with the prescription of the “ q -ordering” (or Weyl ordering) applied in [41], it will be interesting to study whether and how such a prescription works in general.

Remark 2.5. The quantum Weyl group actions on the subfield $\mathbb{C}(h_i, e_i, x, y)$ can be constructed from the quantum curves in [41] without difficulty. In [41], two realizations of the quantum curves, i.e., the “triangular” form and the “rectangular” form were constructed from a heuristic method by consulting previous classical results in [3, 33] and an empirical quantization rule in [36]. The two realizations are related explicitly by a birational transformation, where each simple reflection s_i is given by explicit actions on $\{h_i, e_i\}$, and besides, trivially on $\{x, y\}$ at least in one realization. By composition, the nontrivial actions in one realization are transplanted from the trivial ones in the other and all the actions of s_i in the subfield $\mathbb{C}(h_i, e_i, x, y)$ are obtained. As a result, the actions of s_i are identical to those anticipated from previous works by [19] for $W(D_5^{(1)})$ and [36]³ for $W(D_5^{(1)})$, $W(E_7^{(1)})$. We emphasize that here the quantum Weyl group actions on the tau variables are also obtained. Namely, inspired by the work [38], we have further noticed that the representations can be lifted by including the variables $\{\sigma_i, \tau_i\}$ as in equation (1.1). Since the final result is quite simple and almost identical to the known classical case, we decide to take a quick style of presentation omitting the roundabout derivations. With the quantum Weyl group actions on the tau variables identified, we can rederive the quantum curves from solid arguments.

Example 2.6. For $w = s_{3,2,1,0,2,4,3}$, $e^\lambda := w(e_1) = \frac{h_1^2 h_2}{e_1 e_7 e_9 e_{10} e_{11}}$, we have

$$F_\lambda(x, y) = \left(1 + \frac{e_1 e_7 e_9 e_{10} e_{11}}{h_1^2 h_2} x\right) \left(1 + \frac{q^{-1}}{e_1} x\right) + e_{11} \left(1 + \frac{e_7}{h_1} x\right) \left(1 + \frac{e_9}{h_1} x\right) y. \quad (2.2)$$

For $w = s_{0,3,4,0,2,3,2,1,0,2,4,3}$, $e^\lambda := w(e_{11}) = \frac{h_1^2 h_2^2}{e_1 e_2 e_7 e_8 e_{10}^2 e_{11}}$, we have

$$\begin{aligned} F_\lambda(x, y) = & \frac{x^2 \left(1 + \frac{h_2}{e_{10}} y\right) \left(1 + \frac{q h_2}{e_{10}} y\right)}{e_1 e_2 q^2} + \frac{x}{q} \left(1 + \frac{h_2}{e_{10}} y\right) \left\{ \left(\frac{1}{e_7} + \frac{1}{e_8}\right) \frac{h_1 h_2}{e_1 e_2 e_{10}} y + \left(\frac{1}{e_1} + \frac{1}{e_2}\right) \right\} \\ & + (1 + e_{11} y) \left(1 + \frac{h_1^2 h_2^2}{q e_1 e_2 e_7 e_8 e_{10}^2 e_{11}} y\right). \end{aligned} \quad (2.3)$$

As expected, equations (2.2) and (2.3) reduce to equations (1.3) and (1.4) respectively when $q = 1$.

The representation can be realized as the adjoint actions as follows.

Theorem 2.7. *The actions s_i on variables $X = e_i, h_i, \tau_i, \sigma_i, x, y$ can be written as*

$$\begin{aligned} s_i(X) &= G_i^{-1} r_i(X) G_i, \\ G_0 &= \frac{\left(\frac{h_2}{e_{10}} y; q\right)_\infty^+}{\left(e_{11} y; q\right)_\infty^+}, \quad G_3 = \frac{\left(\frac{1}{e_1} x; q\right)_\infty^+}{\left(\frac{e_7}{h_1} x; q\right)_\infty^+}, \quad G_i = 1 \quad (i \neq 0, 3), \end{aligned}$$

³Note that it is necessary to generalize slightly from [36, 41] to obtain the representation of the affine Weyl group by lifting the constraint on the parameters, since only symmetries of the quantum curve (which is non-affine) were discussed there.

where $(z; q)_\infty^+ = \prod_{i=0}^\infty (1 + q^i z)$ is the q -factorial and r_i is a multiplicatively linear action on $\{h_i, e_i, \sigma_i, \tau_i\}$ defined by $r_i(X) = s_i(X)|_{x=y=0}$, and $r_i(x) = x$, $r_i(y) = y$.

Proof. Put $G = \frac{(\beta y; q)_\infty^+}{(\alpha y; q)_\infty^+}$. By the relation $f(y)x = xf(qy)$ we have⁴

$$G^{-1}r_0(x)G = G^{-1}xG = \frac{(\alpha y)_\infty^+}{(\beta y)_\infty^+} x \frac{(\beta y)_\infty^+}{(\alpha y)_\infty^+} = x \frac{(\alpha qy)_\infty^+}{(\beta qy)_\infty^+} \frac{(\beta y)_\infty^+}{(\alpha y)_\infty^+} = x \frac{1 + \beta y}{1 + \alpha y}.$$

This gives the action $s_0(x)$ when $\alpha = e_{11}$, $\beta = \frac{h_2}{e_{10}}$, i.e., $G = G_0$. Fortunately, the formula $G_0^{-1}r_0(*)G_0$ recovers the correct transformation for the other variables as well. For instance

$$G_0^{-1}r_0(\tau_{10})G_0 = G_0^{-1}\frac{\sigma_2}{\tau_{11}}G_0 = G_0^{-1}\left(G_0\Big|_{\substack{h_1 \rightarrow qh_1, \\ e_{11} \rightarrow qe_{11}}}\right)\frac{\sigma_2}{\tau_{11}} = (1 + e_{11}y)\frac{\sigma_2}{\tau_{11}}.$$

The case $i = 3$ is similar and the other cases are obvious. ■

Remark 2.8. Using the realization s_i in Theorem 2.7, one can give another proof of the Weyl group relations as follows. We consider the most non-trivial case $s_0s_3s_0 = s_3s_0s_3$ as an example. Since

$$\begin{aligned} s_0(X) &= G_0^{-1}r_0(X)G_0, \\ s_3s_0(X) &= G_3^{-1}(r_3G_0^{-1})(r_3r_0X)(r_3G_0)G_3, \\ s_0s_3s_0(X) &= G_0^{-1}(r_0G_3^{-1})(r_0r_3G_0^{-1})(r_0r_3r_0X)(r_0r_3G_0)(r_0G_3)G_0, \end{aligned}$$

we have $s_0s_3s_0(X) = G^{-1}(r_0r_3r_0X)G$, where

$$G = (r_0r_3G_0)(r_0G_3)G_0 = \frac{\left(\frac{h_1h_2}{e_1e_7e_{10}}y\right)_\infty^+}{\left(\frac{h_2}{e_{10}}y\right)_\infty^+} \frac{\left(\frac{1}{e_1}x\right)_\infty^+}{\left(\frac{e_7e_{10}e_{11}}{h_1h_2}x\right)_\infty^+} \frac{\left(\frac{h_2}{e_{10}}y\right)_\infty^+}{(e_{11}y)_\infty^+}.$$

Similarly we have $s_3s_0s_3(X) = \tilde{G}^{-1}(r_3r_0r_3X)\tilde{G}$, where

$$\tilde{G} = (r_3r_0G_3)(r_3G_0)G_3 = \frac{\left(\frac{e_7}{h_1}x\right)_\infty^+}{\left(\frac{e_7e_{10}e_{11}}{h_1h_2}x\right)_\infty^+} \frac{\left(\frac{h_1h_2}{e_1e_7e_{10}}y\right)_\infty^+}{(e_{11}y)_\infty^+} \frac{\left(\frac{1}{e_1}x\right)_\infty^+}{\left(\frac{e_7}{h_1}x\right)_\infty^+}.$$

Due to the relation $r_0r_3r_0 = r_3r_0r_3$, the relation $s_0s_3s_0 = s_3s_0s_3$ is guaranteed if $G = \tilde{G}$. Rescaling $y \rightarrow \frac{e_{10}}{h_2}y$, $x \rightarrow \frac{h_1}{e_7}x$ and putting $a = \frac{h_1}{e_1e_7}$, $b = \frac{e_{10}e_{11}}{h_2}$, the relation $G = \tilde{G}$ reduces to the following identity which may be considered as a version of the quantum dilogarithm identity (see, e.g., [35] and references therein).

Lemma 2.9. *For non-commuting variables $yx = qxy$, we have*

$$\frac{(ay)_\infty^+ (ax)_\infty^+ (y)_\infty^+}{(y)_\infty^+ (bx)_\infty^+ (by)_\infty^+} = \frac{(x)_\infty^+ (ay)_\infty^+ (ax)_\infty^+}{(bx)_\infty^+ (by)_\infty^+ (x)_\infty^+}. \quad (2.4)$$

Proof. By replacements $x \rightarrow -x$ and $y \rightarrow -y$, equation (2.4) can be written as

$$\frac{(ay)_\infty (ax)_\infty (y)_\infty}{(y)_\infty (bx)_\infty (by)_\infty} = \frac{(x)_\infty (ay)_\infty (ax)_\infty}{(bx)_\infty (by)_\infty (x)_\infty}, \quad (2.5)$$

⁴We sometimes omit the base q as $(z)_\infty^+ = (z; q)_\infty^+$. Note that our definition of the q -factorial is different from the conventional one $(z; q)_\infty = \prod_{i=0}^\infty (1 - q^i z)$ by signs, which also appears later.

where $(x)_\infty = \prod_{i=0}^{\infty} (1 - q^i x)$, and we will prove equation (2.4) in this form. We recall the q -binomial identity

$$\frac{(az)_\infty}{(z)_\infty} = \sum_{n \geq 0} \frac{(a)_n}{(q)_n} z^n, \quad (a)_n = \frac{(a)_\infty}{(aq^n)_\infty}, \quad (2.6)$$

which follows by solving the difference equation $f(qz) = \frac{1-z}{1-az} f(z)$ for $f(z) = \frac{(az)_\infty}{(z)_\infty}$ in series expansion. Using equation (2.6) and $yx = qxy$, the factors in equation (2.5) can be reordered as

$$\begin{aligned} \frac{(ay)_\infty}{(y)_\infty} \frac{(ax)_\infty}{(bx)_\infty} &= \sum_{n \geq 0} \frac{(a)_n}{(q)_n} y^n \frac{(ax)_\infty}{(bx)_\infty} = \sum_{n \geq 0} \frac{(a)_n}{(q)_n} \frac{(aq^n x)_\infty}{(bq^n x)_\infty} y^n = \frac{(ax)_\infty}{(bx)_\infty} \sum_{n \geq 0} \frac{(a)_n}{(q)_n} \frac{(bx)_n}{(ax)_n} y^n, \\ \frac{(ay)_\infty}{(by)_\infty} \frac{(ax)_\infty}{(x)_\infty} &= \frac{(ay)_\infty}{(by)_\infty} \sum_{n \geq 0} \frac{(a)_n}{(q)_n} x^n = \sum_{n \geq 0} x^n \frac{(a)_n}{(q)_n} \frac{(aq^n y)_\infty}{(bq^n y)_\infty} = \sum_{n \geq 0} x^n \frac{(a)_n}{(q)_n} \frac{(by)_n}{(ay)_n} \frac{(ay)_\infty}{(by)_\infty}. \end{aligned}$$

Hence, equation (2.5) can be written as

$$(ax)_\infty \sum_{n \geq 0} \frac{(a)_n}{(q)_n} \frac{(bx)_n}{(ax)_n} y^n (y)_\infty = (x)_\infty \sum_{n \geq 0} x^n \frac{(a)_n}{(q)_n} \frac{(by)_n}{(ay)_n} (ay)_\infty. \quad (2.7)$$

Since the both hand sides of equation (2.7) are written in the same ordering in x, y , whether the equality holds or not is independent of the commutation relation of x, y . We will show it in the commutative case, where equation (2.7) can be written as

$$(ax)_\infty {}_2\varphi_1 \left(\begin{matrix} a, bx \\ ax \end{matrix} ; y \right) (y)_\infty = (x)_\infty {}_2\varphi_1 \left(\begin{matrix} a, by \\ ay \end{matrix} ; x \right) (ay)_\infty, \quad (2.8)$$

using the Heine's q -hypergeometric series

$${}_2\varphi_1 \left(\begin{matrix} a, b \\ c \end{matrix} ; x \right) = \sum_{n \geq 0} \frac{(a)_n (b)_n}{(q)_n (c)_n} x^n.$$

Then equation (2.8) can be confirmed via iterative use of the Heine's identity and the trivial symmetry relation

$${}_2\varphi_1 \left(\begin{matrix} a, b \\ c \end{matrix} ; x \right) = \frac{(ax)_\infty}{(x)_\infty} \frac{(b)_\infty}{(c)_\infty} {}_2\varphi_1 \left(\begin{matrix} c/b, x \\ ax \end{matrix} ; b \right), \quad {}_2\varphi_1 \left(\begin{matrix} a, b \\ c \end{matrix} ; x \right) = {}_2\varphi_1 \left(\begin{matrix} b, a \\ c \end{matrix} ; x \right).$$

The former is also obtained from the q -binomial identity. ■

Proposition 2.10. *We put k_1, k_2 as the same as the classical case (1.5),*

$$k_1 = x \frac{\tau_{10}}{\tau_{11}}, \quad k_2 = y \frac{\tau_7 \tau_8 \tau_9}{\tau_1 \tau_2 \cdots \tau_6}.$$

Then k_1, k_2 are $W(E_8^{(1)})$ invariant also in the quantum setting.

Proof. We will check only the nontrivial actions and they go as

$$s_0 \left(x \frac{\tau_{10}}{\tau_{11}} \right) = x \frac{1 + y \frac{h_2}{e_{10}}}{1 + ye_{11}} (1 + ye_{11}) \frac{\sigma_2}{\tau_{11}} \frac{1}{1 + y \frac{h_2}{e_{10}}} \frac{\tau_{10}}{\sigma_2} = x \frac{\tau_{10}}{\tau_{11}},$$

and

$$s_3 \left(\frac{\tau_7}{\tau_1} y \frac{\tau_8 \tau_9}{\tau_2 \cdots \tau_6} \right) = \frac{\tau_7}{\sigma_1} \frac{1}{1 + x \frac{e_7}{h_1}} \frac{\sigma_1}{\tau_7} \left(1 + \frac{x}{e_1} \right) \frac{1 + x \frac{e_7}{h_1}}{1 + \frac{x}{e_1}} y \frac{\tau_8 \tau_9}{\tau_2 \cdots \tau_6} = \frac{\tau_7}{\tau_1} y \frac{\tau_8 \tau_9}{\tau_2 \cdots \tau_6}. \quad \blacksquare$$

Due to this proposition, the actions of $w \in W(E_8^{(1)})$ on x, y can be reduced to the actions on σ_i, τ_i as in the classical case.

3 Non-logarithmic property

From the several examples of the quantum polynomials as in equations (2.2) and (2.3), one observes an interesting factorization in their coefficients, which was utilized in constructing quantum curves in [41]. We will clarify the meaning of such factorizations from the viewpoint of the q -difference operators.

Consider a q -difference equation $D\psi(x) = 0$, $D = \sum_{i=0}^{d_1} x^i A_i(y)$, ($yx = qxy$). We look for a solution $\psi(x)$ around $x = 0$ of the form

$$\psi(x) = x^\rho \sum_{j=0}^{\infty} c_j x^j \quad (c_0 \neq 0).$$

From the coefficient of $x^{\rho+k}$ in the equation $D\psi(x) = 0$, we have

$$\sum_{i+j=k} A_i(q^{\rho+j})c_j = A_k(q^\rho)c_0 + A_{k-1}(q^{\rho+1})c_1 + \cdots + A_0(q^{\rho+k})c_k = 0,$$

where $A_i(y) = 0$ for $i > d_1$. The (multiplicative) exponents $y = q^\rho$ are determined as the zeros of $A_0(y)$. Then the coefficients c_1, c_2, \dots will be determined recursively. For c_k , we have the following cases:

- (1) If $A_0(q^{\rho+k}) \neq 0$, then c_k is uniquely determined from c_0, c_1, \dots, c_{k-1} .
- (2a) If $A_0(q^{\rho+k}) = 0$ and $X_k := A_k(q^\rho)c_0 + A_{k-1}(q^{\rho+1})c_1 + \cdots + A_1(q^{\rho+k-1})c_{k-1} \neq 0$, then the equation for c_k has no solution and we do not have the power series solution (one should consider a solution with logarithmic terms in x).
- (2b) If $A_0(q^{\rho+k}) = 0$ and $X_k = 0$, then the coefficient c_k is free and we still have series solutions with exponents $y = q^\rho, q^{\rho+k}$.

For the last case (2b), the difference operator D admits a non-logarithmic solution around $x = 0$ and $x = 0$ is called “non-logarithmic” singularity of D . Non-logarithmic singularities around $x = \infty$ (or $y = 0$ or $y = \infty$) are defined similarly. If we apply the condition of non-logarithmic singularities to the case with successive exponents, coefficients of the q -difference operator D are constrained strongly by the non-logarithmic properties of its solution as follows.

Proposition 3.1. *For a difference operator $D = \sum_{i=0}^{d_1} x^i A_i(y)$, we have*

- (1) D has non-logarithmic singularities at $x = 0$ with $y = a, qa, \dots, q^{m-1}a \Leftrightarrow A_i(y) \propto \prod_{j=0}^{m-i-1} (y - q^j a)$ for $0 \leq i \leq m-1$,
- (2) D has non-logarithmic singularities at $x = \infty$ with $y = a, q^{-1}a, \dots, q^{-m+1}a \Leftrightarrow A_i(y) \propto \prod_{j=0}^{m-i-1} (y - q^{-j} a)$ for $d_1 - m + 1 \leq i \leq d_1$.

Similarly, for a difference operator $D = \sum_{i=0}^{d_2} B_i(x)y^i$, we have

- (3) D has non-logarithmic singularities at $y = 0$ with $x = a, qa, \dots, q^{m-1}a \Leftrightarrow B_i(x) \propto \prod_{j=0}^{m-i-1} (x - q^j a)$ for $1 \leq i \leq m$,
- (4) D has non-logarithmic singularities at $y = \infty$ with $x = a, q^{-1}a, \dots, q^{-m+1}a \Leftrightarrow B_i(x) \propto \prod_{j=0}^{m-i-1} (x - q^{-j} a)$ for $d_2 - m + 1 \leq i \leq d_2$.

Proof. Consider the case (1) (the other cases are similar). For the non-logarithmic property with successive exponents, the recursion relations for the power series solution

$$\begin{aligned} A_0(y)c_0 &= 0, \\ A_1(y)c_0 + A_0(qy)c_1 &= 0, \\ &\dots\dots\dots \\ A_{m-1}(y)c_0 + \dots + A_0(q^{m-1}y)c_{m-1} &= 0, \end{aligned}$$

should be satisfied termwise with m free coefficients: c_0, \dots, c_{m-1} . From the first relation we have $A_0(y) \propto \prod_{j=0}^{m-1} (y - q^j a)$, and the other factorizations also follow easily. ■

In other words, a q -difference operator $D = \sum_{i=0}^{d_1} x^i A_i(y)$ with boundary coefficients $A_0(y)$, $A_{d_2}(y)$ having zeros successive in powers of q , is non-logarithmic iff suitable parts of the zeros penetrate into the internal coefficients. We have similar properties for a difference operator $D = \sum_{i=0}^{d_2} B_i(x)y^i$ also. The non-logarithmic property of q -difference operators plays important roles in the following characterization of quantum polynomials and also in [43, 45, 50, 52] etc.

4 The F -polynomials

Here we study the quantum analog of the polynomials $F_\lambda(x, y)$ in Proposition 1.2.

Definition 4.1. For each degree/multiplicity data $\lambda = ((d_1, d_2), (m_1, \dots, m_{11})) \in P$, we define a non-commutative polynomial $F = F_\lambda(x, y) = F_\lambda(x, y; \{h_i, e_i\})$ by the following conditions:

$(x)_\lambda$ Collecting terms with the same power of x , the polynomial F takes the form

$$F = \sum_{i=0}^{d_1} x^i \prod_{t=i}^{m_{11}-1} (1 + q^t e_{11} y) \prod_{t=d_1-m_{10}}^{i-1} \left(1 + q^t \frac{h_2}{e_{10}} y\right) U_i(y),$$

where $U_i(y)$ is a polynomial⁵ in y of degree $d_2 - (i - d_1 + m_{10})_+ - (m_{11} - i)_+$.

$(y)_\lambda$ Collecting terms with the same power of y , the polynomial F takes the form

$$F = \sum_{i=0}^{d_2} \prod_{k=1}^6 \prod_{t=i-m_k}^{-1} \left(1 + q^t \frac{1}{e_k} x\right) \prod_{k=7}^9 \prod_{t=0}^{i-d_2+m_k-1} \left(1 + q^t \frac{e_k}{h_1} x\right) V_i(x) y^i, \quad (4.1)$$

where $V_i(x)$ is a polynomial in x of degree $d_1 - \sum_{k=1}^6 (m_k - i)_+ - \sum_{k=7}^9 (i - d_2 + m_k)_+$.

In these conditions, $(x)_+ = \max(x, 0)$ and the empty product is 1: $\prod_{t=a}^b (*) = 1$ ($a > b$).

Remark 4.2. For the $q = 1$ case, it is easy to see that the conditions $(x)_\lambda, (y)_\lambda$ reduce to the conditions specified by the degree/multiplicity data $\lambda = (d_i, m_i)$. Hence the quantum polynomial $F_\lambda(x, y)$ reduces to the classical polynomial $F_\lambda(x, y)$ in Proposition 1.2.

Proposition 4.3. Let Λ be the $W(E_8^{(1)})$ -orbit of $\{E_1, \dots, E_{11}\}$. Then for $\lambda \in \Lambda$, the polynomial $F_\lambda(x, y)$ exists and is unique up to a normalization. We will normalize it by $F_\lambda(0, 0) = 1$.

⁵If there appear many polynomials $U_i(y)$ of the same degree, they should be considered as different ones. This applies to $V_i(x)$ in equation (4.1) as well.

Proof. The conditions $(x)_\lambda, (y)_\lambda$ give linear equations⁶ (vanishing conditions) for $F_\lambda(x, y)$. Counting the numbers of coefficients and equations, the dimension of the solution is given by

$$\dim = (d_1 + 1)(d_2 + 1) - \sum_{k=1}^{11} \frac{m_k(m_k + 1)}{2} = \frac{1}{2}\lambda \cdot \lambda + \frac{1}{2}\lambda \cdot \delta_{\text{Red}} + 1, \quad (4.2)$$

where λ is in equation (1.2), $\text{dot}(\cdot)$ is the intersection pairing and $\delta_{\text{Red}} = 2H_1 + 2H_2 - \sum_{i=1}^{11} E_i$. Then, for $\lambda \in \Lambda$ we have $\dim = 1$, since $\lambda \cdot \lambda = -1$ and $\lambda \cdot \delta_{\text{Red}} = 1$. \blacksquare

We use a notation s_i^* to represent the induced action on the data $\lambda = (d_i, m_i)$ defined by $s_i(e^\lambda) = e^{s_i^*\lambda}$, hence $s_i(\tau^\lambda)|_{x=y=0} = \tau^{s_i^*\lambda}$. It is explicitly given as

$$\begin{aligned} s_0^* &= \{d_2 \mapsto d_1 + d_2 - m_{10} - m_{11}, m_{10} \mapsto d_1 - m_{11}, m_{11} \mapsto d_1 - m_{10}\}, \\ s_1^* &= \{m_8 \leftrightarrow m_9\}, \quad s_2^* = \{m_7 \leftrightarrow m_8\}, \\ s_3^* &= \{d_1 \mapsto d_1 + d_2 - m_1 - m_7, m_1 \mapsto d_2 - m_7, m_7 \mapsto d_2 - m_1\}, \quad s_4^* = \{m_1 \leftrightarrow m_2\}, \\ s_5^* &= \{m_2 \leftrightarrow m_3\}, \quad s_6^* = \{m_3 \leftrightarrow m_4\}, \quad s_7^* = \{m_4 \leftrightarrow m_5\}, \quad s_8^* = \{m_5 \leftrightarrow m_6\}. \end{aligned}$$

The following is the main result of this paper.

Theorem 4.4. *Let $F_\lambda(x, y)$ be a polynomial satisfying the conditions $(x)_\lambda, (y)_\lambda$. Then for each simple reflection $s_i \in W(E_8^{(1)})$, the function $F_{s_i^*\lambda}(x, y)$ defined by*

$$s_i(F_\lambda(x, y)\tau^\lambda) = F_{s_i^*\lambda}(x, y)\tau^{s_i^*\lambda}, \quad \tau^\lambda = \frac{\sigma_1^{d_1}\sigma_2^{d_2}}{\tau_1^{m_1}\dots\tau_{11}^{m_{11}}}, \quad (4.3)$$

is also a polynomial in x, y and satisfy the condition $(x)_{s_i^\lambda}, (y)_{s_i^*\lambda}$. In particular, for $\lambda \in \Lambda$, the unique normalized polynomials $F_\lambda(x, y)$ can be obtained by the actions (4.3) from the initial condition $F_{e_i} = 1$.*

Remark 4.5. The polynomial $F_\lambda(x, y)$ is not a function but a section of a line bundle \mathcal{L}_λ on X , and equation (4.3) can be considered as its trivialization in the commutative case [46, 47]. Theorem 4.4 suggests a non-commutative analog of such a geometric understanding.

Example 4.6. For $e^\lambda = \frac{h_1 h_2}{e_{10} e_{11}}$, the corresponding F_λ has two parameters:

$$F_\lambda = c_0(1 + e_{11}y) + c_1x\left(1 + \frac{h_2}{e_{10}}y\right). \quad (4.4)$$

Then, we have

$$s_3\left(F_\lambda \frac{\sigma_1 \sigma_2}{\tau_{10} \tau_{11}}\right) = \tilde{F}_\lambda \frac{\sigma_1^2 \sigma_2}{\tau_1 \tau_7 \tau_{10} \tau_{11}},$$

where

$$\begin{aligned} \tilde{F}_\lambda &= (c_0 + c_1x)\left(1 + \frac{1}{qe_1}x\right) + \left(1 + \frac{e_7}{h_1}x\right)\left(c_0e_{11} + c_1\frac{h_1 h_2}{e_1 e_7 e_{10}}x\right)y \\ &= c_0(1 + e_{11}y) + x\left\{c_0\left(\frac{1}{qe_1} + \frac{e_7 e_{11}}{h_1}y\right) + c_1\left(1 + \frac{h_1 h_2}{e_1 e_7 e_{10}}y\right)\right\} + c_1\frac{1}{qe_1}x^2\left(1 + q\frac{h_2}{e_{10}}y\right). \end{aligned}$$

We see that the polynomial \tilde{F}_λ gives a general solution for the condition $(x)_{\tilde{\lambda}}, (y)_{\tilde{\lambda}}$, where $e^{\tilde{\lambda}} = s_3(e^\lambda) = \frac{h_1^2 h_2}{e_1 e_7 e_{10} e_{11}}$.

⁶In the commutative case, this is known as the linear system $|\lambda|$.

Proof of Theorem 4.4. We will consider the cases s_0 and s_3 (other cases are obvious).

Case s_0 . Let $F = F_\lambda(x, y)$ be a polynomial satisfying the condition $(x)_\lambda$. We compute the action of s_0 on $F\tau^\lambda$. For F , we have

$$\begin{aligned} F &= \sum_{i=0}^{d_1} x^i \prod_{t=i}^{m_{11}-1} (1 + q^t e_{11}y) \prod_{t=d_1-m_{10}}^{i-1} \left(1 + q^t \frac{h_2}{e_{10}}y\right) U_i(y) \\ &\xrightarrow{s_0} \sum_{i=0}^{d_1} x^i \prod_{t=0}^{i-1} \frac{1 + q^t \frac{h_2}{e_{10}}y}{1 + q^t e_{11}y} \prod_{t=i}^{m_{11}-1} \left(1 + q^t \frac{h_2}{e_{10}}y\right) \prod_{t=d_1-m_{10}}^{i-1} (1 + q^t e_{11}y) \tilde{U}_i(y), \end{aligned}$$

where $\tilde{U}_i(y)$ is a polynomial in y of degree i . For τ^λ , considering only the relevant factors, we have

$$\begin{aligned} \frac{\sigma_1^{d_1} \sigma_2^{d_2}}{\tau_{10}^{m_{10}} \tau_{11}^{m_{11}}} &= \tau_{11}^{-m_{11}} \tau_{10}^{d_1-m_{10}} \left(\frac{\sigma_1}{\tau_{10}}\right)^{d_1} \sigma_2^{d_2} \\ &\xrightarrow{s_0} \prod_{t=0}^{m_{11}-1} \frac{1}{1 + q^t \frac{h_2}{e_{10}}y} \prod_{t=0}^{d_1-m_{10}-1} (1 + q^t e_{11}y) \frac{\sigma_1^{d_1} \sigma_2^{d_1+d_2-m_{10}-m_{11}}}{\tau_{10}^{d_1-m_{11}} \tau_{11}^{d_1-m_{10}}}. \end{aligned}$$

Collecting the factors $(1 + q^t e_{11}y)$ and $(1 + q^t \frac{h_2}{e_{10}}y)$, we have $s_0(F\tau^\lambda) = \tilde{F}\tau^{s_0\lambda}$, where

$$\tilde{F} = \sum_{i=0}^{d_1} x^i \prod_{t=i}^{d_1-m_{10}-1} (1 + q^t e_{11}y) \prod_{t=m_{11}}^{i-1} \left(1 + q^t \frac{h_2}{e_{10}}y\right) \tilde{U}_i(y).$$

Note that here we have applied the formula $\prod_{t=u}^{v-1} (*) [\prod_{t=w}^{u-1} (*) / \prod_{t=w}^{v-1} (*)] = \prod_{t=v}^{u-1} (*)$, which holds for $w \leq \min(u, v)$. Hence, \tilde{F} is a polynomial of bidegree $(d_1, \tilde{d}_2 = d_1 + d_2 - m_{10} - m_{11})$ satisfying the condition $(x)_{\tilde{\lambda}}$ for $\tilde{\lambda} = s_0(\lambda)$. Moreover, \tilde{F} satisfies the condition $(y)_{\tilde{\lambda}}$ also. To confirm this, we note that the condition $(y)_\lambda$ is equivalent to the condition on the top and bottom coefficients of $F = \sum_{i=0}^{d_2} A_i(x)y^i$:

$$A_0 = \text{const} \prod_{k=1}^6 \prod_{t=-m_k}^{-1} \left(1 + q^t \frac{1}{e_k}x\right), \quad A_{d_2} = \text{const} \prod_{k=7}^9 \prod_{t=0}^{m_k-1} \left(1 + q^t \frac{e_k}{h_1}x\right),$$

together with the non-logarithmic properties. For the coefficients $\tilde{A}_0, \tilde{A}_{\tilde{d}_2}$ of \tilde{F} , we have obviously $\tilde{A}_0 = s_0(A_0) = \text{const} A_0$, and we also have

$$\tilde{A}_{\tilde{d}_2} = s_0(A_{d_2}) = \text{const} \prod_{k=7}^9 \prod_{t=0}^{m_k-1} \left(1 + q^t \frac{e_k}{s_0(h_1)} \frac{h_2}{e_{10}e_{11}}x\right) = \text{const} A_{d_2},$$

since $s_0(x) = x \frac{1 + \frac{h_2}{e_{10}}y}{1 + e_{11}y} \rightarrow \frac{h_2}{e_{10}e_{11}}x$ ($y \rightarrow \infty$). Hence, the leading coefficients $\tilde{A}_0, \tilde{A}_{\tilde{d}_2}$ have the required from $(y)_{\tilde{\lambda}}$. Our remaining task is to show that the non-logarithmic property of \tilde{F} is inherited from that of F . Indeed, recall that the s_0 -transformation is realized as the adjoint action $s_0(X) = G_0^{-1}r_0(X)G_0$ with $G_0 = (y \frac{h_2}{e_{10}})_\infty^+ / (ye_{11})_\infty^+$. Then, under the corresponding transformation of the solutions $\psi(y) \xrightarrow{s_0} G_0(y)^{-1}r_0(\psi(y))$, the non-logarithmic property around $y = 0$ is preserved from the regularity of the q -factorial $(z)_\infty^+$. Besides, with the rewriting

$$G_0 = \frac{(y \frac{h_2}{e_{10}})_\infty^+}{(ye_{11})_\infty^+} = C(y)y^\nu \frac{(q/(ye_{11}))_\infty^+}{(q/(y \frac{h_2}{e_{10}}))_\infty^+}, \quad q^\nu = \frac{e_{10}e_{11}}{h_2},$$

we can use $\tilde{G}_0 = y^\nu(q/(ye_{11}))_\infty^+/(q/(y\frac{h_2}{e_{10}}))_\infty^+$ instead of G_0 , since the factor $C(y)$ is a pseudo constant: $C(qy) = C(y)$ and irrelevant for the adjoint action. Hence the non-logarithmic property around $y = \infty$ follows similarly.

Case s_3 . Let $F = F_\lambda(x, y)$ be a polynomial satisfying the condition $(y)_\lambda$. The action of s_3 on two parts of $F\tau^\lambda$ is given by

$$\begin{aligned} & \prod_{t=i-m_1}^{-1} \left(1 + q^t \frac{1}{e_1} x\right) \prod_{t=0}^{i-d_2+m_7-1} \left(1 + q^t \frac{e_7}{h_1} x\right) V_i(x) y^i \\ & \xrightarrow{s_3} \prod_{t=i-m_1}^{-1} \left(1 + q^t \frac{e_7}{h_1} x\right) \prod_{t=0}^{i-d_2+m_7-1} \left(1 + q^t \frac{1}{e_1} x\right) \tilde{V}_i(x) \prod_{t=1}^{i-1} \frac{1 + q^t \frac{e_7}{h_1} x}{1 + q^t \frac{1}{e_1} x} y^i, \end{aligned}$$

and

$$\frac{\sigma_1^{d_1} \sigma_2^{d_2}}{\tau_1^{m_1} \tau_7^{m_7}} = \tau_1^{-m_1} \tau_7^{d_2-m_7} \left(\frac{\sigma_2}{\tau_7}\right)^{d_2} \sigma_1^{d_1} \xrightarrow{s_3} y^{-i} \prod_{t=i-m_1}^{i-1} \frac{1}{1 + q^t \frac{e_7}{h_1} x} \prod_{t=i-d_2+m_7}^{i-1} \left(1 + q^t \frac{1}{e_1} x\right) y^i,$$

where we choose $\tilde{\lambda} = s_3^* \lambda$ with the tilde applying to each component of $\lambda = (d_i, m_i)$. By combining them, the factors $(1 + q^t \frac{1}{e_1} x)$ and $(1 + q^t \frac{e_7}{h_1} x)$ in the coefficient of y^i in $s_3(F\tau^\lambda)$ are

$$\prod_{t=i-d_2+m_7}^{-1} \left(1 + q^t \frac{1}{e_1} x\right) \prod_{t=0}^{i-m_1-1} \left(1 + q^t \frac{e_7}{h_1} x\right),$$

where we have applied the formula $\prod_{t=u}^{v-1} (*) [\prod_{t=v}^{w-1} (*) / \prod_{t=u}^{w-1} (*)] = \prod_{t=v}^{u-1} (*)$, which holds for $\max(u, v) \leq w$. Then, we have $s_3(F\tau^\lambda) = \tilde{F}\tau^{\tilde{\lambda}}$ with

$$\tilde{F} = \sum_{i=0}^{d_2} \prod_{k=1}^6 \prod_{t=i-\tilde{m}_k}^{-1} \left(1 + q^t \frac{1}{e_k} x\right) \prod_{k=7}^9 \prod_{t=0}^{\tilde{m}_k - \tilde{d}_2 + i - 1} \left(1 + q^t \frac{e_k}{h_1} x\right) \tilde{V}_i(x) y^i,$$

and hence, \tilde{F} satisfies the condition $(y)_{\tilde{\lambda}}$ as desired. The condition $(x)_{\tilde{\lambda}}$ can be confirmed using the adjoint action realization of s_3 . \blacksquare

5 Quantum E_8 curve as the Weyl group invariant

Consider a degree/multiplicity data

$$\lambda = (d_i, m_i) = ((6, 3), (1, 1, 1, 1, 1, 1, 2, 2, 2, 3, 3)), \quad (5.1)$$

which is special since it is invariant $w(\lambda) = \lambda$ under $w \in W(E_8^{(1)})$, i.e., $\{\lambda\}$ is a Weyl orbit with only a single element, hence λ is not in the Weyl orbit Λ of Proposition 4.3. We look for the corresponding quantum polynomial $P = P(x, y)$ defined by the conditions $(x)_\lambda, (y)_\lambda$:

$$\begin{aligned} P &= y^{[0]} \prod_{t=0}^2 (1 + q^t e_{11} y) + xy^{[1]} \prod_{t=1}^2 (1 + q^t e_{11} y) + x^2 y^{[2]} (1 + q^2 e_{11} y) + x^3 y^{[3]} \\ &\quad + x^4 y^{[2]} \left(1 + q^3 \frac{h_2}{e_{10}} y\right) + x^5 y^{[1]} \prod_{t=3}^4 \left(1 + q^t \frac{h_2}{e_{10}} y\right) + x^6 y^{[0]} \prod_{t=3}^5 \left(1 + q^t \frac{h_2}{e_{10}} y\right) \\ &= x^{[0]} \prod_{k=1}^6 \left(1 + \frac{1}{qe_k} x\right) + x^{[6]} y + x^{[3]} \prod_{k=7}^9 \left(1 + \frac{e_k}{h_1} x\right) y^2 + x^{[0]} \prod_{k=7}^9 \prod_{t=0}^1 \left(1 + q^t \frac{e_k}{h_1} x\right) y^3, \quad (5.2) \end{aligned}$$

where $x^{[i]}$ (or $y^{[i]}$) represent some polynomials in x (or y) of degree i .

Proposition 5.1. *When the parameters satisfy the constraint*

$$h_1^6 h_2^3 = e_1 e_2 e_3 e_4 e_5 e_6 e_7^2 e_8^2 e_9^2 e_{10}^3 e_{11}^3, \quad (5.3)$$

then the general solution $P(x, y)$ of the condition (5.2) takes the form

$$P(x, y) = c_0 P_0(x, y) + c_1 x^3 y. \quad (5.4)$$

Moreover, when the polynomial $P(x, y)$ is normalized as $P(0, 0) = 1$ and $c_1 \in \mathbb{C}$, then $P(x, y)$ is invariant under the action of $W(E_8^{(1)})$ up to some multiplicative factors, namely

$$s_i(P) = P \quad (i \neq 0, 3), \quad s_0(P) = P \prod_{i=0}^2 \frac{1 + q^i \frac{h_2}{e_{10}} y}{1 + q^i e_{11} y}, \quad s_3(P) = P \frac{1 + \frac{e_7}{q h_1} x}{1 + \frac{1}{q e_1} x}. \quad (5.5)$$

Proof. Consider an auxiliary case

$$\mu = ((d_1, d_2), (m_1, \dots, m_{11})) = ((6, 3), (0, 1, 1, 1, 1, 1, 2, 2, 2, 3, 3)).$$

From equation (4.2), the general solution $F_\mu(x, y)$ for the condition $(x)_\mu, (y)_\mu$ has two linearly independent solutions. We can and we will choose a basis $\{P_0(x, y), x^3 y\}$, where $P_0(x, y)$ is fixed by the conditions: (i) the coefficient of $x^3 y$ in $P_0(x, y)$ is zero, and (ii) $P_0(0, 0) = 1$. By definition $F_\mu(x, 0)$ has 6 roots at $x = a, e_2, \dots, e_6$, where a is determined by $h_1^6 h_2^3 = a e_2 e_3 e_4 e_5 e_6 e_7^2 e_8^2 e_9^2 e_{10}^3 e_{11}^3$ due to the relation between the roots and the coefficients. Now we turn to the case λ in equation (5.1). Compared with the case μ , the case λ demands one more condition $P(e_1, 0) = 0$. However this extra condition is automatically satisfied if $a = e_1$, i.e., the constraint (5.3) is satisfied. Hence, under the constraint (5.3), the general solution $F_\lambda(x, y)$ is given by $F_\mu(x, y)|_{a=e_1}$ which has the desired form (5.4). Equation (5.5) follows from Theorem 4.4 and explicit computation on the monomial $x^3 y$. ■

Corollary 5.2. *The quotient $H(x, y) = P(x, y)x^{-3}y^{-1}$ is invariant under $W(E_8^{(1)})$.*

The quantum curve $\hat{H}(Q, P; \{f_i, g_i, h_i\})$ for E_8 in [41] written in the “rectangular” realization coincides with $H(x, y; \{h_i, e_i\}) = x^{-3}P_0(x, q^{-3}y)y^{-1}$ up to a normalization, by the following change of the variables and parameters,⁷

$$\begin{aligned} (Q, P) &\rightarrow q^{\frac{1}{2}}(x^{-1}, y), & (f_1, f_2, f_3) &\rightarrow h_1^{-1}(e_7, e_8, e_9), & (g_1, g_2) &\rightarrow \left(\frac{1}{e_{11}}, \frac{e_{10}}{h_2}\right), \\ (h_1, \dots, h_6) &\rightarrow e_{11}(e_1, \dots, e_6). \end{aligned}$$

The corresponding tropical curve is a pencil of elliptic curves given in Figure 2.

An explicit form of the polynomial $P_0(x, y)$ is given by

$$\begin{aligned} P_0(x, y) &= \sum_{i=0}^3 C_i(x) y^i, \\ C_3(x) &= q^3 e_{11}^3 \prod_{i=7}^9 \left(1 + \frac{e_i}{h_1} x\right) \left(1 + q \frac{e_i}{h_1} x\right), \\ C_2(x) &= q e_{11}^2 \prod_{i=7}^9 \left(1 + \frac{e_i}{h_1} x\right) \{[3]_q + q x A_{-1} + q \kappa A_1 x^2 + [3]_q \kappa x^3\}, \end{aligned}$$

⁷Note that the symbols h_i have different meanings in $\hat{H}(Q, P; \{f_i, g_i, h_i\})$ [41] and $H(x, y; \{h_i, e_i\})$ here.

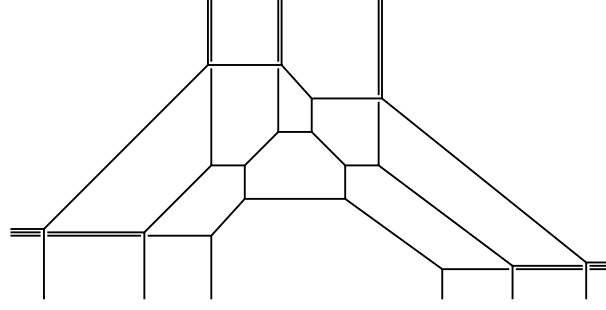


Figure 2. The web diagram corresponding to the curve $H(x, y) = E$ which has 6(single) + 3(double) + 2(triple) asymptotic lines. It has one closed cycle whose size depends on the parameter E , hence (genus) = 1. See, e.g., [3, 32, 33].

$$C_1(x) = e_{11} \left\{ [3]_q + [2]_q A_{-1} x + (\kappa A_1 + A_{-2}) x^2 + \frac{\kappa}{q} (\kappa A_2 + A_{-1}) x^4 \right. \\ \left. + \frac{[2]_q \kappa^2 A_1}{q^2} x^5 + \frac{[3]_q \kappa^2}{q^3} x^6 \right\},$$

$$C_0(x) = \prod_{i=1}^6 \left(1 + \frac{1}{q e_i} x \right),$$

where

$$A_{\pm 1} = \sum_{i=1}^9 a_i^{\pm 1}, \quad A_{\pm 2} = \sum_{1 \leq i < j \leq 9} (a_i a_j)^{\pm 1},$$

$$a_i = e_i \quad (1 \leq i \leq 6), \quad a_i = \frac{h_1}{e_i} \quad (7 \leq i \leq 9),$$

$$[k]_q = \frac{1 - q^k}{1 - q}, \quad \kappa = \frac{e_7 e_8 e_9 e_{10} e_{11}}{h_1^2 h_2}.$$

Remark 5.3. In the context of discrete Painlevé equations, the constraint (5.3) gives the autonomous case where the system admits a conserved curve $H(x, y) = E$.

6 Bilinear equations

The Weyl group representation in Theorem 2.3 can be reformulated as follows.

Proposition 6.1. *Introduce variables $\tau_{1,10}, \tau_{1,11}, \tau_{2,1}, \tau_{2,7}$ instead of x, y, σ_1, σ_2 as*

$$\tau_{1,10} = \frac{\sigma_1}{\tau_{10}}, \quad \tau_{1,11} = x \frac{\sigma_1}{\tau_{11}}, \quad \tau_{2,1} = y \frac{\sigma_2}{\tau_1}, \quad \tau_{2,7} = \frac{\sigma_2}{\tau_7}.$$

Then we have the representation of the Weyl group $W(E_8^{(1)})$ given by

$$s_0 = \left\{ e_{10} \rightarrow \frac{h_2}{e_{11}}, e_{11} \rightarrow \frac{h_2}{e_{10}}, h_1 \rightarrow \frac{h_1 h_2}{e_{10} e_{11}}, \right. \\ \left. \tau_{10} \rightarrow (\tau_{2,7} \tau_7 + e_{11} \tau_{2,1} \tau_1) \frac{1}{\tau_{11}}, \tau_{11} \rightarrow \frac{1}{\tau_{10}} \left(\tau_{2,7} \tau_7 + \frac{h_2}{e_{10}} \tau_{2,1} \tau_1 \right) \right\},$$

$$s_1 = \{e_8 \leftrightarrow e_9, \tau_8 \leftrightarrow \tau_9\}, \quad s_2 = \left\{ e_7 \leftrightarrow e_8, \tau_7 \leftrightarrow \tau_8, \tau_{2,7} \rightarrow \frac{\tau_7 \tau_{2,7}}{\tau_8} \right\},$$

$$\begin{aligned}
s_3 &= \left\{ e_1 \rightarrow \frac{h_1}{e_7}, e_7 \rightarrow \frac{h_1}{e_1}, h_2 \rightarrow \frac{h_1 h_2}{e_1 e_7}, \right. \\
&\quad \left. \tau_1 \rightarrow \left(\tau_{1,10} \tau_{10} + \frac{e_7}{h_1} \tau_{1,11} \tau_{11} \right) \frac{1}{\tau_7}, \tau_7 \rightarrow \frac{1}{\tau_1} \left(\tau_{1,10} \tau_{10} + \frac{1}{e_1} \tau_{1,11} \tau_{11} \right) \right\}, \\
s_4 &= \left\{ e_1 \leftrightarrow e_2, \tau_1 \leftrightarrow \tau_2, \tau_{2,1} \rightarrow \frac{\tau_1 \tau_{2,1}}{\tau_2} \right\}, \quad s_5 = \{e_2 \leftrightarrow e_3, \tau_2 \leftrightarrow \tau_3\}, \\
s_6 &= \{e_3 \leftrightarrow e_4, \tau_3 \leftrightarrow \tau_4\}, \quad s_7 = \{e_4 \leftrightarrow e_5, \tau_4 \leftrightarrow \tau_5\}, \quad s_8 = \{e_5 \leftrightarrow e_6, \tau_5 \leftrightarrow \tau_6\}.
\end{aligned}$$

Proof. The actions written in the new variables are computed as follows.

$$\begin{aligned}
\tau_{10} &\xrightarrow{s_0} (1 + y e_{11}) \frac{\sigma_2}{\tau_{11}} = \left(\frac{\sigma_2}{\tau_7} \tau_7 + e_{11} \frac{\sigma_2 y}{\tau_1} \tau_1 \right) \frac{1}{\tau_{11}} = (\tau_{2,7} \tau_7 + e_{11} \tau_{2,1} \tau_1) \frac{1}{\tau_{11}}, \\
\tau_{11} &\xrightarrow{s_0} \frac{\sigma_2}{\tau_{10}} \left(1 + y \frac{h_2}{e_{10}} \right) = \frac{1}{\tau_{10}} \left(\frac{\sigma_2}{\tau_7} \tau_7 + \frac{h_2}{e_{10}} \frac{\sigma_2 y}{\tau_1} \tau_1 \right) = \frac{1}{\tau_{10}} \left(\tau_{2,7} \tau_7 + \frac{h_2}{e_{10}} \tau_{2,1} \tau_1 \right), \\
\tau_1 &\xrightarrow{s_3} \left(1 + x \frac{e_7}{h_1} \right) \frac{\sigma_1}{\tau_7} = \left(\frac{\sigma_1}{\tau_{10}} \tau_{10} + \frac{e_7}{h_1} \frac{\sigma_1 x}{\tau_{11}} \tau_{11} \right) \frac{1}{\tau_7} = \left(\tau_{1,10} \tau_{10} + \frac{e_7}{h_1} \tau_{1,11} \tau_{11} \right) \frac{1}{\tau_7}, \\
\tau_7 &\xrightarrow{s_3} \frac{\sigma_1}{\tau_1} \left(1 + \frac{x}{e_1} \right) = \frac{1}{\tau_1} \left(\frac{\sigma_1}{\tau_{10}} \tau_{10} + \frac{1}{e_1} \frac{\sigma_1 x}{\tau_{11}} \tau_{11} \right) = \frac{1}{\tau_1} \left(\tau_{1,10} \tau_{10} + \frac{1}{e_1} \tau_{1,11} \tau_{11} \right).
\end{aligned}$$

Other actions are obvious. ■

In order to describe the bilinear equations in the Weyl-group covariant way, we define the tau functions $\tau(\lambda)$ on a certain lattice L as follows.

- (i) For $\lambda \in L_0 = \{e_1, \dots, e_{11}, \frac{h_2}{e_1}, \frac{h_2}{e_7}, \frac{h_1}{e_{10}}, \frac{h_1}{e_{11}}\}$, we put $\tau(e_i) = \tau_i$ ($1 \leq i \leq 11$), $\tau(\frac{h_2}{e_i}) = \tau_{2,i}$ ($i = 1, 7$) and $\tau(\frac{h_1}{e_j}) = \tau_{1,j}$ ($j = 10, 11$).
- (ii) For Weyl-group elements $w \in W(E_8^{(1)})$, we put $\tau(w(\lambda)) = w(\tau(\lambda))$.

From (i) and (ii), one can uniquely determine the functions $\tau(\lambda)$ for any $\lambda = \frac{h_1^{d_1} h_2^{d_2}}{e_1^{m_1} \dots e_{11}^{m_{11}}} \in L$, where L is the Weyl-group orbit of L_0 . For $\lambda \in \Lambda$, this fact is a consequence of Theorem 4.4, and it can be extended for $\lambda \in L$ similarly by using the normalization condition

$$\lim_{x \rightarrow 0} x^{-1} F_\lambda(x, y = 0) = 1, \quad \text{or} \quad \lim_{y \rightarrow 0} y^{-1} F_\lambda(x = 0, y) = 1,$$

for $\lambda = w(\frac{h_1}{e_{11}})$ or $\lambda = w(\frac{h_2}{e_1})$, respectively. (For the other cases we still have $F_\lambda(0, 0) = 1$.)

Corollary 6.2. *The functions $\tau(\lambda)$ satisfy the following relations:*

$$\begin{aligned}
\tau(e_{10}) \tau\left(\frac{h_2}{e_{10}}\right) &= \frac{h_2}{e_{10}} \tau\left(\frac{h_2}{e_i}\right) \tau(e_i) + \tau\left(\frac{h_2}{e_j}\right) \tau(e_j), \\
\tau\left(\frac{h_2}{e_{11}}\right) \tau(e_{11}) &= e_{11} \tau\left(\frac{h_2}{e_i}\right) \tau(e_i) + \tau\left(\frac{h_2}{e_j}\right) \tau(e_j), \\
\tau(e_i) \tau\left(\frac{h_1}{e_i}\right) &= \frac{1}{e_i} \tau\left(\frac{h_1}{e_{11}}\right) \tau(e_{11}) + \tau\left(\frac{h_1}{e_{10}}\right) \tau(e_{10}), \\
\tau\left(\frac{h_1}{e_j}\right) \tau(e_j) &= \frac{e_j}{h_1} \tau\left(\frac{h_1}{e_{11}}\right) \tau(e_{11}) + \tau\left(\frac{h_1}{e_{10}}\right) \tau(e_{10}), \\
\tau\left(\frac{h_2}{e_1}\right) \tau(e_1) &= \dots = \tau\left(\frac{h_2}{e_6}\right) \tau(e_6), \quad \tau\left(\frac{h_2}{e_7}\right) \tau(e_7) = \dots = \tau\left(\frac{h_2}{e_9}\right) \tau(e_9),
\end{aligned}$$

where $1 \leq i \leq 6$ and $7 \leq j \leq 9$. Furthermore, the infinitely many bilinear relations obtained from the above equations via the Weyl group actions also hold.

Proof. This is a simple reformulation of Proposition 6.1. ■

Example 6.3. The s_0 transform of the fourth equation in equation (6.2) with $j = 7$ is

$$\tau\left(\frac{h_1 h_2}{e_7 e_{10} e_{11}}\right) \tau(e_7) = \frac{e_7 e_{10} e_{11}}{h_1 h_2} \tau\left(\frac{h_1}{e_{11}}\right) \tau\left(\frac{h_2}{e_{10}}\right) + \tau\left(\frac{h_1}{e_{10}}\right) \tau\left(\frac{h_2}{e_{11}}\right). \quad (6.1)$$

This can be confirmed by

$$\begin{aligned} \tau\left(\frac{h_1 h_2}{e_7 e_{10} e_{11}}\right) &= \left\{1 + e_{11} y + \frac{e_7 e_{10} e_{11}}{h_1 h_2} x \left(1 + \frac{h_2}{e_{10}} y\right)\right\} \frac{\sigma_1 \sigma_2}{\tau_7 \tau_{10} \tau_{11}}, & \tau(e_7) &= \tau_7, \\ \tau\left(\frac{h_1}{e_{11}}\right) &= x \frac{\sigma_1}{\tau_{11}}, & \tau\left(\frac{h_2}{e_{10}}\right) &= \left(1 + \frac{h_2}{e_{10}} y\right) \frac{\sigma_2}{\tau_{10}}, \\ \tau\left(\frac{h_1}{e_{10}}\right) &= \frac{\sigma_1}{\tau_{10}}, & \tau\left(\frac{h_2}{e_{11}}\right) &= (1 + e_{11} y) \frac{\sigma_2}{\tau_{11}}. \end{aligned}$$

Note that each term in equation (6.1) is a member of two-parameter family $\tau\left(\frac{h_1 h_2}{e_{10} e_{11}}\right)$ (see equation (4.4)) and should satisfy a relation among three of them.

So far, we have derived the bilinear relations as the identities satisfied by the functions $\tau(\lambda)$ defined by the Weyl group actions. Conversely, we can consider the relations as the infinite system of equations viewing $\tau(\lambda)$ ($\lambda \in L$) as infinite unknown variables. We call this over-determined system of equations the *quantum bilinear equations* (denoted by \mathcal{B}). Note that the bilinear system \mathcal{B} is “tropical” (or subtraction free) [51].

Theorem 6.4. *For any initial data $\tau(\lambda)$ ($\neq 0$) ($\lambda \in L_0$), there exists a unique solution for the system \mathcal{B} . And this solution gives the general solution.*

Proof. We already have a solution as given in Corollary 6.2. It has 15 free parameters $\tau_1, \dots, \tau_{11}, \sigma_1, \sigma_2, x, y$ which are enough to fit the 15 initial data $\tau(\lambda)$ ($\lambda \in L_0$). ■

Remark 6.5. In the commutative case ($q = 1$), the space of the solutions of the system \mathcal{B} is of dimension 2 (modulo rescaling of variables τ_i, σ_i) and can be identified with the Okamoto space with coordinates x, y . Hence the system \mathcal{B} can be considered as a quantum analog of the Plücker embedding of the Okamoto space [26]. It will be interesting if the system \mathcal{B} can be obtained from some infinite-dimensional quantum integrable hierarchies. In view of this, we note that one can eliminate variables $\tau\left(\frac{h_1}{e_i}\right)$ ($i = 10, 11$) from the third and fourth equations in Corollary 6.2 to derive the bilinear equations in the standard Hirota–Miwa form (see [48]) such as

$$\begin{aligned} \left(\frac{1}{e_1} - \frac{1}{e_2}\right) \tau(e_3) \tau\left(\frac{h_1}{e_3}\right) + \left(\frac{1}{e_2} - \frac{1}{e_3}\right) \tau(e_1) \tau\left(\frac{h_1}{e_1}\right) + \left(\frac{1}{e_3} - \frac{1}{e_1}\right) \tau(e_2) \tau\left(\frac{h_1}{e_2}\right) &= 0, \\ \left(\frac{1}{e_1} - \frac{1}{e_2}\right) \tau\left(\frac{h_1}{e_7}\right) \tau(e_7) + \left(\frac{1}{e_2} - \frac{e_7}{h_1}\right) \tau(e_1) \tau\left(\frac{h_1}{e_1}\right) + \left(\frac{e_7}{h_1} - \frac{1}{e_1}\right) \tau(e_2) \tau\left(\frac{h_1}{e_2}\right) &= 0. \end{aligned}$$

7 Summary and discussions

In this paper, we studied the quantization of the affine Weyl group of type $E_8^{(1)}$ and obtained the following several results:

- A quantum (non-commutative) version of the affine Weyl group representation is formulated (Theorem 2.3).

- Its realization as adjoint actions is obtained (Theorem 2.7).
- Fundamental polynomials arising from the representation are studied and its characterization is given (Theorem 4.4).
- The quantum curve for $E_8^{(1)}$ in [41] is rederived by the Weyl group symmetry (Section 5).
- The quantum bilinear equations are obtained (Section 6).

Many of the results can be formulated similarly for the cases $D_5^{(1)}$, $E_6^{(1)}$ and $E_7^{(1)}$ as well. Such results are summarized in Appendix A.

One of our motivations for studying the quantum curves is the correspondence between spectral theories and topological strings, as observed in [18, 20, 21, 22]. Namely, the determinant of the spectral operator obtained from the quantum curve is described by the free energy of topological strings on the same geometry, which is captured by the period integrals. After providing the quantum curves and their origins in the affine Weyl groups, we believe that there are many directions to pursue to deepen the correspondence. Here we list some of future problems.

- Given a quantum curve, the study of the *spectral problem* is important. Since the expression is very huge in the exceptional cases $E_n^{(1)}$, the Weyl group symmetry will play a fundamental role to control them as discussed in [41]. It is interesting to start with the study of matrix elements of the spectral operators as in [30].
- After fixing the spectral operators, besides the spectral determinant, we can study various invariant or covariant quantities including the F -polynomials defined above. We believe that the correspondence is clarified from their relations.
- In relation to the spectral problem mentioned above, computation of the quantum period integrals is also an interesting problem [1, 2, 23, 24, 25, 39]. Even in genus one cases they are technical challenges in particular for the fully massive E_6 , E_7 , E_8 cases. Again, we expect that the Weyl groups serve an important role in studying the periods [15, 42].
- It is of course an interesting future direction to generalize our characterization to the cases of spectral operators of higher genus to study the correspondence in [11, 12].
- The tau functions for some classical (commutative) Painlevé equations have the general solutions in terms of the Nekrasov functions (the Kiev formula) in both differential and difference cases (see [16, 27] for example). Their extension to quantum Painlevé equations is an important problem [5, 6].
- In the context of Painlevé equations, the quantum curve appears in two ways: (i) as the conserved quantity for quantum autonomous Painlevé equations discussed in this paper, and (ii) as a certain specialization of the Lax linear equation for classical non-autonomous Painlevé equations [4, 45, 50]. In order to clarify the relation between the Kiev formula and the partition function of topological strings [18, 20], it will be important to understand the relation between (i) and (ii).
- There is a Lens generalization of the discrete Painlevé equation [31] whose identification in the Sakai's classification is not clear so far. It may be related to a quantization where q is a root of unity.
- The Weyl group symmetry (the iWeyl group) for the various (quantum) Seiberg–Witten curves was obtained (see [34, 44] for example). The Weyl-group actions considered in this paper are expected to be a realization of the iWeyl group.

A Cases $D_5^{(1)}$, $E_6^{(1)}$ and $E_7^{(1)}$

Here we will give the results for the cases $D_5^{(1)}$, $E_6^{(1)}$ and $E_7^{(1)}$ which are similar⁸ to the case $E_8^{(1)}$. First, we prepare some notations which are common for all cases.

To describe the Weyl group actions, we put

$$\begin{aligned} s_{i,j} &= \{e_i \leftrightarrow e_j, \tau_i \leftrightarrow \tau_j\}, \\ s_{i,j}^x &= \left\{ e_i \rightarrow \frac{h_2}{e_j}, e_j \rightarrow \frac{h_2}{e_i}, h_1 \rightarrow \frac{h_1 h_2}{e_i e_j}, x \rightarrow x \frac{1 + y \frac{h_2}{e_i}}{1 + y e_j}, \right. \\ &\quad \left. \tau_i \rightarrow (1 + y e_j) \frac{\sigma_2}{\tau_j}, \tau_j \rightarrow \frac{\sigma_2}{\tau_i} \left(1 + y \frac{h_2}{e_i}\right), \sigma_1 \rightarrow (1 + y e_j) \frac{\sigma_1 \sigma_2}{\tau_i \tau_j} \right\}, \\ s_{i,j}^y &= \left\{ e_i \rightarrow \frac{h_1}{e_j}, e_j \rightarrow \frac{h_1}{e_i}, h_2 \rightarrow \frac{h_1 h_2}{e_i e_j}, y \rightarrow \frac{1 + x \frac{e_j}{h_1}}{1 + \frac{x}{e_i}} y, \right. \\ &\quad \left. \tau_i \rightarrow \left(1 + x \frac{e_j}{h_1}\right) \frac{\sigma_1}{\tau_j}, \tau_j \rightarrow \frac{\sigma_1}{\tau_i} \left(1 + \frac{x}{e_i}\right), \sigma_2 \rightarrow \frac{\sigma_1 \sigma_2}{\tau_i \tau_j} \left(1 + \frac{x}{e_i}\right) \right\}. \end{aligned}$$

To specify the form of the F -polynomials for a given data $((d_1, d_2), (m_1, m_2, \dots))$, we put

$$F_{I,J}^x = \sum_{i=0}^{d_1} x^i \prod_{k \in J} \prod_{t=i}^{m_k-1} (1 + q^t e_k y) \prod_{k \in I} \prod_{t=d_1-m_k}^{i-1} \left(1 + q^t \frac{h_2}{e_k} y\right) U_i(y),$$

where $\deg U_i(y) = d_2 - \sum_{k \in I} (i - d_1 + m_k)_+ - \sum_{k \in J} (m_k - i)_+$, and

$$F_{I,J}^y = \sum_{i=0}^{d_2} \prod_{k \in I} \prod_{t=i-m_k}^{-1} \left(1 + q^t \frac{1}{e_k} x\right) \prod_{k \in J} \prod_{t=0}^{m_k-d_2+i-1} \left(1 + q^t \frac{e_k}{h_1} x\right) V_i(x) y^i,$$

where $\deg V_i(x) = d_1 - \sum_{k \in I} (m_k - i)_+ - \sum_{k \in J} (i - d_2 + m_k)_+$. With this notation, the previous result for the $E_8^{(1)}$ case is given by

$$\begin{aligned} s_0 &= s_{10,11}^x, & s_1 &= s_{8,9}, & s_2 &= s_{7,8}, & s_3 &= s_{1,7}^y, & s_4 &= s_{1,2}, \\ s_5 &= s_{2,3}, & s_6 &= s_{3,4}, & s_7 &= s_{4,5}, & s_8 &= s_{5,6}. \end{aligned}$$

Besides, the notation is applicable to all the other lower-rank cases. All these results are consistent with the quantum curves and the Weyl actions given in [41].

Case $E_7^{(1)}$:

- The Weyl group $W(E_7^{(1)})$ corresponding to the Dynkin diagram

$$\begin{array}{c} s_0 \\ | \\ s_1 - s_2 - s_3 - s_4 - s_5 - s_6 - s_7, \end{array}$$

can be realized as

$$\begin{aligned} s_0 &= s_{9,10}^x, & s_1 &= s_{7,8}, & s_2 &= s_{6,7}, & s_3 &= s_{5,6}, \\ s_4 &= s_{1,5}^y, & s_5 &= s_{1,2}, & s_6 &= s_{2,3}, & s_7 &= s_{3,4}. \end{aligned}$$

⁸For the cases $D_5^{(1)}$, $E_6^{(1)}$, $E_7^{(1)}$ one can extend the affine Weyl group by including the automorphisms of the Dynkin diagram. However we will not consider such extensions here.

- Defining conditions for the F -polynomials are given as follows. If we collect terms with the same power of x , the F -polynomials take the form $F_{\{9\},\{10\}}^x$, while if we collect terms with the same power of y , they take the form $F_{\{1,2,3,4\},\{5,6,7,8\}}^y$.
- Under the condition $\frac{h_1^4 h_2^2}{e_1 \cdots e_8 e_9^2 e_{10}^2} = 1$, we have the quantum curve

$$P_{E_7^{(1)}} = \prod_{i=1}^4 \left(1 + \frac{x}{qe_i}\right) + \left\{ e_{10}(1+q) + e_{10} \left(\sum_{i=5}^8 \frac{e_i}{h_1} + \sum_{i=1}^4 \frac{1}{e_i} \right) x + cx^2 \right. \\ \left. + \kappa x^3 \left(\sum_{i=5}^8 \frac{h_1}{e_i} + \sum_{i=1}^4 e_i \right) + \frac{\kappa}{q} (1+q)x^4 \right\} y + e_{10}^2 q \prod_{i=5}^8 \left(1 + \frac{e_i x}{h_1}\right) y^2,$$

where $\kappa = \frac{h_2}{qe_1 e_2 e_3 e_4 e_9}$ and $c \in \mathbb{C}$.

Case $E_6^{(1)}$:

- The Weyl group $W(E_6^{(1)})$ corresponding to the Dynkin diagram

$$\begin{array}{c} s_0 \\ | \\ s_6 \\ | \\ s_1 - s_2 - s_3 - s_4 - s_5, \end{array}$$

can be realized as

$$\begin{aligned} s_0 &= s_{8,9}, & s_1 &= s_{5,6}, & s_2 &= s_{4,5}, & s_3 &= s_{1,4}^y, \\ s_4 &= s_{1,2}, & s_5 &= s_{2,3}, & s_6 &= s_{7,8}^x. \end{aligned}$$

- The F -polynomials take respectively the form of $F_{\{7\},\{8,9\}}^x$ and $F_{\{1,2,3\},\{4,5,6\}}^y$ if we collect the same power of x and y .
- Under the condition $\frac{h_1^3 h_2^2}{e_1 \cdots e_6 e_7^2 e_8 e_9} = 1$, we have the quantum curve

$$P_{E_6^{(1)}} = \left\{ e_8 + e_9 + cx + \frac{h_2}{qe_1 e_2 e_3 e_7} \left(\sum_{i=4}^6 \frac{h_1}{e_i} + \sum_{i=1}^3 e_i \right) x^2 + \frac{h_2(1+q)}{q^2 e_1 e_2 e_3 e_7} x^3 \right\} y \\ + e_8 e_9 \prod_{i=4}^6 \left(1 + \frac{e_i x}{h_1}\right) y^2 + \prod_{i=1}^3 \left(1 + \frac{x}{qe_i}\right).$$

Case $D_5^{(1)}$:

- The Weyl group $W(D_5^{(1)})$ corresponding to the Dynkin diagram

$$\begin{array}{ccccc} s_0 & & s_4 & & \\ | & & | & & \\ s_1 & - & s_2 & - & s_3 - s_5, \end{array}$$

can be realized as

$$s_0 = s_{7,8}, \quad s_1 = s_{3,4}, \quad s_2 = s_{3,7}^y, \quad s_3 = s_{1,5}^x, \quad s_4 = s_{1,2}, \quad s_5 = s_{5,6}.$$

- The F -polynomials take respectively the form of $F_{\{1,2\},\{5,6\}}^x$ and $F_{\{7,8\},\{3,4\}}^y$ if we collect the same power of x and y .
- Under the condition $\frac{h_1^2 h_2^2}{e_1 \cdots e_8} = 1$, we have the quantum curve

$$P_{D_5^{(1)}} = \prod_{i=7}^8 \left(1 + \frac{x}{qe_i}\right) + \left(e_5 + e_6 + cx + \frac{(e_1 + e_2)h_2}{qe_1 e_2 e_7 e_8} x^2\right)y + e_5 e_6 \prod_{i=3}^4 \left(1 + \frac{e_i x}{h_1}\right)y^2.$$

B Standard realizations in commutative case

In Sakai's theory [49], the geometry relevant for the 2nd order discrete/continuous Painlevé equations are classified as in the following list:

elliptic	E_8	
multiplicative	$E_8 \rightarrow E_7 \rightarrow E_6 \rightarrow D_5 \rightarrow A_4 \rightarrow A_{2+1} \rightarrow A_{1+1} \rightarrow A_1 \rightarrow A_0$	$\nearrow A_1$
additive	$E_8 \rightarrow E_7 \rightarrow E_6 \rightarrow$	$D_4 \rightarrow A_3 \rightarrow A_{1+1} \rightarrow A_1 \rightarrow A_0$ $\searrow \quad \quad \quad \searrow$ $\quad \quad \quad A_2 \rightarrow A_1 \rightarrow A_0$

This list is the same as the degeneration scheme of the E -string. The classes of “elliptic”, “multiplicative” and “additive” mean the types of the difference equation and correspond to the gauge theories in 6D/5D/4D (see, e.g., [8, 40]). The cases in the box admit the continuous flows (of the original Painlevé equation), and the relation between their Hamiltonians and the $D = 4$, $SU(2)$ Seiberg–Witten curves was observed in [28]. Symbols A_n , D_n , E_n represent the types of the symmetry (affine in the Painlevé equations) and correspond to the (non-affine) flavor symmetry of the gauge theory.⁹

There are two standard ways to realize the above geometry, namely (i) nine-point blow-up of \mathbb{P}^2 or (ii) eight-point blow-up of $\mathbb{P}^1 \times \mathbb{P}^1$. In the most generic case, these points determine an elliptic curve and we have the elliptic Painlevé equation. Here we will give the multiplicative case in the realizations (i) and (ii) together with their relations.

(i) \mathbb{P}^2 -realization. Consider a parametrization of a point $p_3(u) = (x(u) : y(u) : 1) \in \mathbb{P}^2$

$$x(u) = u, \quad y(u) = \frac{\epsilon_0}{u} - u^2 \quad (u \in \mathbb{P}^1). \quad (\text{B.1})$$

The equations parametrize a cubic curve C_3 (with a node) given by

$$\varphi_3(x, y) = x^3 + xy - \epsilon_0 = 0.$$

The group structure of the curve C_3 is multiplicative, i.e., $3n$ points $p_3(u_i)$ ($i = 1, \dots, 3n$) are intersections of C_3 and a curve of degree n iff $u_1 \cdots u_{3n} = \epsilon_0^n$. Hence, the blow-up of \mathbb{P}^2 at the nine points $p_3(\epsilon_i)$ has the elliptic fibration iff $\epsilon_1 \cdots \epsilon_9 = \epsilon_0^3$.

(ii) $\mathbb{P}^1 \times \mathbb{P}^1$ -realization. Consider a parametrization of a point $p_{2,2}(u) = (f(u), g(u)) \in \mathbb{P}^1 \times \mathbb{P}^1$

$$f(u) = u + \frac{h_1}{u}, \quad g(u) = u + \frac{h_2}{u}. \quad (\text{B.2})$$

⁹Since the gauge theories are associated with the autonomous limit of the Painlevé equations, the affine Weyl groups are reduced to the finite Weyl groups.

The equations parametrize a bidegree (2,2) curve $C_{2,2}$ (with a node) given by

$$\varphi_{2,2}(f, g) = \frac{(f - g)(h_2 f - h_1 g)}{h_1 - h_2} + (h_1 - h_2) = 0.$$

This curve is also multiplicative; $N = 2(m + n)$ points $p_{2,2}(u_i)$ ($i = 1, \dots, N$) are intersections of $C_{2,2}$ and a curve of bidegree (m, n) iff $u_1 \cdots u_N = h_1^m h_2^n$. Hence, the blow-up of $\mathbb{P}^1 \times \mathbb{P}^1$ at the eight points $p_{2,2}(e_i)$ has the elliptic fibration iff $e_1 \cdots e_8 = h_1^2 h_2^2$.

Proposition B.1. *The realizations (i) and (ii) are equivalent through the following birational symplectic transformation of variables (x, y) and (f, g) with the identification of parameters $(\epsilon_0 \dots, \epsilon_9)$ and $(h_1, h_2, e_1, \dots, e_8)$ given by*

$$\begin{aligned} f &= \frac{\frac{\epsilon_0}{\epsilon_1} - \epsilon_1 x - y}{x - \epsilon_1}, & g &= \frac{\frac{\epsilon_0}{\epsilon_2} - \epsilon_2 x - y}{x - \epsilon_2}, \\ h_1 &= \frac{\epsilon_0}{\epsilon_1}, & h_2 &= \frac{\epsilon_0}{\epsilon_2}, & e_1 &= \frac{\epsilon_0}{\epsilon_1 \epsilon_2}, & e_i &= \epsilon_{i+1} \quad (i > 1). \end{aligned} \quad (\text{B.3})$$

Proof. The relation of the parameters $(\epsilon_0 \dots, \epsilon_9)$ and $(h_1, h_2, e_1, \dots, e_8)$ is invertible (it is a “linear” isomorphism written in multiplicative coordinates). Also, by a direct computation, we see that the transformation between (x, y) and (f, g) is birational with the indeterminate points $p_3(\epsilon_1), p_3(\epsilon_2) \in \mathbb{P}^2$ and $p_{2,2}(e_1) \in \mathbb{P}^1 \times \mathbb{P}^1$. It is easy to check the parameterizations (B.1), (B.2) and the curves $C_3, C_{2,2}$ are mapped to each other by the transformation (B.3). Since

$$\omega := \frac{dx \wedge dy}{\varphi_3(x, y)} = \frac{df \wedge dg}{\varphi_{2,2}(f, g)}, \quad (\text{B.4})$$

equation (B.3) gives a symplectic transformation w.r.t. this symplectic form. ■

Using the transformation (B.3), we can derive the actions of affine Weyl group $W(E_8^{(1)})$.

Proposition B.2. *There exists a unique birational symplectic representation of affine Weyl group $W(E_8^{(1)})$ with the following properties:*

(i) *In variables $(x, y, \epsilon_0 \dots, \epsilon_9)$, the action is given by*

$$\begin{aligned} s_0 &= \left\{ \epsilon_0 \rightarrow \frac{\epsilon_0^2}{\epsilon_1 \epsilon_2 \epsilon_3}, \epsilon_1 \rightarrow \frac{\epsilon_0}{\epsilon_2 \epsilon_3}, \epsilon_2 \rightarrow \frac{\epsilon_0}{\epsilon_3 \epsilon_1}, \epsilon_3 \rightarrow \frac{\epsilon_0}{\epsilon_1 \epsilon_2}, x \rightarrow \tilde{x}, y \rightarrow \tilde{y} \right\}, \\ s_i &= \{\epsilon_i \leftrightarrow \epsilon_{i+1}\} \quad (i = 1, \dots, 8), \end{aligned}$$

where \tilde{x} and \tilde{y} are certain rational functions of (x, y) .

(ii) *In variables $(f, g, h_1, h_2, v_1, \dots, v_8)$, the action is given by*

$$\begin{aligned} s_0 &= \{e_1 \leftrightarrow e_2\}, & s_i &= \{e_{i-1} \leftrightarrow e_i\} \quad (i = 3, \dots, 8), \\ s_1 &= \{h_1 \leftrightarrow h_2, f \leftrightarrow g\}, & s_2 &= \left\{ h_2 \rightarrow \frac{h_1 h_2}{e_1 e_2}, e_1 \rightarrow \frac{h_1}{e_2}, e_2 \rightarrow \frac{h_1}{e_1}, g \rightarrow \tilde{g} \right\}, \end{aligned}$$

where \tilde{g} is a certain rational function of (f, g) .

Proof. In the \mathbb{P}^2 realization, we have obvious symmetries $s_i = \{\epsilon_i \leftrightarrow \epsilon_{i+1}\}$ ($i = 1, \dots, 8$) which generate \mathfrak{S}_9 , and also in the $\mathbb{P}^1 \times \mathbb{P}^1$ realization we have $\mathfrak{S}_2 \times \mathfrak{S}_8 = \langle s_1 = \{h_1 \leftrightarrow h_2, f \leftrightarrow g\} \rangle \times \langle s_0 = \{e_1 \leftrightarrow e_2\}, s_i = \{e_{i-1} \leftrightarrow e_i\} (i = 3, \dots, 8) \rangle$ (see Figure 3). By mixing up the actions \mathfrak{S}_9 and $\mathfrak{S}_2 \times \mathfrak{S}_8$, one can obtain the full generators for $W(E_8^{(1)})$. The non-trivial actions s_0



Figure 3. \mathfrak{S}_9 (left) and $\mathfrak{S}_2 \times \mathfrak{S}_8$ (right) subgroups in $W(E_8^{(1)})$.

in (i) and s_2 in (ii) can be obtained from the obvious actions in opposite realization through the relation (B.3). The explicit forms of \tilde{x} , \tilde{y} , \tilde{g} can be determined by

$$s_0(x) = \frac{x\epsilon_0(\epsilon_0 - w)}{\epsilon_0^2 - \epsilon_1\epsilon_2\epsilon_3w}, \quad s_0(w) = w, \quad w = \frac{(x - \epsilon_1)(x - \epsilon_2)(x - \epsilon_3)}{\epsilon_0(x^3 + xy - \epsilon_0)},$$

and

$$s_2\left(\frac{g - (v_1 + \frac{h_2}{v_1})}{g - (v_2 + \frac{h_2}{v_2})}\right) = \frac{f - (v_2 + \frac{h_1}{v_2})}{f - (v_1 + \frac{h_1}{v_1})} \frac{g - (v_1 + \frac{h_2}{v_1})}{g - (v_2 + \frac{h_2}{v_2})}.$$

Thus, we obtain the desired results. ■

Remark B.3. Written in the coordinates (x, w) , the Weyl group representation of $W(E_8^{(1)})$ is the same as that in Section 1 up to a change of the parameters (note that w here corresponds to y in Section 1). In this coordinate, the symplectic form (B.4) takes a simple form $\omega = \frac{dx \wedge dw}{xw}$. This explains the reason why the realization in Section 1 is suitable for quantization.

In closing this appendix, we will give explicit forms of the pencil of the conserved elliptic curves.

In the realization (i) the conserved curve is given by

$$0 = \lambda\varphi_3 + m_1(-x^2\epsilon_0 - xy^2 + y\epsilon_0) - m_2x(xy - \epsilon_0) - m_3x^3 + m_4x^2 - m_5x \\ + m_6 - m_7(x^2 + y)\epsilon_0^{-1} + m_8(x^2y + x\epsilon_0 + y^2)\epsilon_0^{-2} - m_9(-3x^3\epsilon_0 + y^3 + 3\epsilon_0^2)\epsilon_0^{-3},$$

where $\sum_{i=0}^9 m_i z^i = \prod_{j=1}^9 (1 + \epsilon_j z)$ under the constraint $\epsilon_0^3 = \epsilon_1 \cdots \epsilon_9$.

In the realization (ii) the conserved curve is given by

$$0 = \lambda\varphi_{2,2} - m_1[fg(gh_1 - fh_2) + fh_2^2 - gh_1^2] + m_2 \frac{(gh_1 - fh_2)^2}{h_1 - h_2} - m_3(gh_1 - fh_2) \\ + m_4(h_1 - h_2) - m_5(f - g) + m_6 \frac{(f - g)^2}{h_1 - h_2} - m_7 \frac{fg(f - g) - fh_1 + gh_2}{h_1 h_2} \\ + (h_1 - h_2)((f^2 - 2h_1)(g^2 - 2h_2) - h_1^2 - h_2^2),$$

where $\sum_{i=0}^8 m_i z^i = \prod_{j=1}^8 (1 + v_j z)$ under the constraint $h_1^2 h_2^2 = v_1 \cdots v_8$.

Written in the Weierstrass form these curves coincide with the Seiberg–Witten curve for 5D E -string [13]. In the quantum case, we do not know whether such cubic or bi-quadratic form is available or not so far.

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References

- [1] Aganagic M., Cheng M.C.N., Dijkgraaf R., Krefl D., Vafa C., Quantum geometry of refined topological strings, *J. High Energy Phys.* **2012** (2012), no. 11, 019, 53 pages, [arXiv:1105.0630](#).
- [2] Aganagic M., Dijkgraaf R., Klemm A., Mariño M., Vafa C., Topological strings and integrable hierarchies, *Comm. Math. Phys.* **261** (2006), 451–516, [arXiv:hep-th/0312085](#).
- [3] Benini F., Benvenuti S., Tachikawa Y., Webs of five-branes and $\mathcal{N} = 2$ superconformal field theories, *J. High Energy Phys.* **2009** (2009), no. 9, 052, 33 pages, [arXiv:0906.0359](#).
- [4] Bershtein M., Gavrylenko P., Grassi A., Quantum spectral problems and isomonodromic deformations, [arXiv:2105.00985](#).
- [5] Bershtein M., Gavrylenko P., Marshakov A., Cluster integrable systems, q -Painlevé equations and their quantization, *J. High Energy Phys.* **2018** (2018), no. 2, 077, 34 pages, [arXiv:1711.02063](#).
- [6] Bershtein M.A., Shchepochkin A.I., q -deformed Painlevé τ function and q -deformed conformal blocks, *J. Phys. A: Math. Theor.* **50** (2017), 085202, 22 pages, [arXiv:1608.02566](#).
- [7] Bonelli G., Grassi A., Tanzini A., Quantum curves and q -deformed Painlevé equations, *Lett. Math. Phys.* **109** (2019), 1961–2001, [arXiv:1710.11603](#).
- [8] Bonelli G., Lisovyy O., Maruyoshi K., Sciarappa A., Tanzini A., On Painlevé/gauge theory correspondence, *Lett. Math. Phys.* **107** (2017), 2359–2413, [arXiv:1612.06235](#).
- [9] Chekhov L., Mazzocco M., Rubtsov V., Quantised Painlevé monodromy manifolds, Sklyanin and Calabi–Yau algebras, *Adv. Math.* **376** (2021), 107442, 52 pages, [arXiv:2020.10744](#).
- [10] Chen J., Haghighat B., Kim H.-C., Sperling M., Wang X., E -string quantum curve, [arXiv:2103.16996](#).
- [11] Codesido S., Grassi A., Mariño M., Spectral theory and mirror curves of higher genus, *Ann. Henri Poincaré* **18** (2017), 559–622, [arXiv:1507.02096](#).
- [12] Codesido S., Gu J., Mariño M., Operators and higher genus mirror curves, *J. High Energy Phys.* **2017** (2017), no. 2, 092, 53 pages, [arXiv:1609.00708](#).
- [13] Eguchi T., Sakai K., Seiberg–Witten curve for E -string theory revisited, *Adv. Theor. Math. Phys.* **7** (2003), 419–455, [arXiv:hep-th/0211213](#).
- [14] Fomin S., Zelevinsky A., The Laurent phenomenon, *Adv. in Appl. Math.* **28** (2002), 119–144, [arXiv:math.CO/0104241](#).
- [15] Furukawa T., Moriyama S., Sugimoto Y., Quantum mirror map for del Pezzo geometries, *J. Phys. A: Math. Theor.* **53** (2020), 385401, 18 pages, [arXiv:1908.11396](#).
- [16] Gamayun O., Iorgov N., Lisovyy O., Conformal field theory of Painlevé VI, *J. High Energy Phys.* **2012** (2012), no. 10, 038, 24 pages, Erratum, *J. High Energy Phys.* **2012** (2012), no. 10, 183, 1 page, [arXiv:1207.0787](#).
- [17] Grammaticos B., Nijhoff F.W., Ramani A., Discrete Painlevé equations, in *The Painlevé Property*, CRM Ser. Math. Phys., Springer, New York, 1999, 413–516.
- [18] Grassi A., Hatsuda Y., Mariño M., Topological strings from quantum mechanics, *Ann. Henri Poincaré* **17** (2016), 3177–3235, [arXiv:1410.3382](#).
- [19] Hasegawa K., Quantizing the Bäcklund transformations of Painlevé equations and the quantum discrete Painlevé VI equation, in *Exploring New Structures and Natural Constructions in Mathematical Physics*, Adv. Stud. Pure Math., Vol. 61, Math. Soc. Japan, Tokyo, 2011, 275–288, [arXiv:math.QA/0703036](#).
- [20] Hatsuda Y., Mariño M., Moriyama S., Okuyama K., Non-perturbative effects and the refined topological string, *J. High Energy Phys.* **2014** (2014), no. 9, 168, 42 pages, [arXiv:1306.1734](#).
- [21] Hatsuda Y., Moriyama S., Okuyama K., Instanton effects in ABJM theory from Fermi gas approach, *J. High Energy Phys.* **2013** (2013), no. 1, 158, 40 pages, [arXiv:1211.1251](#).
- [22] Hatsuda Y., Moriyama S., Okuyama K., Instanton bound states in ABJM theory, *J. High Energy Phys.* **2013** (2013), no. 5, 054, 22 pages, [arXiv:1301.5184](#).
- [23] Huang M.-X., On gauge theory and topological string in Nekrasov–Shatashvili limit, *J. High Energy Phys.* **2012** (2012), no. 6, 152, 40 pages, [arXiv:1205.3652](#).
- [24] Huang M.-X., Klemm A., Reuter J., Schiereck M., Quantum geometry of del Pezzo surfaces in the Nekrasov–Shatashvili limit, *J. High Energy Phys.* **2015** (2015), no. 2, 031, 52 pages, [arXiv:1401.4723](#).
- [25] Huang M.-X., Sugimoto Y., Wang X., Quantum periods and spectra in dimer models and Calabi–Yau geometries, *J. High Energy Phys.* **2020** (2020), no. 9, 168, 36 pages, [arXiv:2006.13482](#).

- [26] Jimbo M., Miwa T., Monodromy preserving deformation of linear ordinary differential equations with rational coefficients. II, *Phys. D* **2** (1981), 407–448.
- [27] Jimbo M., Nagoya H., Sakai H., CFT approach to the q -Painlevé VI equation, *J. Integrable Syst.* **2** (2017), xyx009, 27 pages, [arXiv:1706.01940](#).
- [28] Kajiwara K., Masuda T., Noumi M., Ohta Y., Yamada Y., Cubic pencils and Painlevé Hamiltonians, *Funkcial. Ekvac.* **48** (2005), 147–160, [arXiv:nlin.SI/0403009](#).
- [29] Kajiwara K., Noumi M., Yamada Y., Geometric aspects of Painlevé equations, *J. Phys. A: Math. Theor.* **50** (2017), 073001, 164 pages, [arXiv:1509.08186](#).
- [30] Kashaev R., Mariño M., Zakany S., Matrix models from operators and topological strings, 2, *Ann. Henri Poincaré* **17** (2016), 2741–2781, [arXiv:1505.02243](#).
- [31] Kels A.P., Yamazaki M., Lens generalisation of τ -functions for the elliptic discrete Painlevé equation, *Int. Math. Res. Not.* **2021** (2021), 110–151, [arXiv:1810.12103](#).
- [32] Kim S.-S., Taki M., Yagi F., Tao probing the end of the world, *Prog. Theor. Exp. Phys.* **2015** (2015), 083B02, 29 pages, [arXiv:1504.03672](#).
- [33] Kim S.-S., Yagi F., 5d E_n Seiberg–Witten curve via toric-like diagram, *J. High Energy Phys.* **2015** (2015), no. 6, 082, 62 pages, [arXiv:1411.7903](#).
- [34] Kimura T., Pestun V., Twisted reduction of quiver W-algebras, [arXiv:1905.03865](#).
- [35] Kirillov A.N., Dilogarithm identities, *Progr. Theoret. Phys. Suppl.* **118** (1995), 61–142, [arXiv:hep-th/9408113](#).
- [36] Kubo N., Moriyama S., Nosaka T., Symmetry breaking in quantum curves and super Chern–Simons matrix models, *J. High Energy Phys.* **2019** (2019), no. 1, 210, 29 pages, [arXiv:1811.06048](#).
- [37] Kuroki G., Quantum groups and quantization of Weyl group symmetries of Painlevé systems, in Exploring New Structures and Natural Constructions in Mathematical Physics, *Adv. Stud. Pure Math.*, Vol. 61, *Math. Soc. Japan*, Tokyo, 2011, 289–325, [arXiv:0808.2604](#).
- [38] Kuroki G., Regularity of quantum tau-functions generated by quantum birational Weyl group actions, [arXiv:1206.3419](#).
- [39] Mironov A., Morozov A., Nekrasov functions from exact Bohr–Sommerfeld periods: the case of $SU(N)$, *J. Phys. A: Math. Theor.* **43** (2010), 195401, 11 pages, [arXiv:0911.2396](#).
- [40] Mizoguchi S., Yamada Y., $W(E_{10})$ symmetry, M-theory and Painlevé equations, *Phys. Lett. B* **537** (2002), 130–140, [arXiv:hep-th/0202152](#).
- [41] Moriyama S., Spectral theories and topological strings on del Pezzo geometries, *J. High Energy Phys.* **2020** (2020), no. 10, 154, 53 pages, [arXiv:2007.05148](#).
- [42] Moriyama S., Nosaka T., Yano K., Superconformal Chern–Simons theories from del Pezzo geometries, *J. High Energy Phys.* **2017** (2017), no. 11, 089, 40 pages, [arXiv:1707.02420](#).
- [43] Nagao H., Yamada Y., Study of q -Garnier system by Padé method, *Funkcial. Ekvac.* **61** (2018), 109–133, [arXiv:1601.01099](#).
- [44] Nekrasov N., Pestun V., Seiberg–Witten geometry of four dimensional $\mathcal{N} = 2$ quiver gauge theories, [arXiv:1211.2240](#).
- [45] Noumi M., Ruijsenaars S., Yamada Y., The elliptic Painlevé Lax equation vs. van Diejen’s 8-coupling elliptic Hamiltonian, *SIGMA* **16** (2020), 063, 16 pages, [arXiv:1903.09738](#).
- [46] Noumi M., Yamada Y., Affine Weyl groups, discrete dynamical systems and Painlevé equations, *Comm. Math. Phys.* **199** (1998), 281–295, [arXiv:math.QA/9804132](#).
- [47] Noumi M., Yamada Y., Birational Weyl group action arising from a nilpotent Poisson algebra, in Physics and Combinatorics 1999 (Nagoya), *World Sci. Publ.*, River Edge, NJ, 2001, 287–319, [arXiv:math.QA/0012028](#).
- [48] Ohta Y., Ramani A., Grammaticos B., An affine Weyl group approach to the eight-parameter discrete Painlevé equation, *J. Phys. A: Math. Gen.* **34** (2001), 10523–10532.
- [49] Sakai H., Rational surfaces associated with affine root systems and geometry of the Painlevé equations, *Comm. Math. Phys.* **220** (2001), 165–229.
- [50] Takemura K., On q -deformations of the Heun equation, *SIGMA* **14** (2018), 061, 16 pages, [arXiv:1712.09564](#).
- [51] Tsuda T., Tropical Weyl group action via point configurations and τ -functions of the q -Painlevé equations, *Lett. Math. Phys.* **77** (2006), 21–30.
- [52] Yamada Y., An elliptic Garnier system from interpolation, *SIGMA* **13** (2017), 069, 8 pages, [arXiv:1706.05155](#).